Observations of flow and sediment entainment on a large gravel-bed river

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Abstract. Constant-discharge reservoir releases on the Trinity River, California, provide an unusual opportunity to unambiguously relate flow and gravel entainment on a large gravel-bed river. Bed shear stress $\tau_0$ was estimated using local observations of depth-averaged velocity. Gravel entainment was measured using large tracer gravel installations. Lateral variability of $\tau_0$ is large, even for straight channels with simple, trough-like geometry. No simple relation exists between local and cross-section mean values of $\tau_0$. Fine grains (less than 8 mm; 20–30% of the bed material) are transported at lower discharges than coarse grains. Scour to the base of the bed surface layer occurs at a dimensionless shear stress $\tau^*_b \approx 0.035$, for $\tau^*_e$ formed using local $\tau_0$ and the median grain size of the gravel portion of the bed. The dimensionless reference transport rate $W^* = 0.002$, often used as a surrogate for the threshold of grain motion, occurs at nearly the same $\tau^*_b$. At smaller $\tau^*_e$ entrainment and transport rates decrease rapidly, becoming vanishingly small at $\tau^*_b \approx 0.031$. Even at very small gravel transport rates, all sizes are transported, although the coarsest sizes are in a state of partial transport in which only a portion of the exposed grains are entrained. Both entainment and cumulative transport observations suggest that maximum scour depth for plane-bed transport is slightly less than twice the surface layer thickness.

Introduction

The problem of predicting the critical river discharge $Q_c$ that initiates sediment movement is of wide importance. Estimates of $Q_c$ are needed to determine sediment transport rates, the frequency and duration of channel-forming flows, the dimensions of stable channels, and the occurrence of bed scour, armoring, downstream fining, and fine sediment infiltration or removal. Because $Q_c$ depends on local channel properties, the entrainment problem is commonly posed in terms of a critical bed shear stress $\tau_*$ for incipient motion, which may be defined as a function of sediment properties alone. For application, an estimate of $\tau_*$ must be coupled with an estimate of the bed shear stress $\tau_0$ as a function of discharge, channel geometry, and hydraulic roughness.

A number of factors limit the accuracy of estimates of $Q_c$ in large gravel-bed rivers, preventing routine application from readily observable properties of the channel geometry and sediment. One problem is uncertainty in determining values of $\tau_*$ from the bed size distribution. Ideally, $\tau_*$ should be estimated as a function of easily measurable properties of the bed sediment, particularly its size distribution, but problems of method, measurement, and accuracy give rise to uncertainty in the prediction of $\tau_*$ for mixed-size sediments [Wilcock, 1988, 1992a]. Another obstacle is the typically large spatial variability of both flow and grain size in large gravel-bed rivers. In simple equilibrium transport fields, such as flumes or small rivers with uniform geometry, spatial averages of $\tau_0$ and grain size may be used to estimate $\tau_*$ and $Q_c$. The spatial variation of flow and grain size in larger rivers ensures that this approach will be inaccurate locally; the net effect of local error may be quite large, particularly when transport occurs in part of the channel even though the section-average $\tau_0$ is less than $\tau_*$. Error in either $\tau_*$ or $\tau_0$ is compounded by the steep, nonlinear relation between $\tau_0$ and the sediment transport rate, which can produce large errors in calculated transport rate from small errors in either $\tau_*$ or $\tau_0$.

A third obstacle to predicting $Q_c$ concerns measurement logistics and accuracy: in large rivers it is generally not possible to measure flow and sediment entrainment with sufficient accuracy at a temporal and spatial resolution appropriate to the elementary physics of the problem. Clearly, both spatial variability and measurement problems contribute directly to uncertainty in developing predictive relations for $\tau_*$. Finally, the stochastic nature of gravel transport ensures that no single value of $Q_c$ exists. Grain motion near $\tau_*$ occurs in sporadic, brief events, separated by relatively long periods of immobility. The range of grain sizes typically present on a gravel bed, combined with the spatial variability in grain size and bed topography, ensures that uniform entrainment will not occur at a single $Q$ but will consist of individual grain displacements whose frequency varies spatially and increases with $Q$. This local variability points to aspects of grain entrainment with importance beyond defining a threshold for measurable transport. In particular, the proportion of the bed surface mobilized during a flow and the entrained proportion of individual size fractions exposed on the bed surface control sediment
exchange between the transport, bed surface, and subsurface, and therefore changes in the bed surface size distribution and size-selective erosion or deposition.

We measured local flow and gravel entrainment on two reaches of the Trinity River, a large gravel-bed river in northern California. Unusually favorable field conditions were provided by three trial reservoir releases which had nearly constant discharge, were scheduled in advance, and were immediately preceded and followed by very low discharges permitting direct observation and sampling of the gravel bed. The entrainment produced by each release could be determined without ambiguity by measuring the proportion of grains removed from large tracer gravel installations. This new approach provides estimates of the size and scour depth of entrained sediment and the proportion of the bed surface mobilized. The local shear stress $\tau_{oi}$ over the tracer gravels was determined from observations of depth-averaged velocity, providing a basis for generalizing our results through comparison with observations made in flumes or small rivers.

Incipient gravel motion is often represented as the value of $\tau_{ci}$ that produces a small reference transport rate [e.g., Parker et al., 1982; Wilcock, 1988; Wilcock and Southard, 1988; Ashworth and Ferguson, 1989; Kuhnle, 1992]. To provide a comparison with our entrainment observations, we measured gravel transport rates using gravel traps and bed load samplers. Although the reference-transport estimate of $\tau_{ci}$ is clearly defined and measurable, it provides no indication of the proportion of the bed contributing to the transport. The tracer gravel observations provide such a measure, and a comparison of tracer entrainment and transport rates demonstrates the degree of bed mobilization associated with different fractional transport rates.

Critical Shear Stress and Bed Surface Mobilization

There has been considerable discussion and some disagreement about appropriate values for $\tau_{ci}$ in coarse mixed-size sediments. It is well known that $\tau_{ci}$ increases directly with grain size for unisize sediments in the gravel size range. The same relation has also been shown to hold for the median grain size $D_{50}$ of many mixed-size sediments [Wilcock, 1992a]. An exception occurs for gravel beds with a large proportion of sand, giving a strongly bimodal size distribution and apparently reducing $\tau_{50}$ relative to unimodal size mixtures [Wilcock, 1993].

The critical shear stress $\tau_{ci}$ for individual grain sizes $D_i$ within mixed-size sediment depends strongly on the relative grain size within the mixture (e.g., $D_i/D_{50}$). Size-dependent hiding, exposure, and resistance cause $\tau_{ci}$ to be reduced for larger grains and increased for smaller grains relative to comparable values of $\tau_{ci}$ in unisize sediment. Much of the recent work (and disagreement) has focused on whether relative size effects entirely cancel absolute size effects, producing values of $\tau_{ci}$ that are independent of grain size [e.g., Parker et al., 1982; Andrews, 1983; Carling, 1983; Komar, 1987; Wilcock and Southard, 1988; Ashworth and Ferguson, 1989; Church et al., 1991; Wilcock, 1993] (see also recent review by Gomez [1995]).

Discrepancy among $\tau_{ci}$ results may be attributed to the spatial and temporal variability of gravel transport [Gomez, 1983; Iseya and Ikeda, 1987; Kuhnle and Southard, 1988; Pitlick, 1988; Wathen et al., 1995]. To the difficulty of measuring values of $\tau_{ci}$ appropriate to the observed entrainment, and to differences in the methods used to estimate $\tau_{ci}$ [Wilcock, 1998].

Estimates of $\tau_{ci}$ are typically based on observations of gravel transport and therefore provide little guidance regarding the proportion of the bed surface that is entrained, the mobilized proportion of each size fraction, and the depth of grain exchange between the bed and transport. The depth of grain exchange is needed to define an active layer thickness for models of selective transport and to estimate transport rates from observations of the cumulative displacement of tracers. The mobilized proportion of each size on the bed surface influences both the rate and size distribution of sediment exchange between bed and transport, which controls the evolution of the bed surface size distribution and therefore the rates of selective transport, downstream fining, and armoring. The proportion of the bed surface entrained also determines the quantity of finer sediments that may be introduced into, or removed from, the bed surface when coarse surface clasts are entrained. This is the application underlying the field work described here: removal of infiltrated fine sediment was one goal of the trial reservoir releases.

For sediments with a wide range of grain sizes (exceeding a factor of $\sim 16$), it has been observed that a range of transport exists over which the transport rates of the coarser fractions can be orders of magnitude smaller than those of the finer fractions, when each is scaled by its presence in the bed [Wilcock, 1992b; Wilcock and McArdell, 1993; Wathen et al., 1995]. Wilcock and McArdell [1993] proposed that the relative decrease in transport rate with increasing grain size can be attributed, in part, to the fact that a portion of the coarser grains on the bed surface remain immobile over the duration of a flow. This condition, termed partial transport, has importance regarding not only fractional transport rates but any process (such as selective deposition or flushing of subsurface fines) that depends on the proportion of the bed surface entrained. Confirmation of a state of partial transport requires direct observation of the mobilized proportion of the grains exposed on the bed surface. Such observations have been made in the laboratory using the unusual artifice of a sediment bed in which each size fraction was painted a different color [Wilcock and McArdell, 1993]. This work confirmed the existence of a state of partial transport and demonstrated its dependence on grain size and flow strength [McArdell and Wilcock, 1994].

The tracer gravel method described in this paper provides an opportunity to estimate the mobilized proportion of the bed surface and the depth of sediment exchange. Because transport rates were also measured, we can compare tracer entrainment and transport rates to determine the degree of bed mobilization associated with different transport rates.

Study Site

The Trinity River drains 7640 km$^2$ of steep, dissected terrain in the Klamath Mountains of northwestern California. The basin is mostly forested, although much of it has been clearcut since 1950. Runoff from the uppermost 1860 km$^2$ of the basin was impounded by Trinity Dam (and its reregulating reservoir, Lewiston Dam) beginning in 1961, as part of the U.S. Bureau of Reclamation Central Valley project. Our flow and transport observations were made along two reaches, Poker Bar and Steelbridge, located 15 km and 20 km downstream of Lewiston Dam, respectively (Figure 1). Overall, the river has a broadly sinuous pool-and-riffle channel, although both study reaches have nearly straight channels with simple, gradually varying topography. Most observations were focused on one section within each reach (PB2 and SB3C, Figure 2), and the sections immediately upstream and downstream.
Area of detail

CALIFORNIA

LEWISTON DAM

STEEL BRIDGE STUDY SITE

POKER BAR STUDY SITE

Figure 1. Location map of the Trinity River downstream of Trinity and Lewiston Dams.

At Poker Bar the channel is ~35 m wide and rectangular in section, with a deeper trough along the right bank (Figure 2a). The river banks are nearly vertical and composed of fine-grained material (<0.5 mm) deposited along the margins of the much wider active channel that existed before the Trinity and Lewiston Dams were closed in 1963. Extremely low discharges following dam closure permitted vegetation to become established within the former active channel and the banks have apparently been built during occasional tributary floods carrying high concentrations of fine sediment [Wilcock et al., this issue]. Bankfull discharge for the present channel is approximately 75 m³/s (Figure 2a).

At Steelbridge the river is split into two channels by an island that developed from vegetation encroachment and fine-sediment accretion on a midchannel bar deposited after dam closure. The study section is in the smaller right channel, which carries a nearly constant 32% of the discharge for all flows observed by us. Here, the channel is ~20 m wide with steep, fine-grained banks similar to those at Poker Bar (Figure 2b).

The bed at the study sections is composed of sand, gravel, and cobble. Material finer than 8 mm composes between 20 and 30% of the bed, forming a matrix that completely fills the interstices of the larger framework grains. The fine material is derived from decomposed granitic terrain drained by tributaries downstream of the dam. It has a distinctly lighter color and nearly all of it is between 0.5 mm and 8 mm in size. By filling coarse pore spaces the fine material limits habitat for aquatic invertebrates and juvenile fish and is thought to limit salmonid spawning success by blocking fry emergence from the bed. Sediment finer than 0.5 mm is transported through as wash load, except for the small fraction deposited on the channel banks at high stage. A number of factors, including sampling logistics, the distinction between matrix and framework bed material, and a difference in ecological role, suggest that material finer than 8 mm may be treated as a separate population from the gravel and cobble. By local convention, which we...

Figure 2. Map of surveyed channel boundaries and cross sections for the (a) Poker Bar and (b) Steelbridge study reaches. Cross-section topography shown for sections PB2 and SB3C and adjacent sections. All lengths in meters. The plotted water level corresponds to \( Q = 76 \) m³/s. A small channel carrying less than 1% of flow exits the left side of the Poker Bar reach between PB1B and PB2. At Steelbridge, flow is divided into two channels separated by an island; the study reach is the right channel, which carries approximately 32% of the discharge. SB4 marks the downstream end of the island. Cross sections are not to scale. The channels are generally wide, shallow, and trough-like with gradually varying topography.
follow here as a matter of convenience, sand and gravel are defined as material finer or coarser than 8 mm, respectively.

Bed-material size distributions, based on large bulk samples taken before and after the 1992 release, are given for both study sections in Figure 3. Figures 3a and 3b present the cumulative size distribution for the entire bed; Figures 3c and 3d present the size distributions truncated at 8 mm. The samples were taken by working a 50-cm oil drum into the bed and extracting the material within. Before-and-after samples were taken at three locations along PB2 (26, 28.5, and 30 m, Figure 2a) with sample depth between 27 and 40 cm, which is 3–5 times the thickness of $D_{90}$ (size for which 90% of the distribution is finer). Along SB3C, four pairs of samples were taken (11, 12.5, 14.5, and 16.5 m, Figure 2b) with sample depths between 19 and 29 cm, which is 2–3 times the thickness of $D_{90}$ at that section. Because turbidity within the drum prevented observation of the bed during sampling, no attempt was made to sample separately the bed surface and subsurface. Pebble counts were also made before and after the releases along each section shown in Figure 2, typically with 100 counts per section. At PB2, 200 counts were made between 24 and 32 m, and 100 counts were made between 15 and 24 m and between 32 and 37 m. At SB3C, 200 counts were made between 3 and 18 m, and 100 counts were made between 11 and 16.5 m.

From bulk samples the median size $D_{50}$ of the bed is 22 mm at PB2 and 37 mm at SB3C using the entire size distribution and 36 mm and 58 mm, respectively, for the size distribution truncated at 8 mm. Values of $D_{90}$ from pebble counts over the same part of the section were 29 mm at PB2 and 44 mm at SB3C. These values are only slightly larger than the bulk sample values, suggesting that surface coarsening is relatively weak at the study sections. Values of $D_{90}$ for pebble counts and bulk samples are nearly the same at both sections.

Little decrease in sand proportion occurred over the duration of the releases (Figure 3), a result of the quantity of sand present in the bed upstream of the study sections and the limited duration of the releases [Wilcock et al., this issue]. The sand size distributions show some spatial variability, whereas the gravel distributions are relatively constant at each section.

Flow Observations

Releases from Lewiston Dam were made from May 28 to June 2, 1991, June 10 to 19, 1992, and April 13 to May 4, 1993 (Figure 4). At our study sections we observed a constant dis-charge for 3 days in 1991 (76.0 m$^3$/s), 4 days in 1992 (164 m$^3$/s), and 2 days in 1993 (80 m$^3$/s, Table 1).

During the releases, velocity and depth were measured from rafts moved along the study sections using a network of climbing ropes (see Wilcock et al. [1995] for detail). Most observations were made at sections PB2 and SB3C, with additional transects made at the sections immediately upstream and downstream. Velocity was measured with Price AA current meters. Except in 1992 at Poker Bar, the meters were mounted on wading rods up to 3 m long. The rods were held off the bow with the rod base on the channel bottom. All observations in a vertical profile were made without moving the rod to ensure accurate relative placement of the meter. Greater flow depths in 1992 at Poker Bar required that the current meters be mounted on a cable with a 100-lb weight and moved in the vertical with a crane-and-reel assembly. Cable mounting permits less spatial control, particularly in the horizontal, than rod mounting.

Complete velocity transects were made on at least two different days across sections PB2 and SB3C in 1991 and 1992 when tracer gravels were in place and across section PB2 in

Precise estimates than using either a single near-bed velocity only grain-scale roughness and flow depth much greater than logarithmic variation with elevation throughout the water file [Wilcock et al., 1994; Wilcock, 1996]. The relation used here observation or the slope of the near-bed vertical velocity pro-

the flow conditions we observed, this method provides more estimates of $\tau_0$ were made using the flow depth $h$, the depth-integrated velocity $U$, and an estimate of the bed roughness. Estimates of $\tau_0$ were made using the flow depth $h$, the depth-integrated velocity $U$, and an estimate of the bed roughness $Z_0$ in a flow resistance relation. An evaluation of alternative methods for estimating $\tau_0$ demonstrated that for the flow conditions we observed, this method provides more precise estimates than using either a single near-bed velocity observation or the slope of the near-bed vertical velocity profile [Wilcock et al., 1994; Wilcock, 1996]. The relation used here is the depth-integrated form of the log law applied throughout the flow depth

$$\frac{U}{u_*} = \frac{1}{\kappa} \log \left( \frac{h}{\varepsilon Z_0} \right)$$

where $u_*$ is bed shear velocity ($=\tau_0/p)^{1/2}$, $p$ is the water density, $\kappa$ is von Karman’s constant, taken to be 0.4, and $\varepsilon$ is the base of the natural logarithms. The resulting $\tau_0$ estimates were found to have the same mean value but more precision than estimates of $\tau_0$ made from the slope of the semilogarithmic velocity profile for 22 replicate profiles with at least six observations in the lower half of the flow. Although the latter estimate has lower precision, it does not require an independent estimate of bed roughness. Therefore it was used to provide an estimate of $Z_0$ as a function of visual estimates of $D_{90}$ at each station. The relation $Z_0 = 0.095D_{90}$ was found to minimize the squared difference in $u_*$ between the two estimates for 75 velocity profiles. This result is almost identical to $Z_0 = 0.1D_{95}$, found by Whiting and Dietrich [1990] in a more detailed study of boundary shear stress and roughness.

Values of $\tau_0$ and unit discharge $q = Uh$ (normalized by the section mean $Q/B$, where $B$ is channel width) are given in Figure 5 for the 1991 and 1992 releases. Also shown on Figure 5 are the locations of the tracer gravels and gravel traps. Observations for individual days are shown with light lines, and the mean values carried forward in the analysis are shown by heavy lines with symbols. Measurement accuracy is suggested by the relatively small variation between plots of $q$ for different days with the same discharge (Figures 5a and 5b) and by the fact that the largest difference in total discharge calculated for periods with constant stage was always <5% (based on three groups of five $Q$ observations and one group of two). Somewhat greater scatter is evident in the cable-mounted observations made at Poker Bar in 1992, although even in that case, it is evident that the error is smaller than lateral variation in $q$ or $\tau_0$.

The lateral distribution of flow is somewhat more uniform at PB2 than SB3C, where flow becomes increasingly concentrated in a central portion of the channel as discharge increases (Figure 5b). A large lateral variability is evident in $\tau_0$. The variation of $\tau_0$ is nearly threefold within the central 15 m of the nearly rectangular SB3C channel (neglecting the regions within approximately 4 m of the banks). This lateral variability makes it clear that estimates of local gravel entrainment and transport, as well as total transport through the section, require local flow observations, even in straight channels with simple geometry.

Figure 6 provides an approximate comparison between section-averaged and local values of mean velocity, flow depth,
unit discharge, and \( \tau_0 \). Here, local values are averages over the central portion of the section containing tracer gravels. Section-average calculations use the flow area above the main channel because flow in the heavily vegetated overbank regions had negligible velocity. Section-averaged flow depth \( H \) is calculated as \( A/B \), where \( A \) is the flow area of the channel. Section-averaged \( \tau_0 \) is calculated as \( p g H S \), where \( S \) is the slope of the energy grade line calculated using section-averaged \( U \) and the observed water surface elevation at the main section and the sections immediately upstream and downstream. Section-averaged \( \tau_0 \) is also calculated by integrating \( \tau_0 \) across the section. These values agree fairly closely with those calculated as \( p g H S \) (Figures 6g and 6h), suggesting that only a weak lateral component exists in the flow. The anomalous large value of \( p g H S \) at \( Q = 76 \) m\(^3\)/s at SB3C is probably related to uncertainty in the values of mean channel velocity at each section because the flow is beginning to exceed the bank level and some leakage of flow from the left to the right channel may occur over the island.

Local velocity increases with \( Q \) at both sites, although the increase relative to the section average is much larger at SB3C, where the mean velocity shows only a weak increase with \( Q \) (Figures 6a and 6b), probably as the result of a backwater from the flow merging with the larger left channel at SB4. A similar pattern is evident in the local unit discharge: both \( q_l \) and \( \partial q_l / \partial Q \) are larger at SB3C (Figures 5, 6e and 6f). At both sites, local values of \( \tau_0 \) are larger than section-averaged values and increase more rapidly with discharge. At PB2 the mean local value increases to a value 30% greater than the section mean, while at SB3C the mean local value is 65% larger at the highest flow observed. Clearly, if estimates of grain entrainment over the central portion of the channels were based on the section-averaged \( \tau_0 \), the entrainment would be grossly underestimated. Differences in estimated \( Q_c \) of the order of a factor of 2 are possible at PB2, with even larger differences at SB3C. Given the very strong nonlinear relation between \( \tau_0 \) and gravel transport, such errors in estimating \( Q_c \) could produce errors of several orders of magnitude in estimated transport rates.

**Gravel Movement**

**Tracer Gravels**

Our primary observations of grain entrainment were made using large tracer gravel installations on the main study sections. After surveying the bed elevation, a metal cylinder was inserted into the bed, all sediment down to the bottom of the sampler was excavated (similar to the method of McNeil and Ahnell [1960]), and its size distribution was determined. In 1991 we used a sampler with a diameter of 30 cm and sample depths up to 30 cm, giving sample sizes between 13 and 30 kg. We enlarged the sample size in 1992 by using a sampler with a diameter of 59 cm and sampling as deep as 40 cm (Table 2).
Figure 6. Variation with $Q$ of (a) and (b) velocity, (c) and (d) flow depth, (e) and (f) unit discharge, and (g) and (h) bed shear stress at sections PB2 and SB3C. Both section average and local values shown. Local values are averages over the tracer gravels (lateral station 23.5–35.5 m for PB2; 10–17 m at SB3C; see Figure 5). Section-averaged values exclude overbank flow. Section averaged $\tau_0$ is calculated using $\rho g H S$; also shown is mean $\tau_0$ obtained by integrating local $\tau_0$ across each section.

These samples contained between 112 and 281 kg and provide our best measure of the bed size distribution.

The sediment was returned to the cylinder as a homogeneous mix and filled to the presample elevation. All sizes larger than 16 mm were replaced with distinctly marked tracer grains, matching the number and shape of grains in each $1/2\phi$ size class. During the 1991 release we replaced the entire bulk sample with pure white quartz clasts. During the 1992 release we replaced a representative portion of the sampled sediment with painted grains. The tracer clasts were placed in the center of the bulk sample location, using a 25-cm-diameter plastic bucket with the bottom removed, and the annulus was filled with the balance of the bulk sample. This approach was used to facilitate recovery of all of the remaining tracer clasts after the release.

When replacing the tracer sediment, we attempted to replicate the bulk density of the natural bed, as could be approximately determined by walking over tracers and undisturbed material. Nonetheless, the sampling disturbance raises the possibility that entrainment of tracer and undisturbed grains might differ. Although we cannot independently confirm that the structure and entrainment of the tracers mimicked that of the natural bed, several lines of evidence suggest that the tracer entrainment is representative. First, active spawning has been observed in the study reaches in the 5 years between the start of this study and the previous gravel-transporting flow in 1986, suggesting that the bed should be relatively well-mixed with little internal structure. This is supported by our observation of little or no armoring at the study sections. Second, tracer entrainment in both 1991 (partial surface mobilization) and 1992 (complete surface mobilization) was consistent with visual estimates of entrainment made possible by the fact that a
dark stain develops on the exposed side of surface grains, so that mobilized grains typically appear clean and lighter in color. From these observations we assume, but cannot prove, that the tracer entrainment was comparable to local entrainment of undisturbed material. Application of this tracer method to channels with an armored, imbricated bed would be more difficult, unless a means were found to replicate the surface structure.

After the trial releases, we resampled the tracers to determine the number and size of immobile grains, from which the entrained mass was determined. The depth of bed scour, or exchange depth $dx$, was calculated as the sample depth multiplied by the mass proportion of grains entrained (Table 2). Because grain size and sample depth vary from sample to sample, comparison among samples requires a consistent basis for scaling $dx$, for which the thickness of the bed surface layer is appropriate. Using $D_{90}$ as the scaling length, the scaled exchange depth is calculated as

$$
\frac{dx}{D_{90}} = \frac{M_e}{M_t} \frac{d_s}{D_{90}}
$$

where $M_e$ is the mass of grains entrained, $M_t$ is the total mass of tracer gravels, and $d_s$ is the sample depth.

A fractional exchange depth $d_{x_i}$ is calculated as the entrained proportion of grains of each size times the sample depth. Scaled by $D_{90}$, this is

$$
\frac{d_{x_i}}{D_{90}} = \frac{(M_{ei})}{(M_{ti})} \frac{d_s}{D_{90}}
$$

where $(M_{ei})$ and $(M_{ti})$ are the entrained and total mass of grains of size $i$. Because $(M_{ei})/(D_{90}d_s)$ approximates the mass of grains of size $i$ in the surface layer, values of $d_{x_i}/D_{90}$ less than one may also be interpreted as the mobilized proportion of surface grains of each size.

To account for differences in grain size from sample to sample, $d_{x_i}/D_{90}$ is plotted in Figure 7a as a function of the Shields parameter $r_g = q_{th}(\rho_s - \rho)D_g g$, where $\rho_s$ and $\rho$ are the sediment and water density, $g$ is acceleration of gravity, and $D_g$ is the median size of the gravel portion of the bed. $D_g$ is used in preference to $D_{90}$ of the entire bed because the sand content (and therefore $D_{90}$) may be expected to be more transient than that of the gravel. The effect of using $D_{50}$ of the entire bed in $r_g$ would be a relatively uniform shift of the points in Figure 7 to the right by roughly 60%.

Direct comparison of the entrainment produced by the two releases implies that both flows were of sufficient duration to mobilize nearly all entrainable grains, so that a longer release would not produce measurably larger values of $d_{x}/D_{90}$. Although we cannot demonstrate that such a limit was reached, the mobilized proportion of individual fractions has been observed to asymptotically approach a limiting value in a laboratory study of grain entrainment [McArdell and Wilcock, 1994; P. R. Wilcock and B. W. McArdell, Partial transport of a sand/gravel sediment, submitted to Water Resources Research, 1996].

Scaled exchange depth $d_{x}/D_{90}$ shows a consistent increase with $r_g$ (Figure 7a). If $d_{x}/D_{90}$ is assumed to increase rapidly for $r_g$ just above a threshold $r_{th}$ at which $d_{x}/D_{90} = 0$, and asymptotically approach a limiting value of $d_{x}/D_{90}$ at large $r_g$, an appropriate function is

$$
\frac{dx}{D_{90}} = a \left(1 - \frac{r_{th}}{r_g}\right)^b
$$

For $r_{th} = 0.031$, values of $a = 1.7$ and $b = 0.26$ minimize the sum of squared differences between (4) and the six 1992 observations in Figure 7a. This suggests a limiting scour depth on

![Figure 7](image-url)

**Table 2. Tracer Gravel Installations**

<table>
<thead>
<tr>
<th>Sample</th>
<th>Depth $d_{x_i}$</th>
<th>Mass, kg</th>
<th>Pebble Count</th>
<th>Exchange Depth $d_{x_i}$</th>
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<td>9</td>
<td>14</td>
<td>100</td>
<td>0.18</td>
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<td>14</td>
<td>28</td>
<td>100</td>
<td>0.28</td>
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<tr>
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<td>3</td>
<td>22</td>
<td>120</td>
<td>0.18</td>
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<td>15</td>
<td>120</td>
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</tr>
<tr>
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<td>120</td>
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</tr>
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<td>112</td>
<td>100</td>
<td>10.4</td>
</tr>
<tr>
<td>1992 SB3C-14.5</td>
<td>21</td>
<td>132</td>
<td>100</td>
<td>8.8</td>
</tr>
<tr>
<td>1992 SB3C-16.5</td>
<td>19</td>
<td>120</td>
<td>100</td>
<td>10</td>
</tr>
</tbody>
</table>

*No scour observed, tracers buried to a depth of $\sim D_{90}$. 

* $D_{90}$ is used in preference to $D_{90}$ of the entire bed. 

* $r_{th} = 0.031$, values of $a = 1.7$ and $b = 0.26$ minimize the sum of squared differences between (4) and the six 1992 observations in Figure 7a. This suggests a limiting scour depth on...
the order of $1.7D_{90}$ for plane-bed transport. A mean exchange depth of $d_{1/3}/D_{90} \approx 1.0$, or scour to the base of the bed surface layer, occurs at $\tau^*_{0} \approx 0.035$, which is similar to many estimates of the dimensionless shear stress for incipient motion of the median size of mixed size sediment [Wilcock, 1992a, 1993]. The threshold between negligible grain entrainment and entrainment of most of the bed surface occurs over a narrow range of $\tau^*_{0}$ of the order of 10–15%.

For all cases with $d_{1/3}/D_{90} > 0.1$ (SB3D in 1991, SB3C and PB2 in 1992), all sizes are entrained and fractional exchange depths fall within a factor of 2, with the very largest fractions tending to have the smallest entrainment (Figure 7b). In 1992 at SB3C, sizes larger than 90 mm ($\sim D_{72}$) are in a state of partial transport ($d_{1/3}/D_{90} < 1$), whereas all finer sizes and $d_{1/3}/D_{90}$ take values $\geq 1$. In 1992 at PB2, sizes larger than 32 mm ($\sim D_{60}$) fall below $d_{1/3}/D_{90}$, whereas the finest two fractions have $d_{1/3}/D_{90} > 1.8$. In 1991, the finest sizes observed at PB2 (26 mm) and at SB3C (19 mm) show weak partial transport ($0.2 < d_{1/3}/D_{90} < 0.4$), whereas all the coarser sizes are essentially immobile and $d_{1/3}/D_{90} < 0.1$. In 1991, $\tau_{01}$ at SB3C was slightly larger than at SB3C, giving $0.25 < d_{1/3}/D_{90} < 0.5$ (Figure 7a), and grains of all sizes finer than $D_{90}$ were moved.

The 1991 discharge of $Q = 76$ m$^3$/s produced very little entrainment (all $d_{1/3}/D_{90} < 0.5$, most $d_{1/3}/D_{90} < 0.1$), whereas in 1992, $Q = 164$ m$^3$/s entrained the bed to a depth of roughly one surface layer or greater ($0.8 < d_{1/3}/D_{90} < 1.6$) and mobilized all sizes up to 128 mm ($D_{98}$ at SB3C). These observations support a direct specification of the discharge needed to mobilize the bed surface, although this estimate is specific to our study reach and comes at the expense of extensive field work, the logistics of which depended on a prescheduled hydrograph with a simple rectangular shape.

**Transport Rates**

Gravel transport rates were measured at $Q = 80$ m$^3$/s in 1993 using five wooden boxes buried flush with the streambed across PB2. The boxes were 80 cm long, 12 cm wide, and 10.5 cm deep and extended across 15 m of section (Figure 5 and Table 3). Sand was visibly in motion within the traps, so an unknown proportion was swept downstream and the boxes are assumed to act as efficient traps only for grains coarser than 8 mm. The traps were first visited 14 days after the release began, at which time all traps were less than half full. The volume and grain size of trapped material were visually estimated, and the boxes were swept clean. Sediment continued to accumulate in the boxes until the end of the release, providing a second sample with a duration of 68 hours. The accumulated sediment was then removed by hand, weighed, and sieved. The accumulation rates from both sample periods were comparable, but only the second period is used here because the mass and size distribution of the trapped sediment were measured directly.

The gravel transport rates in 1993 were uniformly very small, with a mean transport rate of 0.0045 g/ms (Table 3). The catch included one clast in the 64–90.5 mm size class and from five to more than 300 grains in the smaller size classes. Although the largest of the five samples is 30 times the smallest, there is no obvious lateral trend in the transport rates and much of the trap-to-trap variability may be attributed to the very small transport rates and the fact that the sample size is sensitive to the presence or absence of individual large clasts whose mass is on the order of the total for each trap. As a result, we use the aggregate transport for all five traps. As discussed below, the measurement of such small transport rates proves very useful in estimating the relation between gravel transport rate and $\tau_{01}$.

Transport rates were also measured with Helley-Smith samplers in 1992 and 1993. A sampler with a 15.2-cm orifice was used in 1992. Because of the weight of the sampler and the large water depth and velocity the sampler was deployed from a crane-operated cable mounted on a wooden raft. It was not always possible to directly observe the sampling. Samples were collected during the peak discharge of 164 m$^3$/s and during a 12-hour steady flow of 103 m$^3$/s on the rising limb of the hydrograph. A sampler with a 3-inch (7.62 cm) orifice was used.

---

**Table 3. Poker Bar Bed Load Sampling**

<table>
<thead>
<tr>
<th>Date</th>
<th>Discharge, m$^3$/s</th>
<th>Station</th>
<th>Total Number of Samples</th>
<th>Total Sample Duration</th>
<th>Gravel Transport Rate, g/ms</th>
<th>Sand Transport Rate, g/ms</th>
<th>Number of Censored Samples</th>
<th>Gravel Transport Rate, g/ms</th>
<th>Sand Transport Rate, g/ms</th>
<th>Censored Gravel Transport Rate, g/ms</th>
<th>Censored Sand Transport Rate, g/ms</th>
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<tr>
<td>April 27–30, 1993</td>
<td>80</td>
<td>19</td>
<td>2</td>
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<td>...</td>
<td>...</td>
<td>...</td>
<td>...</td>
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</tr>
<tr>
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<td>22</td>
<td>2</td>
<td>4 min</td>
<td>...</td>
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<td>...</td>
<td>...</td>
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<td>...</td>
</tr>
<tr>
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<td>2</td>
<td>4 min</td>
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<td>...</td>
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</tr>
<tr>
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<td>30</td>
<td>106.1</td>
<td>51.9</td>
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in 1993. The lighter sampler and smaller flow depths permitted the sampler to be deployed on a rigid rod, and the collection of all samples was directly observed. Results are presented here for samples made across the same 15-m portion of PB2 that incorporates the gravel traps. To provide some measure of the possible oversampling (through bed scooping) or undersampling (when the base of the sampler is not flush with the bed surface) to which the Helley-Smith sampler is susceptible, the Helley-Smith transport rates were recalculated using approximately one-half of the samples, excluding the largest 25% and smallest 25% of the observed transport rates for each station. These values are presented as “censored transport rates” on Table 3.

For the Helley-Smith samples the ratio of maximum to minimum station gravel transport rate is approximately 25 for $Q = 103$ m$^3$/s and 3.6 for 164 m$^3$/s; for sand transport rates, the ratios are 10 for $Q = 80$ m$^3$/s, 2.6 for $Q = 103$ m$^3$/s and 5.4 for 164 m$^3$/s. These are typical values for the spatial and temporal variability of gravel transport [Kuhnele and Southard, 1988; Pitlick, 1988; Gomez et al., 1989; Wathen et al., 1995]. In lieu of a statistical analysis of the relatively small number of samples, we assume no better than order-of-magnitude accuracy in the mean transport rates reported on Table 3, based on the variability in transport from station to station and between total and censored samples. The least confidence is associated with the observations at $Q = 103$ m$^3$/s and the sand transport at $Q = 80$ m$^3$/s, which are based on the smallest number of samples. Despite the large variation from sample to sample, we place somewhat greater confidence in the very small aggregate transport rates for the gravel traps because of their long duration and greater sampling width.

Total and fractional transport rates for each flow are summarized on Figure 8. Transport rate is plotted as a function of the mean $T_0$ over the sampled width. Different relations are evident for the total gravel and sand transport rates. At the smallest flow, sand transport exceeds that of gravel by 4 orders of magnitude. Gravel transport rate increases far more rapidly than that of sand, so that the two transport rates are comparable at the highest flow observed. For individual fractions within the sand and gravel, fractional transport rates differ by 1 order of magnitude or less and overlap only at $Q = 164$ m$^3$/s, for which the total sand and gravel transport rates are nearly identical.

The gravel transport rates are well matched by the Parker [1979] transport relation, which is a power approximation of the Einstein relation at low shear stresses. In the units of Figure 8, this relation is

$$q_b = \frac{\tau_0^{1.5}}{17.8} \left( 1 - 0.85 \frac{\tau}{\tau_0} \right)^{1.5}$$

(5)

where $\tau_0$ is the reference shear stress that produces a small reference transport rate and serves as a surrogate for the critical shear stress $\tau_c$. The reference transport rate is given as $W^* = 0.002$, where $W^* = [q_b(s - 1)g]/[\rho_s u^3]$, $s$ is $\rho_s/\rho$, taken to be 2.7, and $\rho_s$ is the sediment density, taken to be 2700 kg/m$^3$. For the specified values of $s$ and $\rho_s$, $W^*_c$ is given in the units of Figure 8 as $q_b = 10^{-5} \frac{o}{\tau_0}$, $\tau_0$ is the only parameter in (5) requiring specification, and $\tau_0 = 22.5$ Pa is fitted to the higher-quality 1993 gravel trap observations (Figure 8). Because of the very steep form of (5) at small transport rates the value of $\tau_0$ is insensitive to the measured value of transport rate, as long as it is smaller than the reference transport. (For the same reason, $\tau_0$ is relatively insensitive to the choice of $W^*_c$, as long as it is small.) With the fitted $\tau_0$, the resulting transport relation fits the larger 1992 gravel transport rates very well (Figure 8); $\tau_0 = 22.5$ Pa corresponds to $\tau_0 = 0.039$, which falls within a narrow range of values found to represent $\tau_0$ for the median size of a wide range of mixed-size sediments [Wilcock, 1992a].

A transport relation for sand is difficult to demonstrate because all observed transport rates are much larger than the reference transport rate. The total sand transport is consistent with a 3/2 power relation between $q_s$, and $T_0$ (Figure 8), as would be expected for fully mobilized transport at large values of $\tau_0/\tau_c$ [Yalin, 1977; Parker and Klingeman, 1982; Wilcock and McArdell, 1993].

Comparison of Entrainment and Transport Rates

Comparison of fractional entrainment $d_{w/D_0}$ and transport rates $q_{bi}/f_i$ at PB2 may be made by plotting each as a function of grain size (Figure 9). Fractional entrainment rates for SB3C are also shown on Figure 9a for comparison, although no observations of $q_{bi}$ were made. At PB2, the tracer observations at $Q = 76$ m$^3$/s and the gravel trap observations at $Q = 80$ m$^3$/s show that transport rates of all sizes are comparable, but very small, and little tracer entrainment occurs for any size. For $Q = 164$ m$^3$/s, gravel $q_{bi}/f_i$ are within a factor of 3 for all sizes (Figure 9b) and $d_{w/D_0}$ are within a factor of 2 (Figure 9a).
The degree of entrainment associated with different transport rates is of particular interest regarding $Z_r$, which is determined as the value of $Z_o$ producing a small reference transport rate, but serves as a surrogate for incipient motion conditions. The value of $\tau_i = 22.5$ Pa gives $\tau^*_i = 0.039$, corresponding to an exchange depth $d_{x}/D_{90} \approx 1.15$ (Figure 7a) and is associated with entrainment of all sizes, with some partial transport evident for the coarsest fractions (Figure 7b). This is an interesting and useful result. Two entirely different measures of incipient motion (the reference transport rate and mean scour to a depth of $D_{90}$) are shown to occur at nearly the same flow. At $\tau_0/\tau_i$ of 0.85 to 0.9, both entrainment and transport rates for the gravel fractions drop to essentially zero, as suggested by the form of the Parker [1979] transport relation (5). In the region $0.85 < \tau_0/\tau_i < 1.0$, a state of partial transport exists, and both $d_{x}/D_{90}$ and $q_{bi}/f_i$ increase very rapidly with $\tau_i$.

A more direct comparison between entrainment and transport rate may be made if transport is expressed as a cumulative mass per unit width ($\Sigma q_b = q_b T$, where $T$ is flow duration). This is plotted on Figure 10 as a function of fractional entrainment per unit area $E_{ai}$; both $\Sigma q_b$ and $E_{ai}$ are scaled by $f_i$. The curves shown on Figure 10 correspond to the variation with $\tau_0$ of $d_{x}/D_{90}$ and $q_b$ given in (4) and (5), respectively (Figures 7a and 8). $E_{ai}$ is calculated using (4) with $\tau^*_i = 0.031$, giving $d_{x}/D_{90}$ as a function of $\tau^*_i$, from which $E_{ai}$ is calculated as

$$E_{ai} = \left( \frac{d_{x}}{D_{90}} \right) D_{90} \rho_b (1 - n)$$

where $n$ is the bed porosity, taken to be 0.25, and the values $D_{90} = 92$ mm and $\rho_b = 2700$ kg/m$^3$ are used; $q_b$ is calculated as a function of $\tau_0$ using (5) with $\tau_i = 22.5$ Pa; $\tau^*$ is calculated using $D = 39$ mm, a size slightly larger than $D_{50}$ of the gravel portion of the bed. For a specified $\tau_0$, calculation of $\Sigma q_b$ requires specification of $T$. Two values of $T$ (4 and 16 days) are shown on Figure 10; the smaller value corresponds to the $T$ associated with the smaller values of $E_{ai}$. The apparent effect of $T$ decreases at larger values of $E_{ai}$, for which a large difference in $\tau^*$ makes little difference in $d_{x}/D_{90}$ (Figure 7a) and therefore $E_{ai}$.

Although a match is evident between the curves and data on Figure 10, no physical significance is attached to this beyond a claim that the two independent measurements of grain motion are consistent. The variation of $\Sigma q_b$ with $E_{ai}$ at small transport rates is highly sensitive to the choice of $\tau^*_i$, $\tau_i$, and the value of grain size used to form $\tau^*$ in (5). Further, the form of the curves at small $E_{ai}$ is sensitive to the variation of $d_{x}/D_{90}$ with $\tau^*_i$, which is not well constrained by the data (Figure 7a). Indeed, the plotted relation is likely incorrect in this region because, in the limit of vanishing transport, $\Sigma q_b$ will result from single, rare grain displacements of short length, so that $\Sigma q_b$ will approach $E_{ai}$ multiplied by the length of an individual displacement, which should scale with grain size [Drake et al., 1988]. This is shown as $\Sigma q_b = E_{ai} D_{90}$ on Figure 10.

Despite the uncertainty in the relation between $\Sigma q_b$ and $E_{ai}$, the merit of Figure 10 is to show a direct comparison between the two independent measures of grain motion and the empir-
ical relations representing them. At large $\tau^*$ the asymptotic approach to a limiting value of $E_\infty$ of the order of 300 kg/m$^2$ is required by (4) and corresponds to $d_j/D_{90} = 1.7$. Figure 10 also implicitly demonstrates the influence on entrainment of time, which is often not considered in calculations of $\tau_0$ or $Q_c$. If the relation between $\Sigma q_t$ and $E_\infty$ is relatively insensitive to $\tau_0$, then the magnitude of $E_\infty$ will depend on both $\tau_0$ (which determines $q_{t0}$) and time (which determines $\Sigma q_t$), suggesting that the same degree of entrainment can be produced by different $\tau_0$ if the flow duration is suitably chosen. Further, Figure 10 suggests that the influence of time is negligible at large $\Sigma q_t$, for which $E_\infty$ is a constant, and becomes increasingly important at smaller $\Sigma q_t$.

Summary and Conclusions

Trial reservoir releases on the Trinity River, California, provided an unusual opportunity to observe gravel entrainment under structured conditions of steady flow preceded and followed by very low flows, which permitted access to the river bed. Estimates of local shear stress $\tau_{0l}$ were made from observations of local depth-averaged velocity. A large lateral variability in $\tau_{0l}$ was observed, even for sections of relatively simple, trough-like geometry in a nearly straight reach. More critically, the variation of $\tau_{0l}$ with discharge differed from that of the section-average $\tau_{0b}$, demonstrating that the latter cannot be modified by a simple coefficient to give $\tau_{0l}$ and therefore estimate local entrainment and transport.

Large tracer gravel installations that incorporated the bed surface and subsurface were used to measure the entrained proportion of the bed surface, the entrainment of individual grain sizes, and the depth of scour. A scaled exchange depth, $d_j/D_{90}$, was found to vary consistently with a Shields parameter $\tau^*$ defined using $\tau_{0l}$ and the median size of the gravel fraction of the bed material. A threshold value of $\tau^*$, $\tau_{th}^* \approx 0.031$, is associated with the onset of measurable entrainment. $d_j/D_{90}$ increases rapidly at small $\tau^*/\tau_{th}^*$ and asymptotically approaches a limiting value of $\approx 1.7$ at large $\tau^*$, suggesting a limiting exchange depth for plane-bed transport of slightly less than $2D_{90}$. Mean entrainment to the base of the bed surface layer, $d_j/D_{90} \approx 1$, occurs at $\tau^* \approx 0.035$, a typical value for incipient motion in well-controlled experiments in smaller channels.

The transition from complete immobility to entrainment of the entire surface occurs over a narrow range of $\tau^*$ of the order of 10–15%. This range is characterized by partial transport, in which only a portion of the grains of a given size on the bed surface are mobile. For $d_j/D_{90} > 0.1$, all sizes are entrained and fractional exchange depths fall within a factor of 2, with the larger sizes tending to have the smallest entrainment.

Transport rates were measured with gravel traps for one flow with small transport rates and with bed load samplers at three discharges. Different transport relations are observed for the fine and coarse portions of the bed using a size boundary of 8 mm, which corresponds approximately to the threshold between matrix and framework sizes in the bed. Measurable gravel transport begins at a higher $\tau_0$ than that for sand, increases more rapidly with $\tau_0$, and reaches comparable transport rates at $\tau_0$ approximately 70% larger than that producing the onset of measurable gravel entrainment.

Gravel transport rates can be represented by the Parker [1979] transport relation, the only fitted parameter for which is the reference shear stress, a surrogate for the critical shear stress for incipient motion. A dimensionless reference shear stress of $\tau_{1g}^* = 0.039$, which is consistent with observations of $\tau_{1g}^*$ in smaller channels, was determined using very small transport rates observed in the gravel traps.

Observations of entrainment and transport were found to be consistent and provide an indication of the degree of entrainment associated with small transport rates. Both entrainment and transport rates for the gravel fractions drop to essentially zero for $\tau_{1g}^*/\tau_{th}^* < 0.1$, a state of partial transport exists, and both $d_j/D_{90}$ and $q_{wfl}$ increase very rapidly with $\tau_0$. The reference transport rate of $\tau_{1g}^*$ corresponds to grain exchange to a depth slightly greater than $D_{90}$ and is associated with entrainment of all grain sizes, with partial transport evident for the coarsest fractions. Two entirely different measures of incipient motion (mean scour to a depth of $\approx D_{90}$ and the reference transport rate) are observed to occur at nearly the same flow.

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