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## Determining Optimal Discharge and Optimal Penstock Diameter in Water Turbines

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## Determining Optimal Discharge and Optimal Penstock Diameter in Water Turbines

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### ABSTRACT

*Minimizing water consumption for producing hydropower is critical given that overuse of flows for energy production may result in a shortage of flows for other purposes such as irrigation and navigation. This paper presents a dimensional analysis for finding optimal flow discharge and optimal penstock diameter when designing impulse and reaction water turbines for hydropower systems. The objective of this analysis is to provide general insights for minimizing water consumption when producing hydropower. This analysis is based on the geometric and hydraulic characteristics of the penstock, the total hydraulic head, and the desired power production. As part of this analysis, various dimensionless relationships between power production, flow discharge, and head losses were derived. These relationships were used to draw general insights on determining optimal flow discharge and optimal penstock diameter. For instance, it was found that for minimizing water consumption, the ratio of head loss to gross head should not exceed about 15%. An example of application is presented to illustrate the procedure for determining optimal flow discharge and optimal penstock diameter for an impulse turbine. It is worth mentioning that this paper presents part of the material published by the author in Leon and Zhu (2014).*

**Keywords:** Design, dimensional analysis, efficiency, hydropower, optimal power, turbine

### 1. INTRODUCTION

The world energy consumption will grow by 56% between 2010 and 2040 (U.S. Energy Information Administration 2013). As world population continues to grow and the limited amount of fossil fuels begins to diminish, there is an increasing demand to exploit renewable sources of energy. In the United States, about 9% of all energy consumed in 2012 was from renewable sources (U.S. Institute for Energy Research 2012). While this is a relatively small fraction of the U.S. energy supply, in 2012, the United States was the world's largest consumer of renewable energy from geothermal, solar, wood, wind, and waste for electric power generation, producing almost 25% of the world's total (U.S. Institute for Energy Research 2012). This institute also reports that in 2012, 30% of the renewable energy in the U.S. was from hydropower. This means that only about 3% of all energy consumed in the United States was from hydropower.

Globally, hydropower accounted for 16% of all global electricity production in 2007, with other renewable energy sources totaling 3% (Schumann et al. 2010). Hence, it is not surprising that when options are evaluated for new energy developments, there is a strong impulse toward fossil fuel or nuclear energy as opposed to renewable sources. However, as hydropower schemes are often part of a multipurpose water resources development project, they can often help to finance other important functions of the project (IEA Hydro, 2000). In addition, hydropower provides benefits that are rarely found in other sources of energy. In fact, dams built for hydropower schemes, and their associated reservoirs, provide human well-being benefits such as securing water supply, flood control, and irrigation for food production and societal benefits such as increased recreational activities and improved navigation (IEA Hydro, 2000).

Furthermore, hydropower, due to its associated reservoir storage, can provide flexibility and reliability for energy production in integrated energy systems. The storage capability of hydropower systems can be seen as a regulating mechanism by which other diffuse and variable renewable energy sources (wind, wave, solar) can play a larger role in providing electric power of commercial quality (Schumann et al. 2010). While development of all the remaining hydroelectric potential could not hope to cover total future world demand for electricity, implementation of the remaining potential can make a vast contribution to improving living standards in the developing world (South America, Asia and Africa), where the greatest potential still exists (U.S. Institute for Energy Research 2012).

Minimizing water consumption for producing hydropower is critical given that overuse of flows for energy production may result in a shortage of flows for other purposes such as irrigation or navigation (Leon and Zhu 2014). The present work was motivated when the author was unable to find in literature a theoretical framework for determining optimal flow discharge and optimal penstock diameter for the design of impulse and reaction turbines. Recently, Pelz (2011) provided a theoretical approach for determining the upper limit for hydropower gained by a water wheel or turbine per unit width in a rectangular open-channel. This is somewhat different for impulse and reaction turbines, as in the latter turbines, the flow in the penstock is pressurized.

This paper aims to provide general insights on determining optimal flows and optimal penstock diameters when designing impulse and reaction turbines for hydropower systems. This paper is divided as follows. First, dimensionless relationships between power production, flow discharge, and head losses are derived. Second, these relationships are used to draw general insights on determining optimal flow discharge and optimal penstock diameter. Third, an example of application for determining optimal flows when designing impulse turbines is presented. Finally, the key results are summarized in the conclusion.

## 2. DIMENSIONAL ANALYSIS FOR OPTIMAL FLOW DISCHARGE, OPTIMAL HEAD LOSSES, AND OPTIMAL POWER

The electric power,  $P$ , in Watts (W), can be determined by the following equation:

$$P = \eta \gamma Q (H_g - h_L) \quad (1)$$

where  $\gamma (= \rho g)$  is the specific weight of water in  $\text{kg}/(\text{m}^2 \text{s}^2)$ ,  $Q$  is flow discharge in  $\text{m}^3/\text{s}$ ,  $H_g$  is gross head in m,  $h_L$  is sum of head losses in m,  $\rho$  is water density in  $\text{kg}/\text{m}^3$ ,  $g$  is acceleration of gravity in  $\text{m}/\text{s}^2$ , and  $\eta$  is overall hydroelectric unit efficiency, which in turn is the product of turbine efficiency ( $\eta_t$ ) and generator efficiency ( $\eta_g$ ). In all derivations presented in this paper, it is assumed that  $\eta (= \eta_t \eta_g)$  is constant.

For an impulse turbine (see Figure 1), the sum of head losses can be written as

$$h_L = \frac{Q}{2gA_2^2} \left[ f \frac{L}{D_2} + \sum k_{1-2} + k_N \left( \frac{A_2}{A_N} \right)^2 \right] \quad (2)$$

where  $L$ ,  $D_2$ , and  $A_2$  are length, diameter, and cross-sectional area of penstock, respectively. In addition,  $f$  is friction factor,  $\sum k_{1-2}$  is the sum of local losses in penstock due to entrance, bends, penstock fittings and gates,  $A_N$  is nozzle area at its exit (section 3 in Figure 1), and  $k_N$  is nozzle head loss coefficient, which is given by (e.g., Brater and King 1976)

$$k_N = \frac{1}{C_V^2} - 1 \quad (3)$$

where  $C_V$  is nozzle velocity coefficient. According to Dixon (2005),  $C_V$  varies between 0.98 and 0.99 for a typical Pelton turbine nozzle.

For a reaction turbine (see Leon and Zhu 2014), the sum of head losses can be written as

$$h_L = \frac{Q}{2gA_2^2} \left[ f \frac{L}{D_2} + \sum k_{1-2} + k_N \left( \frac{A_2}{A_d} \right)^2 \right] \quad (4)$$

where  $A_d$  is the draft tube cross-sectional area at its outlet (section Leon and Zhu 2014). The expression inside the brackets in Eqs. (2) and (4) is dimensionless, and it is denoted herein as

$$C_L = \begin{cases} f \frac{L}{D_2} + \sum k_{1-2} + k_N \left( \frac{A_2}{A_N} \right)^2 & \text{for an impulse turbine} \\ f \frac{L}{D_2} + \sum k_{1-2} + \left( \frac{A_2}{A_d} \right)^2 & \text{for a reaction turbine} \end{cases} \quad (5)$$

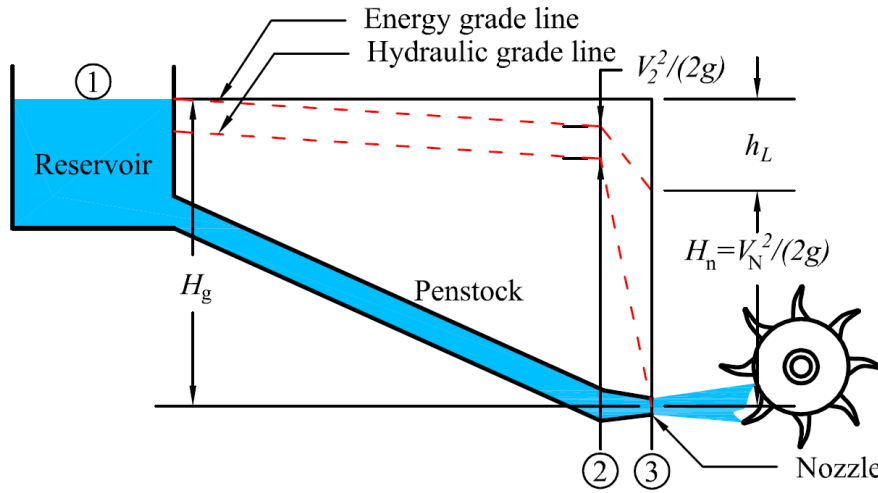


Figure 1. Sketch of an impulse turbine

Hence, the total head losses in Eq. (2) and Eq. (4) is equal to the product of  $C_L$  and  $Q^2/(2g(A_2)^2)$  and, thus, Eq. (1) can be written as

$$P = \eta \gamma Q \left( H_g - C_L \frac{Q^2}{2gA_2^2} \right) \quad (6)$$

For generalizing the findings in this paper, a dimensionless relationship between power and flow discharge is sought. To achieve this, Eq. (6) is divided by a reference power ( $P_r$ ).  $P_r$  is assumed to be the maximum power that can be generated using a reference discharge ( $Q_r$ ) and a fixed gross head and penstock geometry (constant  $C_L$ ). For maximum power, the turbine and generator efficiencies need to be 100% (i.e.,  $\eta_t = 100\%$  and  $\eta_g = 100\%$ ). Also, maximum power for a fixed penstock geometry can be obtained by setting  $dP/dQ$  in Eq. (6) equal to zero, which gives

$$h_L = H_g/3 \quad (7)$$

The reference flow discharge  $Q_r$  can be obtained by using Eq. (7) and the energy equation between the reservoir and the nozzle exit for an impulse turbine or between the reservoir and the tailrace for a reaction turbine, which gives

$$Q_r = 2A_3 \sqrt{\frac{1}{3}gH_g} \quad (8)$$

where  $A_3$  is given by

$$A_3 = \begin{cases} A_N & \text{for an impulse turbine} \\ A_d & \text{for a reaction turbine} \end{cases} \quad (9)$$

Substituting Eq. (7) and Eq. (8) into Eq. (1) gives the following relation for the reference power ( $P_r$ ):

$$P_r = \frac{4}{3}\gamma H_g A_3 \sqrt{\frac{1}{3}gH_g} \quad (10)$$

Note that  $Q_r$  and  $P_r$  (Eqs. 8 and 10) are a function of the penstock properties and the gross head only. Dividing each side of Eq. (6) by  $P_r$  (Eq. 10) and defining  $P/P_r$  as  $P_+$  and  $Q/Q_r$  as  $Q_+$ , and after some algebra, the following dimensionless relationship between power and discharge is obtained:

$$P_+ = \eta \left[ \frac{3}{2}Q_+ - C_L \left( \frac{A_3}{A_2} \right)^2 Q_+^3 \right] \quad (11)$$

Denoting with  $\beta$  the product of  $C_L$  and  $(A_3/A_2)^2$ , Eq. (11) can be rewritten as

$$P_+ = \eta \left( \frac{3}{2}Q_+ - \beta Q_+^3 \right) \quad (12)$$

where

$$\beta = \begin{cases} \left( \frac{A_N}{A_2} \right)^2 \left( f \frac{L}{D_2} + \sum k_{1-2} + k_N \left( \frac{A_2}{A_N} \right)^2 \right) & \text{for an impulse turbine} \\ \left( \frac{A_d}{A_2} \right)^2 \left( f \frac{L}{D_2} + \sum k_{1-2} + \left( \frac{A_2}{A_d} \right)^2 \right) & \text{for a reaction turbine} \end{cases} \quad (13)$$

In practice, the ratios  $A_N/A_2$  and  $A_d/A_2$  in Eq. (13) are typically kept constant, which means that  $\beta$  varies as a function of  $f$ ,  $L$ ,  $D_2$ , and the coefficients of local head losses ( $\Sigma k$ ). In most applications, friction losses are more important than local head losses, that is  $fL/D_2 \gg \Sigma k$ . Also,  $L$  is typically constant as it is restricted by topographic conditions. In addition,  $f$  does not show significant variation as a function of discharge or penstock diameter. Recall that for a given penstock diameter,  $f$  is independent of the Reynolds number for fully developed turbulent flows, which is the case of most penstock flows. Hence,  $\beta$  is more or less inversely proportional to the penstock diameter. The variation of  $P_+$  with respect to  $Q_+$  for a fixed  $\beta$  can be obtained by differentiating  $P_+$  with respect to  $Q_+$  in Eq. (12), which gives

$$\frac{dP_+}{dQ_+} = \eta \left( \frac{3}{2} - 3\beta Q_+^2 \right) \quad (14)$$

The maximum dimensionless power for a fixed  $\beta$  can be obtained by setting  $dP_+/dQ_+$  in Eq. (14) equal to zero. The maximum power occurs when

$$(Q_+)_{\max} = \sqrt{\frac{1}{2\beta}} \quad (15)$$

The maximum dimensionless power for a fixed  $\beta$  is obtained by substituting  $Q_+$  from Eq. (15) in Eq. (12), which gives

$$(P_+)_{\max} = \eta \sqrt{\frac{1}{2\beta}} \quad (16)$$

In most applications,  $\beta$  should range between 0.01 and 1.0 for impulse turbines and between 10 and 1000 for reaction turbines. Likewise,  $C_L$  should range between 1 and 100 for both impulse and reaction turbines. Even though  $\beta$  is used throughout the entire paper, as shown later, only  $C_L$  is needed for design purposes. Figures 2 and 3 plot  $Q_+$  versus  $P_+$  in Eq. (12) for typical ranges of  $\beta$  for impulse and reaction turbines, respectively. An overall hydroelectric unit efficiency ( $\eta$ ) of 0.8 was used for plotting these figures. As can be observed in Figures 2 and 3, the change in power production in relation to change in flow discharge ( $\Delta P_+/\Delta Q_+$ ) for each dimensionless curve has a positive and negative gradient. For optimizing power production, only the positive gradient is of interest ( $\Delta P_+/\Delta Q_+ > 0$ ). To visualize changes in power production in relation to changes in flow discharge, five ratios of  $dP_+/dQ_+$  in Eq. (14) are plotted in Figures 2 and 3. Note in Figures 2 and 3 that for a given  $\beta$ , the positive range of  $dP_+/dQ_+$  varies from  $(3/2)\eta$  to 0. Note also that  $dP_+/dQ_+$  changes rapidly near  $(Q_+)_{\max}$  and that in the positive range of  $dP_+/dQ_+$ , the maximum relative power  $P_+$  occurs for the maximum relative flow discharge  $Q_+$ .

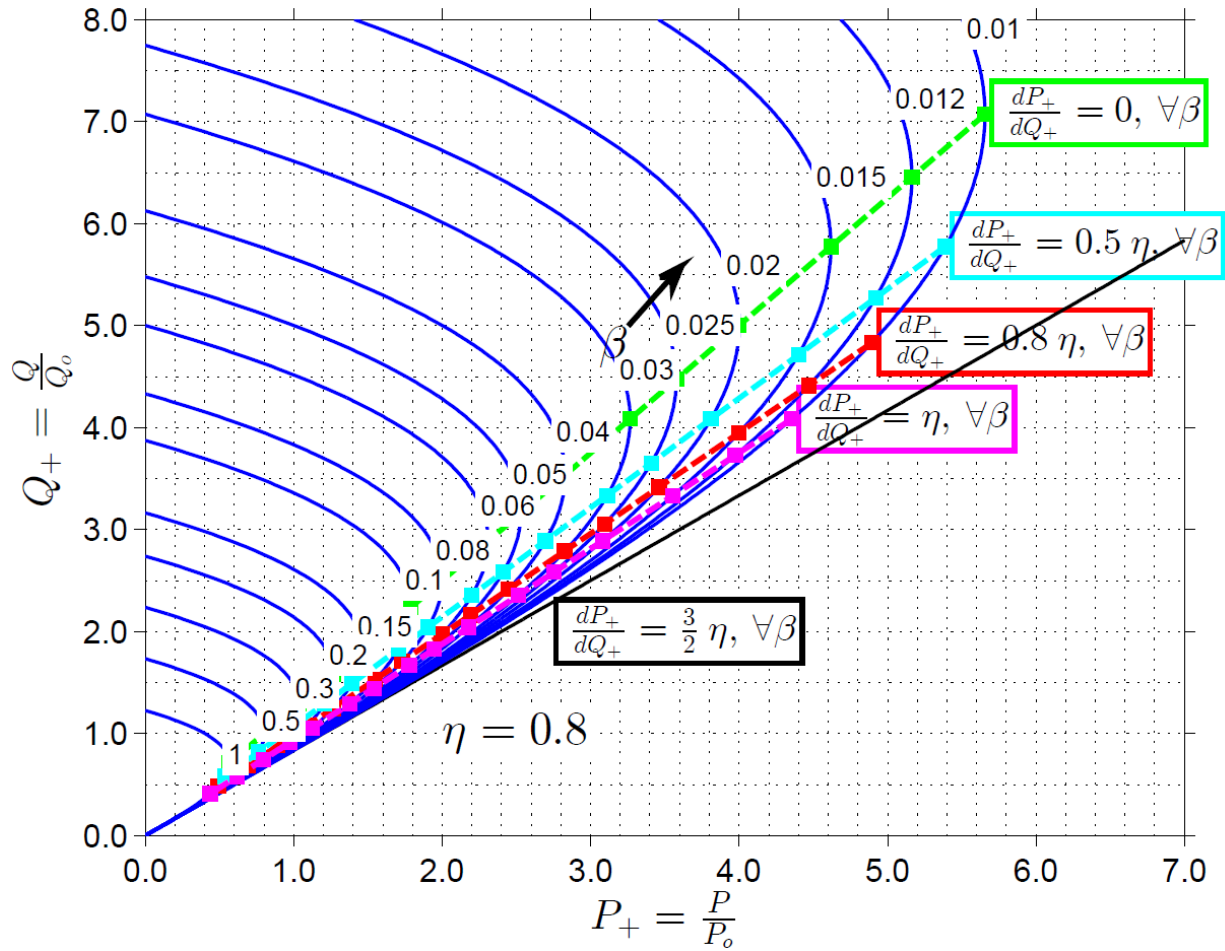


Figure 2. Dimensionless discharge ( $Q_+$ ) versus dimensionless power ( $P_+$ ) for  $\eta = 0.8$  and a typical range of  $\beta$  for impulse turbines

For minimizing water consumption to produce a given amount of hydropower, it is necessary that  $dP_+/dQ_+$  in Eq. (14) is close to its maximum value  $(3/2)\eta$ . Note in Figures 2 and 3 that for each curve between approximately  $dP_+/dQ_+ = (3/2)\eta$  and  $dP_+/dQ_+ = 0.8\eta$ , the increase in dimensionless power ( $P_+$ ) is approximately linear with increase in dimensionless discharge ( $Q_+$ ). Note also in these figures that for  $dP_+/dQ_+$  smaller than about  $0.8\eta$ , the increase in  $P_+$  is small compared to the increase in  $Q_+$ . Herein, to minimize water consumption, the optimal lower

limit of  $dP_+/dQ_+$  is set to  $0.8\eta$ . Substituting  $dP_+/dQ_+ = 0.8\eta$  into Eq. (14) gives the following upper limit for the dimensionless flow discharge:

$$(Q_+)_{\text{opt upper}} = \sqrt{\frac{7}{30\beta}} \quad (17)$$

The corresponding upper limit for the dimensionless power is

$$(P_+)_{\text{opt upper}} = \eta \frac{19}{15} \sqrt{\frac{7}{30\beta}} \quad (18)$$

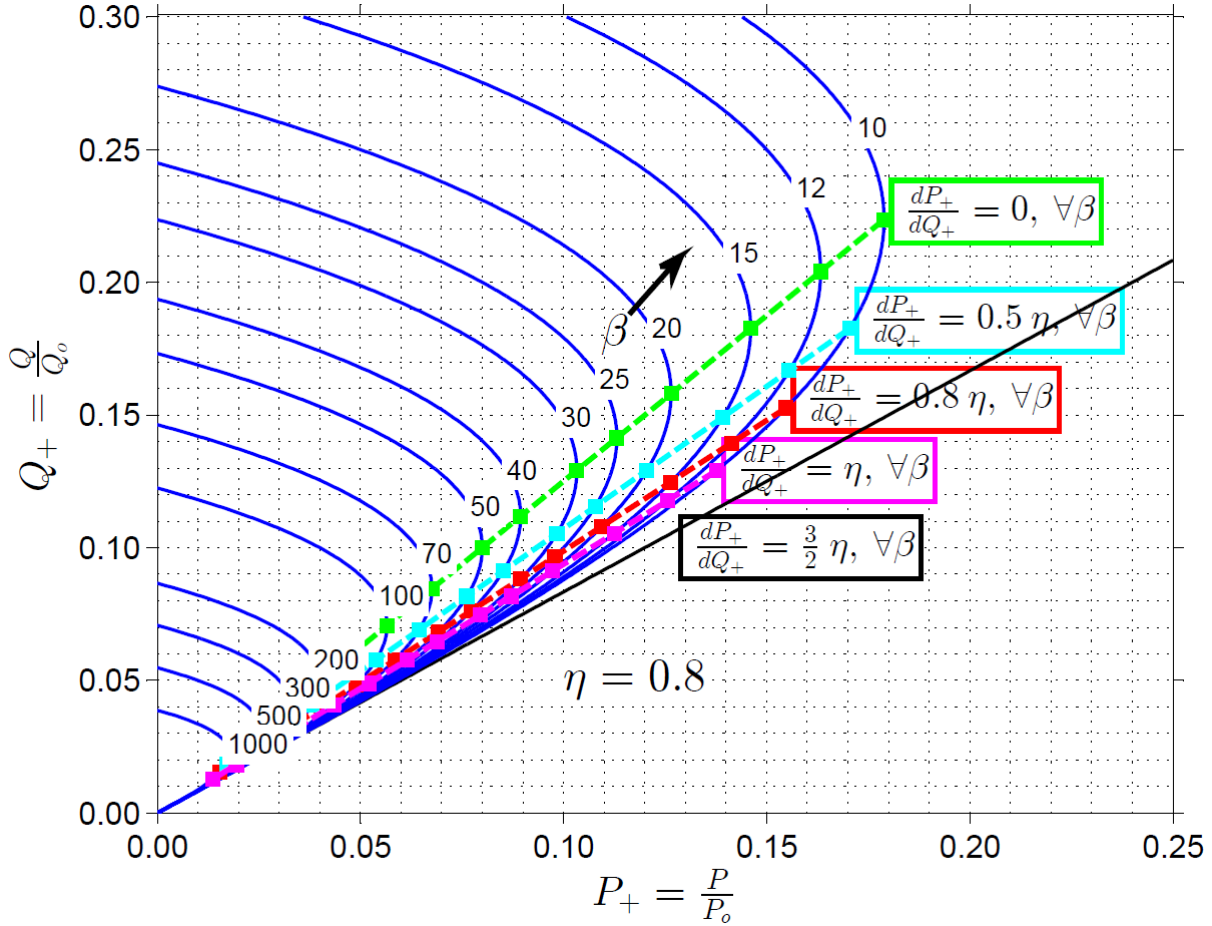


Figure 3. Dimensionless discharge ( $Q_+$ ) versus dimensionless power ( $P_+$ ) for  $\eta = 0.8$  and a typical range of  $\beta$  for reaction turbines

The optimal dimensionless head loss ( $h_{L+} = h_L/H_g$ ) can be obtained by assuming that the optimal upper limit for the flow discharge is  $Q_+ = [7/(30\beta)]^{1/2}$  (Eq. 17). In Eq. (12), dividing the second term of the right-hand side (RHS) by the first term of the RHS gives

$$h_{L+} \leq \frac{2}{3}\beta Q_+^2 \quad (19)$$

Substituting  $(Q_+)_{\text{opt upper}} = [7/(30\beta)]^{1/2}$  into Eq. (19) gives

$$h_{L+} \leq \frac{7}{45} \quad (20)$$

Eq. (20) shows that for minimizing water consumption, the ratio of head loss to gross head ( $h_{L+} = h_L/H_g$ ) should not exceed 15.6%. The 15.6% ratio also provides the threshold for the optimal penstock diameter. Losses higher than 15.6% mean that a small penstock diameter is used. The 15.6% ratio is about half of that derived for maximum power and maximum flow discharge, which is 33.3%. This means that the optimal conditions for producing power do not correspond to those that use maximum flow discharge for a given  $\beta$ . This can be better understood by observing Figures 2 and 3, in which  $dP_+/dQ_+$  decreases rapidly near  $(P_+)_{\max}$  for all  $\beta$ . So far the analysis assumed that  $\beta$  is constant, and, hence, the penstock diameter ( $D_2$ ). For the influence of changing the penstock diameter on power production, the reader is referred to Leon and Zhu (2014).

For practical applications, the derived dimensionless relationships are made non-dimensionless. For instance, the optimal upper limit of the flow discharge can be obtained by combining Eqs. (8) and (17), which, after some algebra, gives

$$Q_{\text{opt}} = \frac{2}{3} A_2 \sqrt{\frac{7}{10} \frac{g H_g}{C_L}} \quad (21)$$

Similarly, the optimal upper limit of the power can be obtained by combining Eqs. (10) and (18), which, after some algebra, gives

$$P_{\text{opt}} = \frac{76}{135} \eta \gamma H_g A_2 \sqrt{\frac{7}{10} \frac{g H_g}{C_L}} \quad (22)$$

When designing a turbine, it is necessary to specify either the flow discharge to use or the desired electric power. These cases are presented below.

## 2.1. $P$ is Specified

If  $P$  is specified, the optimal upper limit of the flow discharge can be obtained by combining Eqs. (21) and (22), which gives

$$Q_{\text{opt}} = \frac{45}{38} \left( \frac{P}{\eta \gamma H_g} \right) \quad (23)$$

The optimal penstock diameter (or  $C_L$ ) can be determined from Eq. (21), which gives

$$\frac{(C_L)_{\text{opt}}}{A_2^2} \leq \frac{14}{45} \frac{g H_g}{Q^2} \quad (24)$$

where  $Q$  in Eq. (24) is the same as  $Q_{\text{opt}}$  in Eq. (23).

## 2.2. $Q$ is Specified

If  $Q$  is specified, the optimal upper limit of the power can be obtained by combining Eqs. (21) and (22), which gives

$$P_{\text{opt}} = \frac{38}{45} \eta \gamma H_g Q \quad (25)$$

In this case, the optimal penstock diameter can still be determined using Eq. (24). It is pointed out that the proposed methodology for determining the optimal flow discharge and optimal penstock diameter does not account for



cavitation. Reaction turbines (not impulse turbines) are subjected to cavitation. In reaction turbines, cavitation may occur at the outlet of the runner or at the inlet of the draft tube where the pressure is considerably reduced (Dixon 2005). In order to determine whether cavitation will occur in any portion of a reaction turbine, the Thoma's cavitation factor ( $\sigma$ ) is compared with the critical cavitation factor ( $\sigma_c$ ). If the value of  $\sigma$  is greater than  $\sigma_c$ , cavitation will not occur in the turbine under analysis, where  $\sigma_c$  is a function of the specific speed of the turbine ( $N_s$ ). Because  $N_s$  is not used in the proposed methodology, the occurrence of cavitation cannot be determined using the utilized parameters. The occurrence of cavitation in reaction turbines needs be checked after using the proposed methodology. Following, an example of application for determining optimal flow discharge and optimal penstock diameter for an impulse turbine is presented. For an example of application of reaction turbines, the reader is referred to Leon and Zhu (2014).

### 3. EXAMPLE OF APPLICATION FOR AN IMPULSE TURBINE

The site, penstock and nozzle characteristics for this example are as follows:

1. Gross head ( $H_g$ ) = 200 m
2. Penstock length ( $L$ ) = 500 m
3. Ratio of penstock cross-sectional area to nozzle cross-sectional area at its outlet ( $A_2/A_N$ ) = 16
4. Nozzle velocity coefficient ( $C_V$ ) = 0.985
5. Sum of local losses in penstock due to entrance, bends, penstock fittings and gates ( $\Sigma k_{1-2}$ ) = 1.5
6. Roughness height of penstock material ( $\mathcal{E}$ ) = 0.045 mm (commercial steel)
7. Kinematic viscosity ( $\nu$ ) =  $10^{-6}$  m<sup>2</sup>/s
8. Turbine efficiency ( $\eta_t$ ) = 82%
9. Generator efficiency ( $\eta_g$ ) = 90%

#### 3.1. Case A1: $Q$ is Specified

In this case, it is assumed that the design flow  $Q$  is 0.6 m<sup>3</sup>/s, and it is desired to know the optimal hydropower that can be extracted using this flow. First, it is necessary to determine the optimal penstock diameter. From Eq. (24),

$$\frac{(C_L)_{\text{opt}}}{A_2^2} = 1693.8272 \text{ m}^{-4} \quad (26)$$

where  $C_L = 500f/D_2 + 1.5 + k_N(16^2)$ .

The nozzle coefficient is determined using Eq. (3), which gives  $k_N = 0.0307$ . The friction factor ( $f$ ) is determined using the explicit Swamee-Jain equation, which is given by

$$f = \frac{0.25}{\left[ \log_{10} \left( \frac{\epsilon}{3.7D_2} + \frac{5.74}{\text{Re}^{0.9}} \right) \right]^2} \quad (27)$$

where  $\epsilon$  is the roughness height and Re is the Reynolds number. The Reynolds number is defined as  $VD_2/\nu$ , where  $V$  is the flow velocity. Note that when  $Q$  is known,  $f$  and  $C_L$  are functions of  $D_2$  only. Solving for  $D_2$  in Eq. (26) gives  $D_2 = 0.3968$  m. In practice, a penstock with an internal diameter equal or slightly larger than 0.3968 m (397 mm) would be selected. Assuming that a schedule 80 steel pipe is required due to structural considerations, an 18-in outside diameter pipe would be selected. For this pipe, the wall thickness is 0.938 in, and, hence, the internal diameter is 16.124 in (409.5 mm). For this pipe diameter, the value of  $C_L$  is 25.35. This value can be used to determine the dimensionless head loss as follows (e.g., see Eq. 6).

$$h_{L+} = C_L \frac{Q^2}{2gH_g A_2^2} = 0.134 \text{ or } 13.4\% \quad (28)$$

which satisfies the inequality in Eq. (20) (< 15.6%). The electric power that can be extracted from this system can be determined using Eq. (6), which gives,

$$P = 0.82 \times 0.90 \times 1000 \times 9.8 \times 0.6 \times \left( 200 - 25.35 \times \frac{0.6^2}{2 \times 9.8 \times 0.1317^2} \right) = 751421 \text{ W} = 751.4 \text{ kW}$$

### 3.2. Case A2: $P$ is Specified

In this case, assume that  $P$  is 100 kW, and it is desired to determine the optimal flow discharge and optimal penstock diameter to produce this power. In this case, first, the optimal discharge is determined using Eq. (23) as follows:

$$Q_{\text{opt}} = \frac{45}{38} \left( \frac{100,000}{0.82 \times 0.90 \times 1000 \times 9.8 \times 200} \right) = 0.082 \text{ m}^3/\text{s} \text{ (82 L/s)}$$

The optimal pipe diameter (inside diameter) can be determined in a similar way to Case A1, which gives 0.176 m.

To facilitate the calculations, a MATLAB hydropower calculator was developed for which the graphical user interface (GUI) is shown in Figure 4. As can be observed in this figure, the consumption of flow is optimized in the linear region because the amount of power is proportional to the amount of flow used. Right before the large positive gradient in each curve, both the flow discharge and the penstock diameter are optimized. The hydropower calculator is available at <http://web.engr.oregonstate.edu/~leona/Codes/Hydropower/>.

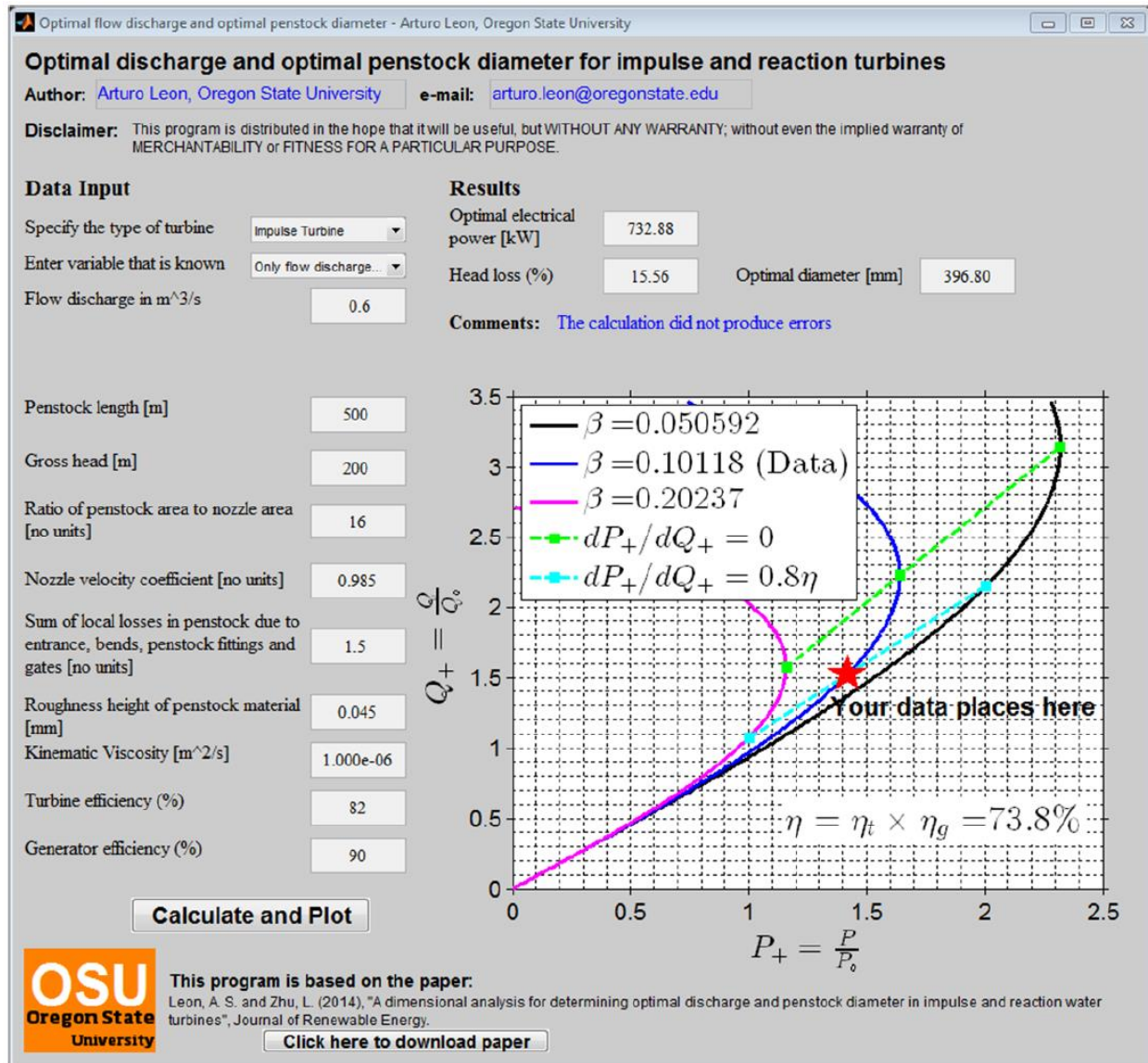


Figure 4. Graphical User Interface (GUI) of hydropower calculator

## 4. CONCLUSIONS

This paper presents a dimensional analysis for determining optimal flow discharge and optimal penstock diameter when designing impulse and reaction turbines for hydropower systems. The aim of this analysis is to provide general insights for minimizing water consumption when producing hydropower. The key findings are as follows:

1. The analysis is based on the geometric and hydraulic characteristics of the penstock, the total hydraulic head, and the desired power production. This analysis resulted in various dimensionless relationships between power production, flow discharge, and head losses.
2. The derived relationships were used to draw general insights on determining optimal flow discharge and optimal penstock diameter. For instance, it was found that for minimizing water consumption, the ratio of head loss to gross head ( $h_l/H_g$ ) should not exceed about 15%.
3. To facilitate the calculations, a MATLAB hydropower calculator was developed which is available at <http://web.engr.oregonstate.edu/~leona/Codes/Hydropower/>.

4. Overall, the present analysis is general and can be used for determining optimal design flow and penstock diameter when designing impulse and reaction turbines.

## 5. ACKNOWLEDGMENTS

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