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A THERMOMECHANICAL ANALYSIS OF AN ULTRASONIC BONDING
MECHANISM

by

Chunbo (Sam) Zhang

A dissertation submitted in partial fulfillment
of the requirements for the degree

of

DOCTOR OF PHILOSOPHY

in

Mechanical Engineering

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2011

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Abstract

A Thermomechanical Analysis of An Ultrasonic Bonding Mechanism

by

Chunbo (Sam) Zhang, Doctor of Philosophy

Utah State University, 2011

Major Professor: Dr. Leijun Li
Department: Mechanical and Aerospace Engineering

A systematic experimental and numerical combined study of the thermomechanical bonding mechanisms in the ultrasonic welding (UW) process was conducted. A fully coupled thermomechanical finite element model has been built to fully understand the evolution and coupling between the in-process thermomechanical variables. The severe, localized, plastic deformation at the bond region is believed to be the major phenomenon for bond formation in ultrasonic welding. The influences of substrate dimensions on bond formation were studied and explained with an analytical vibration model. The formation of banded and cyclic stress-strain maxima in the substrate was found to be caused by superposition of vibrations. A push-pin type, combined experimental and numerical, method has been developed, validated, and applied to quantitatively determine the bond strength of UW parts. The best bond strength produced using the set of process parameters in this study was 75% of the ultimate tensile strength of the base material (Al3003-H18). Effects of UW parameters (normal pressure, vibration amplitude, and travel velocity) on bond strength have been characterized. Due to the weak vertical bond strength of UW parts, the pressurized post-weld heat treatment (PWHT) approach is originally proposed to improve the bond strength. The results show that the modified bond strength, up to 96% of the strength

of the base material, can be achieved under the optimum parameters of 2.5 MPa pressure, 450 °C temperature, and 1.5 h time.

(121 pages)

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Chunbo (Sam) Zhang

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Chapter 1

Introduction

1.1 Ultrasonic Welding Description

Ultrasonic welding (UW), as a solid-state joining process, uses an ultrasonic energy source (usually with a frequency of 20 kHz or above) to induce oscillating shears between the faying surfaces to produce metallurgical bonds between a wide range of metal sheets [1,2], thin foils [3], semiconductors [4], plastics [5], glass [6], and ceramics [7]. In contrast to traditional fusion welding processes, ultrasonic welding has several inherent advantages [3,8] derived from its solid-state process characteristics, and has been in use as a versatile joining method in the electronics, automotive, and aerospace industries since the 1950s.

Recent combination of ultrasonic metal seam welding and CNC milling has resulted in a new additive manufacturing process known as ultrasonic consolidation (UC) (Figure 1.1) [9,10]. By continuously welding layers of metal foil to previously deposited material, during which the profile for each layer is created by contour milling, UC is able to build-up complex, multi-functional 3-D objects. Objects with complex internal features, objects made up of multiple materials, and objects integrated with wiring, fiber optics, sensors and instruments can thus be directly fabricated [3,11,12]. Surface contaminants, such as oxides, are believed to be fractured and displaced, and atomically clean surfaces are brought into intimate contact under modest pressures. Researchers are still struggling to understand the complicated interaction of cyclic motion, large deformation, friction-like shear, localized thermal effect and oxides/contaminants dispersion on the interface.

Most conventional joining processes use some form of liquid-to-solid material phase transformation to achieve the transition from a feedstock form to a finished component. This transformation, and the elevated temperatures associated with melting, put practical limits on the range of materials that can be deposited, prohibits embedding of temperature-

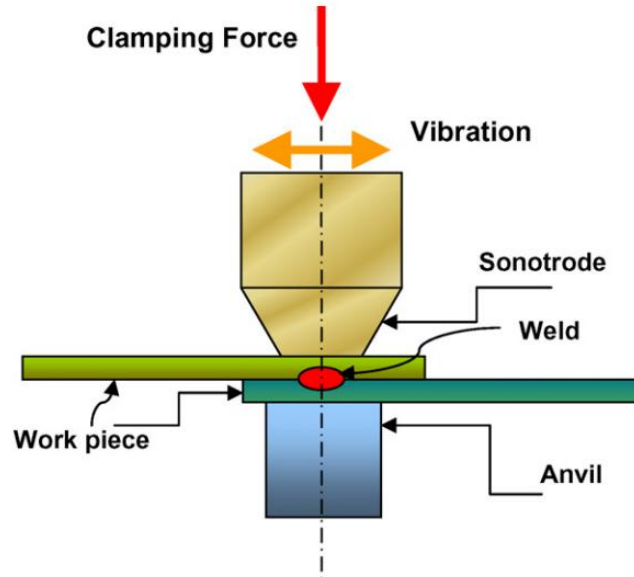


Fig. 1.1: Schematic of ultrasonic welding process.

sensitive devices. As a result, ultrasonic welding has several inherent advantages over other additive manufacturing and solid freeform fabrication processes [3, 8, 12, 13], including: (1) Temperature-sensitive devices (e.g. electronics, sensors, actuators, etc.) can be embedded within the structure, as it is built, without thermal damage. (2) Solid state processing facilitates retention of nonequilibrium microstructures produced during prior processing. (3) No atmosphere control is required to inhibit molten metal oxide formation. (4) Residual stresses and the resultant dimensional distortion may be reduced due to the lack of liquid-solid transformations. (5) Higher deposition rates are achievable at a lower overall energy consumption. (6) Dissimilar metals can be joined without melting mixing of the materials.

One unique aspect of UW is that highly localized plastic flow around embedded structures (e.g. ceramic fibers and wire meshes) is possible, resulting in sound physical/mechanical bonding between the embedded material and the matrix material [12–14]. Although the exact mechanisms by which this occurs are only just beginning to be investigated, this capability can be utilized in several ways, including: (1) manufacture of fiber-reinforced metal matrix composite with structural fibers for localized stiffening, (2) embedding of optical

fibers for communication and sensing, (3) embedding of shape memory fibers for actuation, and (4) embedding of wire meshes for planar or area stiffening.

1.2 Bonding Mechanisms of UW

Although numerous researchers have been studying the bonding mechanisms of ultrasonic welding for over 50 years, the process is still arguably the least understood welding process. Various mechanisms have been proposed for ultrasonic welding including inter-diffusion, re-crystallization, plastic deformation, work hardening, breaking of contaminant, generation of heat by friction and plastic deformation, and even melting [15, 16]. Diffusion has been observed at the interface between copper and aluminum welds for an extended period of welding time [17, 18]. And it has also been found that diffusion occurs along grain boundaries rather than in the bulk of the material [19]. Kreye [20] examined the microstructure at the weld interface with TEM and claimed that the very small grain size observed in a thin layer could only be explained by melting and solidification. However, Harthorn [21] and Heymann [22] concluded that neither diffusion nor re-crystallization could be responsible for the joint formation of ultrasonic welding after comparing low frequency vibration welding with ultrasonic welding of aluminum and examining the copper and soft iron ultrasonic welding. A great deal of plastic deformation and metal flow occur across the interface, and flow lines, evidence of extensive plastic deformation, are visible in the bond zone [3, 14, 23]. For the ultrasonic welding of aluminum foil, plastic flow occurs in a narrow interfacial zone about 10-20 microns in width [3]. In this region, new subgrain structures form across the bond zone. When the relative motion at the beginning of the welding cycle cleans the surfaces and plastically deforms asperities [24], microwelds – areas in which the friction exceeds the flow stress level of the material and plastic metal flow has started – occur immediately between points of contact of the adjacent surfaces, and spread out until a sufficient weld area is built up [15, 16]. Zhou et al. have investigated the effects of process parameters on bond formation in thermosonic gold ball bonding on a copper substrate at ambient temperatures with scanning electron microscopy [25]. They concluded that a relative motion existed at the bonding interface as microslip at lower pow-

ers, transitioning into gross sliding at higher powers. Researchers in a variety of fields have observed that ultrasonic excitation of metals can produce an apparent reduction in the yield strength, and enhancement in the plastic flow of metals [26,27].

Due to the difficulties in definitively characterizing the ultrasonic welding process, the mechanism(s) for bond formation in UW process are still under debate. The challenges for measurement approach and instrumentation come from the fact that: (1) the contact surface where the bonding occurs is invisible to observers; (2) the measurement devices or sensors tend to damage the contact surface and interfere with the bonding process; (3) ultrasonic bonding is highly localized (a few millimeters), and transient (typically 20 kHz vibration); and (4) the high spatial resolution (micron-scale vibration amplitude) is required. To overcome those obstacles and to gain insights on the UW process, numerous researchers have developed analytical [2,14,28] and numerical models [10,29–36].

1.3 Modeling of UW

During the last decade, ultrasonic welding has become a popular technique for joining thermoplastic polymers. Many studies were conducted modeling the process. Senchenkov et al. [37,38] studied the problem of vibrations and the dissipative heating of a viscoelastic prism by a waveguide. Benatar and Gutowski [39] modeled the ultrasonic welding of thermoplastics using a five-part model that included mechanics and vibration of the parts, viscoelastic heating, heat transfer, flow and wetting, and intermolecular diffusion. Roylance et al. [40] outlined numerical simulation methods useful in understanding and developing UW for polymers. Verderber [41] implemented algorithms in an explicit FE analysis code applicable to the unique features of ultrasonic welding analysis. Senchenkov and Zhuk [42] studied the two-dimensional problem of planar oscillations of plates under cyclic loading. The model developed was applied to simulation of vibration of sonotrodes for ultrasonic welding of plastics. A similar study by Mikhailenko and Franovskii [43] proposed a two-dimensional FE model that linked thermoelectric processes in acoustic systems.

Numerical modeling of ultrasonic welding of metals and alloys began recently, with the development of emerging new technologies that employ UW in ultrasonic consolidation

[3,44,45]. Compared with ultrasonic welding of polymers, ultrasonic welding of metals seems to be more challenging to model. In metals, vibration and heat have effects on dislocation dynamics, metallurgical transformations, and associated variations in thermomechanical properties [46].

Gao and Doumanidis [14,28] analyzed the mechanics of metal ultrasonic welding. They developed a 2-D, quasi-static/dynamic, elasto-plastic numerical model of the stress/strain field by FE analysis. In the study by de Vries [2], mechanics-based models were developed, along with a model for temperature generation and effect on mechanical properties of welded material. His models were capable of calculating the interface forces that were verified experimentally. Siddiq and Ghassemieh [35,36] studied the change in the friction work at the weld interface by simulating UW of metals based on a phenomenological material model, which incorporated both surface (friction) and volume (plasticity) “softening” effects. Ding et al. [33,34] analyzed deformation and stress distributions in wire and bond pad during ultrasonic wire bonding using 2-D and 3-D finite element models. A coupled temperature-displacement FE analysis performed by Huang and Ghassemieh [31] found the oscillation of stress in the substrate to lag behind the ultrasonic vibration by about 0.1 cycle of ultrasonic wave. Yadav and Doumanidis [32] performed the FE analysis of two layers of aluminum foil subjected to certain welding conditions - with surface contact resistance calibrated by experiments. The results revealed a moderate temperature rise that was believed to be sufficient for metal bonding via vacancy diffusion.

1.4 Result Summary of Previous UW Research

1.4.1 Vibration

The vibration of the sonotrode with constant pressure produces cyclic sliding at the bonding interface, along with periodic variation of the elastic/plastic strain/stress state, as illustrated in Figure 1.2 [28]. The resultant shear stress at the interface is helpful to break the oxides/contaminants and to generate intimate contact between the faying surfaces. The overall effect of ultrasonic vibration of the bonding interface can be summarized by

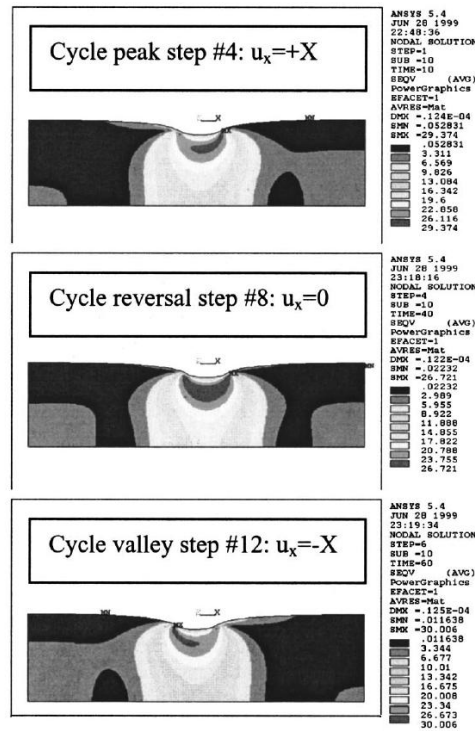


Fig. 1.2: Equivalent stress field ϵ_{eq} at three steps of one ultrasonic cycle.

the so-called “acoustic softening” of metals. Based on a phenomenological model, the stress required to initiate metal plastic deformation decreased significantly with ultrasonic excitation. The material softening effect has been studied by evaluating the sonotrode displacement in the height direction, illustrated in Figure 1.3 [31]. It is found that the sonotrode moves downward continuously during the whole UW process, under a constant load on the sonotrode. A similar effect was observed experimentally. This softening effect becomes more significant if the friction coefficient of the interface is higher.

1.4.2 Friction Work and Temperature

For those areas with the slipping contact condition, the shear stress results in the generation of friction heat. The friction work is affected by the vibration amplitude, illustrated in Figure 1.4, but has a non-linear relation with the applied pressure [36]. The roughness of contact surface decreases under the existence of cyclic shear stress from sonotrode vibration

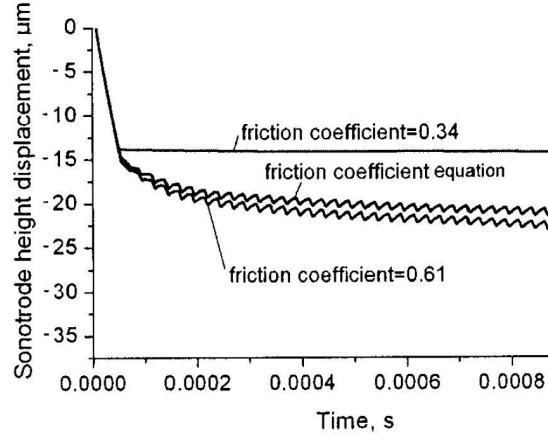


Fig. 1.3: Displacements of the sonotrode versus time along the height direction using different friction models. The external force applied on the sonotrode is 80N. Friction coefficient equation is defined as: $\mu = \frac{E}{p}[\epsilon_{xx}(t) + \nu]$, where E and ν are the modulus and Poisson ratio, and p is constant pressure.

so as to facilitate the interface ultrasonic bonding.

Due to the heat generation from friction work, the temperature on the bonding interface increases during UW. The generated heat can dissipate quickly across the entire contact surface for aluminum alloys. Some simulation results show the maximum temperature to be located at the foil/sonotrode interface (Figure 1.5), resulted from severe plastic deformation [36]. These simulation results seem to match the experimental observation of temperature measurement by high-speed thermal camera [2]. The consequence of temperature rise is to reduce the stress state of the contact interface by decreasing the material properties, and the friction coefficient. Thus, the friction work decreases with temperature rise.

1.4.3 Plastic Deformation

A large von Mises plastic strain may result in bonded areas. Plastic deformation is believed to affect the bond formation in the following ways: (1) it will disperse the surface oxide layer; (2) it will generate heat to make the atoms more thermally active; (3) it will deform the crystals and grains on the interfaces and generate a high dislocation density, which will also provide potential energy for grain boundary migration. Therefore, the plastic

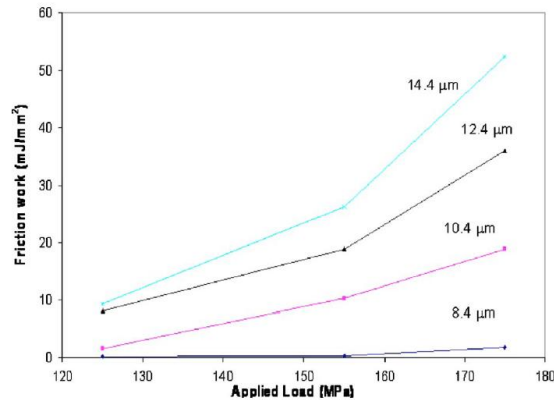


Fig. 1.4: Friction work at foil/substrate interface, velocity = 27.8mm/s.

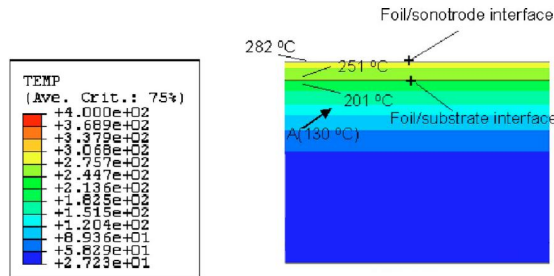


Fig. 1.5: Temperature in the weld specimen (amplitude = 8.4 μm, velocity = 27.8mm/s and pressure = 125MPa).

deformation from the simulation model can be quantitatively correlated with the bonded area (bond strength) from experiments.

Another quantitative approach for building the relationship between the simulated and experimental data is to measure the dimension of dynamic recrystallized grains. From experiments, it is found the crystals may have been dynamically recrystallized into nano-sized grains during processing (Figure 1.6) [13]. At the beginning of bonding, the plastic strain is relatively small. The highest plastic strain occurs at the edge of the contact surface. Subsequently, the plastic deformation in the center of the contact surface exceeds the level of plastic strain at the edge and becomes the highest. This phenomenon is observed from the fracture surface (Figure 1.7) [2].

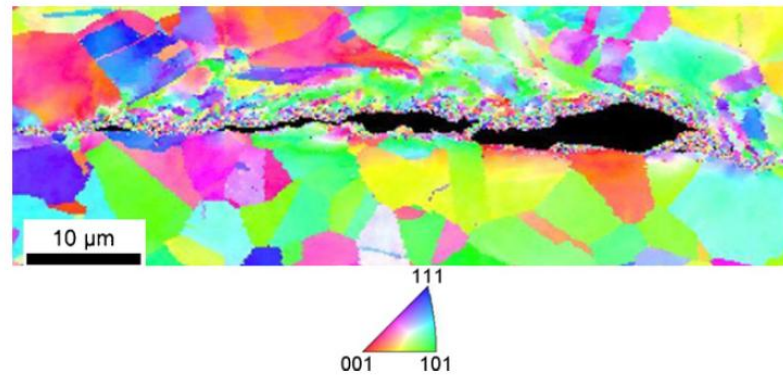


Fig. 1.6: Inverse pole figure of an unconsolidated portion of the Ni-Ni interface. Note the extremely fine grains that are present along the defect boundaries.

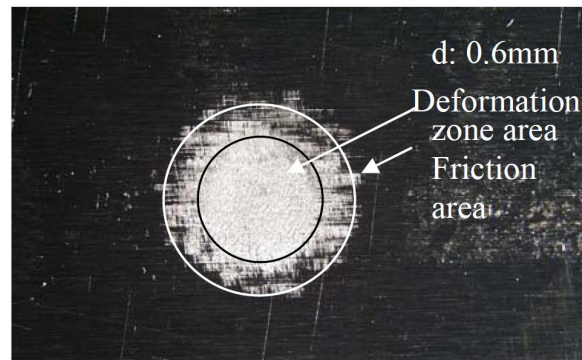


Fig. 1.7: Fracture interface of Al 6061 parts by UW after peel apart.

1.4.4 Bond Strength

The peeling test has been widely used to measure the bond strength of UW-made laminated structures by examining the peak load [47]. The bond strength is significantly influenced by process parameters. The bond strength increases with the vibration amplitude and preheat temperature, decreases with the sonotrode traveling speed, but shows a non-linear relation with normal pressure [48, 49].

Indicators from both experiment and simulation are employed to predict the true bond strength. Linear weld density (LWD), defined as the length of bonded interface divided by the total interface length under consideration (expressed in percentage), has a good linear relation with bond strength (Figure 1.8), but a significant increase of LWD only resulted

small increase in the bond strength [47]. Friction work from simulation has also been used to estimate the bond strength of UW parts [35]. The drawback of friction work as an energy bond strength indicator is the neglect of contributions of plastic heat and vibration energy.

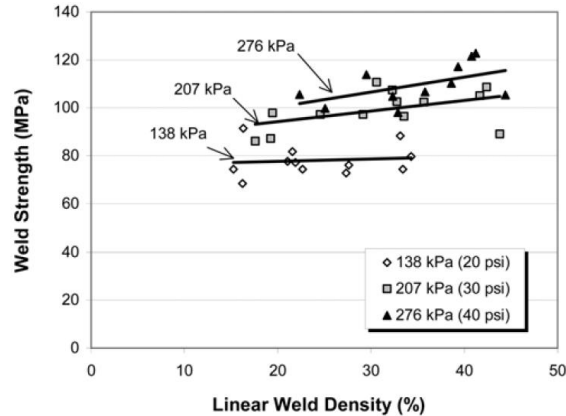


Fig. 1.8: Correlation of weld strength with LWD of welded Al 6061 specimens.

1.5 Research Plan and Potential Impact

Based on the literature survey and preliminary study of this thesis work, metals under ultrasonic bonding simultaneously experience the complex and interactive processes of friction, heat generation, material softening and diffusion under the elevated temperature, and material flow and hardening. Therefore, a fundamental understanding of this ultrasonic thermomechanical bonding phenomena requires a systematic approach covering all these aspects. However, previous experimental studies of UW, based on the post-weld observations and analysis, resulted in different, and even contradictory conclusions regarding the bonding mechanisms because of the “black box” approach in the understanding of bond formation. The reported models of UW focused on either the single mechanical/thermal behavior, or one-way coupled effect of ultrasonic bond formation. These models are still unable to fully interpret the physical, thermomechanical bonding phenomena. No direct relevance of the model was provided for a mechanistic understanding. It is also notable that none of the reported modeling work provided validation, presumably due to the difficulties of strain/stress, temperature, and displacement measurements in real-time during

ultrasonic welding [23,31].

Therefore, there is no universally-acknowledged theory for bond formation of UW currently. This work first identified the ultrasonic bonding process as a thermomechanical coupled problem, and proposed that a bond mechanism study has to be focused on a complete understanding and interpretation of the thermomechanical coupling history at the bonding interface. This work has developed an integrated set of innovative methods for this purpose. A 3D numerical model for simulating the UW process has been developed. A scheme for quantitatively studying the fully thermomechanical coupling effects, especially the calculation of the effect of plastic deformation, has been designed and realized. The data of temperature-dependent friction coefficient have been measured. The approach of quantitative bond strength determination and potential of bond strength improvement with a pressurized post-weld heat treatment have been proposed. This work has successfully solved the problem of non-linear multi-physics analysis and the function of plastic deformation work in UW, and has expanded the capabilities of current commercial FE software packages. With respect to the theoretical contributions to the bonding mechanism of UW, the fully thermomechanical coupling effect during bond formation has been qualitatively described, the roles and contributions of in-process thermomechanical variables to the ultrasonic bonding have been identified and validated with experiments, and finally a possible bonding mechanism for ultrasonic metal welding is proposed. In the study of influences of substrate dimensions on the ultrasonic bonding, the travel, reflection and interference of ultrasonic waves within the substrate have been found and proved to play a critical role in ultrasonic bond formation, which may even cause an interface de-bonding problem in the worst case.

This thesis work has followed the following plan of studies: (1) A series of simulation will be conducted for different process parameters. Then the simulation results will be analyzed so that the relationship among the process parameters (pressure, amplitude, traveling speed), in-process variables (contact surface shear stress, sliding distance, heat generation, temperature and elastic vibration), and the final result (plastic deformation)

will be revealed. (2) The push-pin tests will be designed and conducted for each simulation condition to obtain the bond strength. A detailed analysis of fracture surface will be conducted to find out the local/bulk distribution of the bonded area-ratio and bond strength. (3) The correlation of plastic deformation and bond strength will be established based on a detailed and quantitatively study of local/bulk plastic deformation and bond strength data. (4) The influence of substrate geometry on bond formation will be studied. (5) The bonding mechanism for ultrasonic metal welding will be proposed. The bond strength will be predictable if given a set of process parameters. (6) The pressurized post-weld heat treatment tests will be conducted to explore the potential bond strength improvement method.

Welding process simulation on a thermomechanical fully coupled model provides quantitative information on the bonding interface and overcomes the difficulty in data acquisition in the real manufacturing process, particularly for a process in which the phenomena are localized, non-equilibrium, and coupled. Bond strength obtained by a specially designed experiment serves as the strong verification of the simulation work. This thesis work will have the following deliverables:

1. The complete and accurate temperature-dependent material properties of Aluminum 3003-H18, including the friction coefficient, Young's modulus, stress-strain curves;
2. A 3D thermomechanical fully coupled finite element model;
3. The systematic quantitative identifications of transient histories, distributions and interactions of in-process variables, which include normal stress, shear stress, plastic deformation, heat generation and temperature;
4. A general approach of bond strength determination, and the identifications of bond strength value for different process parameters;
5. The effects of substrate geometry on ultrasonic bond formation, and the ultrasonic energy generation and transfer;
6. The comprehensive insights of thermomechanical coupled phenomena and ultrasonic bonding mechanism;

7. The potential of bond strength improvement and the optimum conditions of pressurized post-weld heat treatment.

As a result of this work, a more complete understanding of the physical phenomena that govern ultrasonic joining has been developed. This will lead to more effective control of ultrasonic welding operations. As a solid-state process with many inherent benefits over other joining technologies, ultrasonic welding has the potential to dramatically affect a number of key industries, including defense, aerospace, automotive, and general manufacturing. The use of ultrasonic welding to form fiber-reinforced metal matrix composites from engineering material matrices, to form dissimilar metal thermal management devices with optimized thermal expansion and conductivity, or to form complex components with embedded functionality, has the potential to provide design breakthroughs for the electronics, aerospace and transportation industries, among others. This proposal also provides a general solution strategy concerning the bonding of coupled multi-field and non-melting process, such as friction-stir welding.

Chapter 2

Methodology and Theoretical Background

2.1 Methodology of UW Modeling

2.1.1 Fully Coupled Thermomechanical Model

Based on the above literature review, it is seen that the reported models of metal UW focused on either the single mechanical/thermal behavior, or one-way coupled effect of ultrasonic bonding formation. However, the ultrasonic bond formation is believed to be a complex physical thermomechanical coupled bonding phenomena. This thesis work first considered the mechanism of ultrasonic bonding as a thermomechanical coupled problem and tried to numerically interpret this phenomena with a new 3D fully coupled thermomechanical non-linear FE model based on ANSYS. Compared with previous numerical models of UW, this developed model is capable of analyzing the non-linear transient fully coupled effects of shear stress, friction heat flux, plastic deformation and plastic heat flux simultaneously. The heat generation by plastic deformation is first introduced into the bond mechanism study and UW numerical modeling.

Figure 2.1 (a) shows the 3-D FE model for ultrasonic welding, with the sonotrode located at the center of the baseplate. Since this model has two symmetric planes and a large difference of dimensions exists between the foil (0.1 mm thick) and baseplate (25 mm thick), the model with quarter size shown in Figure 2.1 (b) is used. The contact region has a refined meshing shown in Figure 2.1 (c). This reduction of dimension size has been validated by a mesh study. The first-layer welding process is simulated by this model, so that simulation results are not affected by the influence of substrate height on bond formation. The dimensions of actual contact areas for different process conditions have been measured and contact elements are defined at contact surface.

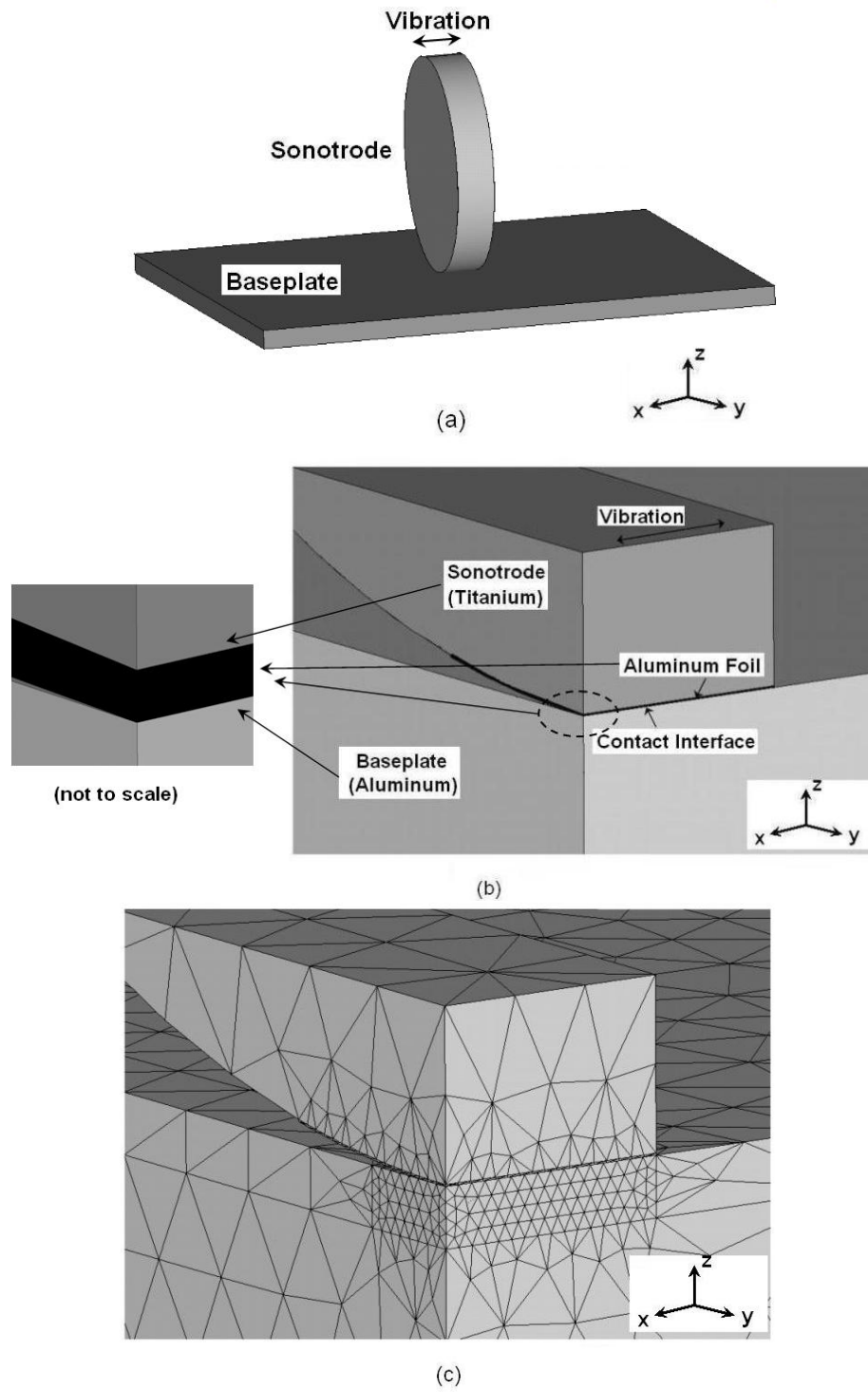


Fig. 2.1: The 3-D fully coupled thermomechanical FE model for ultrasonic welding.

The multi-physics analysis of commercial FE softwares in particular the GUI is not able to satisfy the needs of simulating the fully thermomechanical coupled problem of UW. Their limitations are from two aspects: 1) The multi-physics analysis has to proceed under the condition of linear analysis; 2) The heat flux from plastic deformation can not be directly calculated in the current FE software. Therefore, a new customized algorithm is developed to extend the functions of material non-linear analysis (plastic deformation) and related thermal effects (heat generation from plastic deformation) to the thermomechanical multi-physics analysis. The fully coupled thermomechanical analysis algorithm for the above model has been developed. Definitions of variables and the iteration process for determining the in-process variable matrices for each ultrasonic vibration cycle are shown in Figure 2.2.

(1) A thermal-structural analysis is conducted first to obtain the temperature field. Variables from the last step ($[\epsilon_{ts}]_{j-1}$, $[\sigma_{ts}]_{j-1}$, and $[T]_{j-1}$) are input as the initial conditions. Plastic heat $[PH]_{j-1}$ from the last step is applied as the thermal load. Materials properties $[MP]_{j-1}$ and boundary conditions $[BC]_j$ are also provided. The $[BC]_j$ include the fixed/symmetric constraints and vibrational motion of the sonotrode. The output parameters of this coupled thermomechanical analysis are $[\epsilon_{ts}]_j$, $[\sigma_{ts}]_j$, and $[T]_j$, which will replace the old values from the last step. Based on $[T]_j$, $[MP]_{j-1}$ is also updated. The friction heat flux is calculated from the shear stress and slide distance of the contact interface. The temperature field $[T]_j$ results from friction and plastic heat fluxes, as well as the heat transfer.

(2) A structural analysis is conducted second to obtain the stress and plastic deformation fields. The variables from the last step ($[\epsilon_s]_{j-1}$ and $[\sigma_s]_{j-1}$) are input as the initial conditions. Temperature $[T]_j$ is applied as the thermal load. The outputs of this analysis are $[\epsilon_s]_j$ and $[\sigma_s]_j$, which will replace the values from the last step. The increase of plastic deformation is used to calculate the plastic heat $[PH]_j$.

During numerical simulation, at each vibration cycle, most of material properties, such as modulus of elasticity, yield strength and friction coefficient, are temperature sensitive, so these varying material properties under the temperature of the previous time step must

be known and updated before getting the solution of the current cycle.

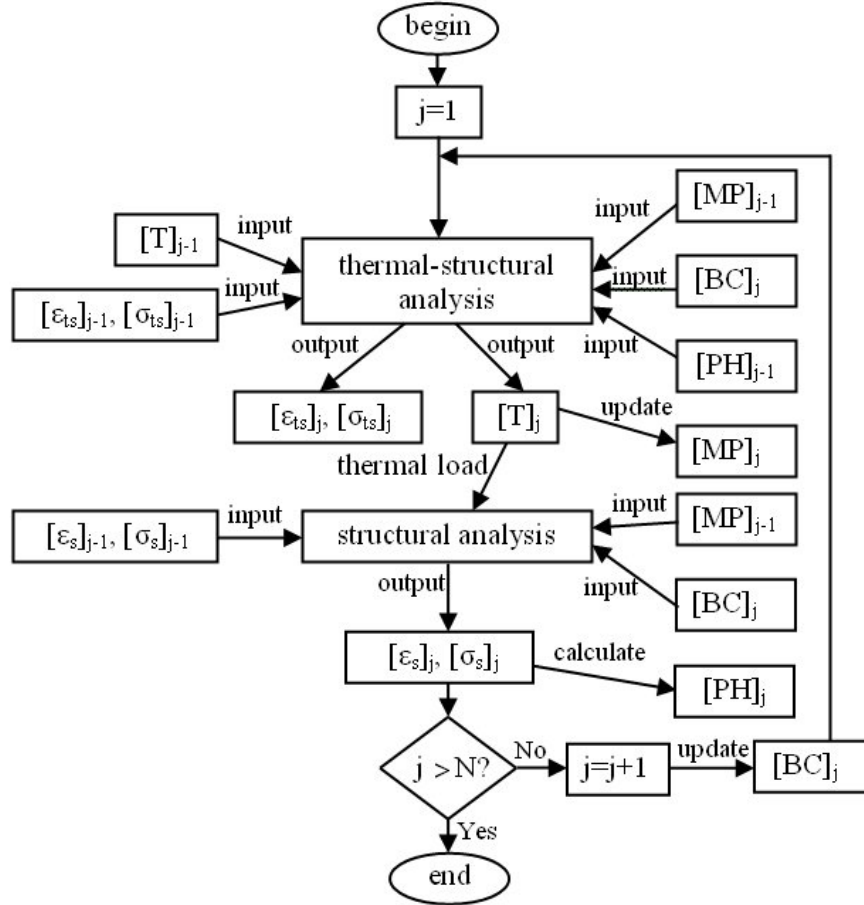


Fig. 2.2: Thermomechanical coupled analysis diagram for ultrasonic welding process: j is the j th number of vibration cycle; N is the total number of vibration cycle; $[T]$ is the temperature matrix; $[\epsilon_{ts}]$, $[\sigma_{ts}]$, and $[\epsilon_s]$, $[\sigma_s]$ are the strain and stress matrixes for thermal-structural and structural analysis, respectively; $[MP]$, $[BC]$ and $[PH]$ are the matrixes for material properties, boundary conditions and plastic heat flux.

2.1.2 Temperature-dependent Material Properties

Tensile tests have been conducted on the Gleeble 1500D thermomechanical simulator for measuring the temperature-dependent mechanical properties of aluminum foil (Al 3003-H18) used in this study. The dimensions of test samples are $22.0 \times 10.1 \times 0.1$ mm, and the tensile load rate is 1.2 mm/min. Specimens were heated to various temperatures and held for 3 min for thermal equilibrium, and tensile pulled. The results are listed in Table 2.1. It is seen that the properties of aluminum material are very sensitive to the temperature. At the preheating temperature (150 °C), the mechanical properties decrease to the half or even less of those at room temperature. The measured yield strength data of aluminum foil are different, but close to those from the reference. The composition and thermal properties for Al 3003-H18 are shown in Tables 2.2 and 2.3 [50].

Table 2.1: Temperature-dependent Mechanical Properties of Al 3003-H18

Temp. (°C)	Modulus of Elasticity (GPa)	Yield Strength (MPa)	Yield Strength* (MPa)
25	53.2	227.3	185
50	28.8	197.3	-
100	26.8	131.2	145
150	22.7	70.0	-
200	16.9	68.3	62
250	16.42	32.85	-
300	14.5	29	17
350	13.26	26.5	-

*: Data come from Ref. [50].

Table 2.2: Material Components Properties of Al 3003-H18 (%)

Al	Cu	Fe	Mn	Si	Zn
96.7-99.0	0.050-0.200	≤ 0.700	1.00-1.50	≤ 0.600	≤ 0.100

2.1.3 Friction Coefficient

As a friction-based welding process, there is no doubt that the friction coefficient at the faying interfaces plays a key role by affecting the stress/strain state and surface heat

Table 2.3: Thermal Properties of Al 3003-H18

CTE ($\mu\text{m}/\text{m}\cdot^{\circ}\text{C}$)	Specific Heat Capacity ($\text{J}/\text{g}\cdot^{\circ}\text{C}$)	Thermal Conductivity ($\text{W}/\text{m}\cdot\text{K}$)
25.1	0.893	155

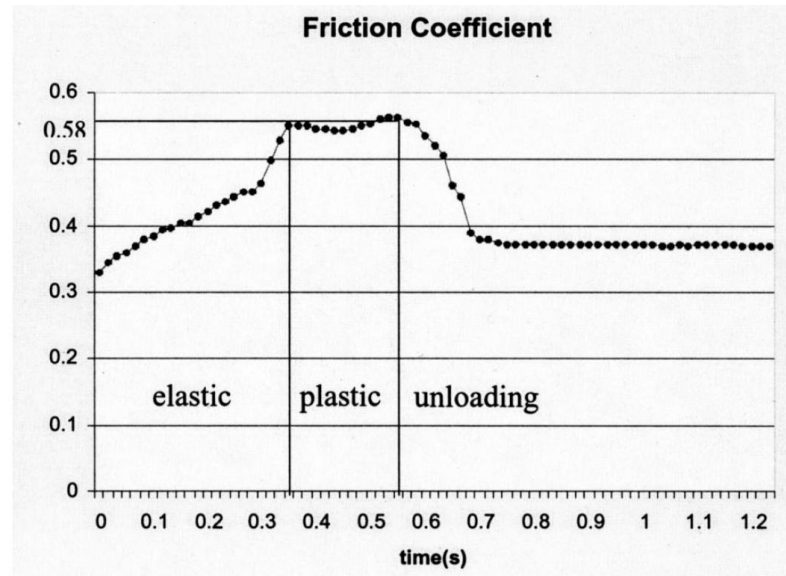


Fig. 2.3: Experimentally identified history of friction coefficient of Al 1100 foils during one weld.

generation in UW. It can be said that the ultrasonic bond formation stems from the shear stress on the interface, which is mainly determined by the friction coefficient, since the existence of shear stress on the interface generates the friction heat flux, and consequent plastic deformation and heat flux. Gao and Doumanidis [28] experimentally determined the friction efficient history of Al 1100 foils during one cycle (Figure 2.3). Siddiq and Ghassemieh [35,36] numerically described the friction coefficient varying with temperature and vibration cycle by curve-fitting the published experimental data. Based on Coulomb's friction law, the stick-slip contact condition can be identified with the relation of $|\tau_{fric}|$ and $\mu \cdot p$. If $|\tau_{fric}| < \mu \cdot p$, the friction is of the stick-type, and if $|\tau_{fric}| = \mu \cdot p$, the friction is of the slip-type.

Regarding the thermomechanical analysis of UW process, the friction coefficient varying

with the vibration cycle is believed to be caused by the elevated temperature on the bonding interface. However, the data of temperature-dependent friction coefficient is still unavailable, and the previous UW models ignored the effect of temperature on the friction coefficient. Therefore, a novel approach for measuring the Al-Al temperature-dependent friction coefficient has been designed and conducted on the Gleeble 1500D. With this method, the obtained temperature-dependent friction coefficient can improve the accuracy of the UW simulation.

Figure 2.4 shows the device setup. The fixed anvil is covered with foil of Al 3003-H18 and another foil is put between them. A normal force (N) is applied on the top anvil. When a certain uniform temperature is achieved by the heating and controlling system of the Gleeble, the middle foil is pulled out at a certain rate (sliding speed). The sliding speeds between 1 and 1000 mm/s are tested. The friction occurs on the top and bottom surfaces of the foil pulled out. Therefore, the Al-Al friction coefficient at the testing temperature can be calculated by $\mu = F/2(N + W)$, where F is the measured pulling force, and W is the top anvil's weight and pressure. Friction coefficient of Al-Al varying with temperature is shown in Figure 2.5. There is no significant effect of sliding speeds, in the range tested, on the friction coefficient. The model in this paper has included the material nonlinear effects by using the data in the temperature-dependent mechanical properties tables above. In the transient dynamic analysis, the mechanical properties for different locations on the contact surface are different depending on the local temperature and are input to the model for each step.

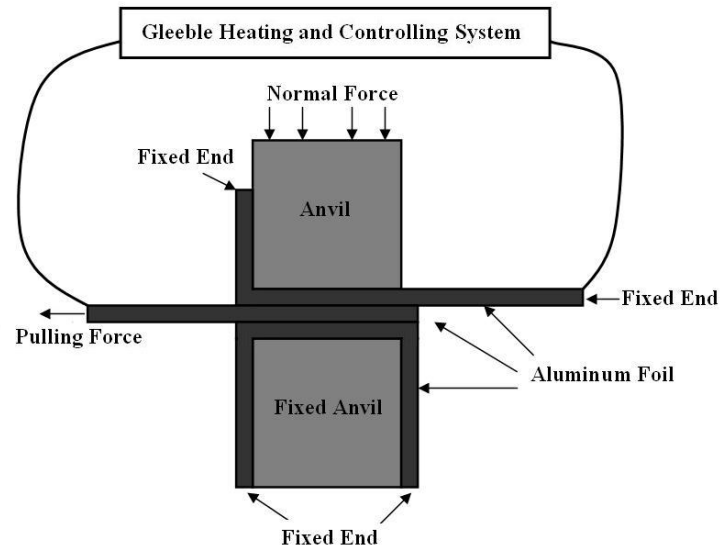


Fig. 2.4: Schematic of Al-Al friction coefficient measurement on Gleeble 1500D.

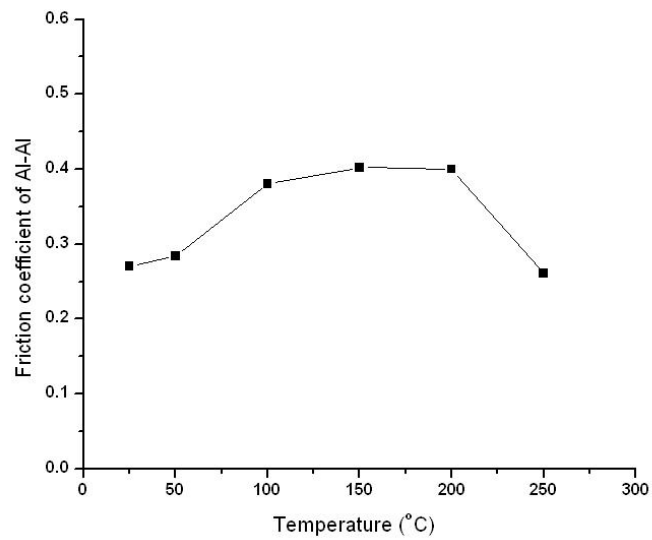


Fig. 2.5: Al-Al friction coefficient varying with temperature.

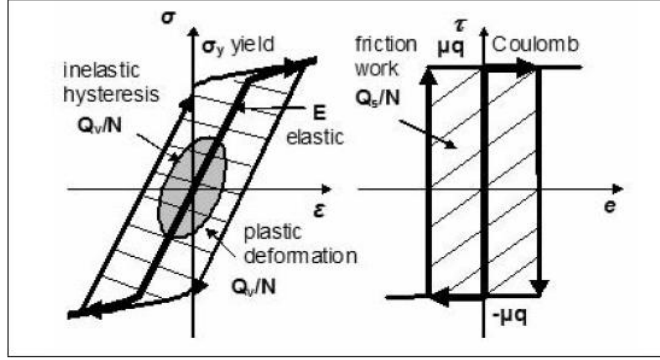


Fig. 2.6: Mechanical work done during each vibration cycle of UW.

2.1.4 Heat Generation in UW

The ultrasonic mechanical energy, subsequently converted to thermal energy, can be determined by the local cyclic stress (σ), strain (ϵ), friction shear (τ_x), and component slip (e_x) at the interface surface. Heat is generated (a) locally in the control volume, ($Q_v(x, y, z)$), by inelastic hysteresis and plastic deformation, and (b) at the interface surfaces, ($Q_s(x, y)$), by friction, during each cycle at the ultrasonic frequency, f , as shown in Figure 2.6 [2, 32, 33].

$$Q_v(x, y, z) = f \oint \sigma_{eq}(\epsilon) d\epsilon_{eq} \quad (2.1)$$

$$Q_s(x, y) = f \oint \tau_x(\epsilon_x) de_x = N \oint \mu \sigma_z(\epsilon_x) de_x \quad (2.2)$$

where σ_{eq} and ϵ_{eq} are equivalent stress and strain, and σ_z is normal compressive stress.

It is first tried to embed the plastic work in the numerical thermomechanical coupled analysis of UW and quantitatively study the role of plastic work in ultrasonic bond formation. The above equations for calculating heat generation assume the mechanical state as a function of location only, not considering time and temperature influences. This thesis proposed a new method of calculating the plastic work, which uses the plastic work rate (\dot{W}^{pl}), proportional to the plastic strain rate, to account for the effects of location, temperature

and time in UW.

$$\dot{W}^{pl}(x, y, z, t, T) = \sigma_j(x, y, z, t, T) \cdot \dot{\epsilon}_j^{pl}(x, y, z, t, T) \quad (2.3)$$

where $j = 1, 2, 3$. It is assumed that some of the plastic work converts to heat, while the rest is stored as energy of crystal defects accompanying plastic deformation. The fraction of the plastic work ($\dot{Q}^{pl}/\dot{W}^{pl}$) that is converted to thermoplastic heating (\dot{Q}^{pl}) is 0.33 for aluminum [51].

2.1.5 Initial and Boundary Conditions

The UC process parameters, including pressure load (1800 N), preheating temperature (150°C), vibration amplitude (16 μm), and vibration frequency (20 kHz), are incorporated in the model as boundary conditions. The pressure load and vibrational displacement are applied to the interface of foil and baseplate. A triangle wave is used to approximate the sinusoidal waveform of the ultrasonic vibration. The measured contact area under 1800 N is 4.7 mm along the direction of the sonotrode's moving, and the sonotrodes speed in the UC process is set at 28 mm/s. For a given point underneath the sonotrode, the maximum time of contact is approximately 0.16 s, or about 3000 cycles of vibration. Therefore, we assume a stationary sonotrode in the model, and the maximum number of cycles to be modeled is 3000. A fixed boundary condition for displacement is applied to the baseplate's bottom surface and all corners except the surface contacting the sonotrode. Symmetrical boundary conditions are used on two symmetrical planes. A uniform initial preheating temperature is applied to the parts around the bonding surface, including the sonotrode, baseplate, and foil. No convective heat loss from baseplate and sonotrode to air is assumed because the UC device is in an enclosed space in which there is no significant air flow. The radiation heat loss is also ignored due to the low temperature range.

2.2 Theoretical Background

2.2.1 Governing Equations of Thermomechanical Processes

In the modeling of UW processes, the governing dynamic equation of displacement (u) for a linear structure [52] is:

$$[M] \{\ddot{u}\} + [C] \{\dot{u}\} + [K] \{u\} = \{L\} \quad (2.4)$$

where $[M]$ is the structural mass matrix, $[C]$ is the structural damping matrix, $[K]$ is the structural stiffness matrix, \ddot{u} is the nodal acceleration vector, \dot{u} is the nodal velocity vector, $\{u\}$ is the nodal displacement vector, and $\{L\}$ is the applied load vector.

The basic constitutive equations between stress (σ) and strain (ϵ) are as follows:

$$\{\sigma\} = [D] \{\epsilon^{el}\} \quad (2.5)$$

$$\{\epsilon^{el}\} = \{\epsilon\} - \{\epsilon^{pl}\} - \{\epsilon^{th}\} \quad (2.6)$$

where $\{\sigma\}$ is the stress vector, $[D]$ is the elastic stiffness matrix which is a function of temperature, $\{\epsilon^{el}\}$ is the elastic strain vector, $\{\epsilon\}$ is the total strain vector, $\{\epsilon^{pl}\}$ is the plastic strain vector, and $\{\epsilon^{th}\}$ is the thermal strain vector.

Conductive heat transfer can be considered using the following governing equation:

$$\rho c \left(\frac{\partial T}{\partial t} + (\vec{v} \cdot \nabla) T \right) = \dot{q} + \nabla \cdot ((\vec{K} \cdot \nabla) T) \quad (2.7)$$

where ρ is the density, c is the specific heat, T is the temperature, t is the time, \vec{v} is the velocity vector for mass transport of heat, \dot{q} is the heat generation rate per unit volume, \vec{K} is the conductivity vector, and $\nabla = \frac{\partial}{\partial x} + \frac{\partial}{\partial y} + \frac{\partial}{\partial z}$.

Ultrasonic welding has been identified as a two-way coupled thermomechanical problem for: (1) the thermal field affects the mechanical field, because most of the material properties, such as modulus of elasticity, yield strength, and friction coefficient, are tem-

perature sensitive; and (2) the mechanical field (friction and plastic deformation) generates heat, which affects the thermal field. Therefore, modeling of the UW process demands a coupled-field analysis method. Because closed-form solutions for coupled-field equations are difficult to obtain, structural and thermal fields are generally treated separately in analytical modeling of UW [2, 14, 28].

Recently, commercial FE packages have expanded their capabilities in solving complex multi-physics problems; many researchers found it convenient to study involved processes that are not solvable analytically. In the thermomechanical coupled analysis, the FE matrix equation for mechanical and thermal fields, coupled by the thermoelastic constitutive equations, is as follows [52]:

$$\begin{bmatrix} [M] & [0] \\ [0] & [0] \end{bmatrix} \begin{Bmatrix} \{ \ddot{u} \} \\ \{ \ddot{T} \} \end{Bmatrix} + \begin{bmatrix} [C] & [0] \\ [C^{tu}] & [C^t] \end{bmatrix} \begin{Bmatrix} \{ \dot{u} \} \\ \{ \dot{T} \} \end{Bmatrix} + \begin{bmatrix} [K] & [K^{ut}] \\ [0] & [K^t] \end{bmatrix} \begin{Bmatrix} \{ u \} \\ \{ T \} \end{Bmatrix} = \begin{Bmatrix} \{ L \} \\ \{ Q \} \end{Bmatrix} \quad (2.8)$$

where $[M]$ is the element mass matrix, $[C]$ is the element structural damping matrix, $[K]$ is the element stiffness matrix, $\{u\}$ is the displacement vector, $\{F\}$ is the sum of the element nodal force, $[C^t]$ is the element specific heat matrix, $[K^t]$ is the element diffusion conductivity matrix, $\{T\}$ is the temperature vector, $\{Q\}$ is the sum of the element heat generation load and element convection surface heat flow vectors, and $[K^{ut}]$ is the element thermoelastic stiffness matrix ($[K^{ut}] = -\int_{vol}[B]^T\{\beta\}\{\nabla\{N\}^T\}d\{vol\}$), where $[B]$ is the strain-displacement matrix, and $[C^{tu}]$ is the element thermoelastic damping matrix ($[C^{tu}] = -T_o[K^{ut}]^T$, and T_o is the absolute reference temperature).

2.2.2 Material Behavior in UW

Yield Rule and Plastic Deformation

The von Mises yield criterion has been widely applied in plastic analysis [29, 32]. If the von Mises equivalent stress is less than the material yield strength at temperature, the stress state is elastic and no plastic strain is computed. If the stress exceeds the material yield strength, the plastic strain ϵ^{pl} is calculated by

$$\{d\epsilon^{pl}\} = \lambda \left\{ \frac{\partial F}{\partial \sigma} \right\} \quad (2.9)$$

where λ determines the amount of plastic straining, and F (yield function) is given by

$$F = |\underline{\sigma} - \underline{\alpha}| - (\sigma^0 + R) = 0 \quad (2.10)$$

where $\underline{\sigma}$ is stress tensor, $\underline{\alpha}$ is the back stress tensor due to kinematic hardening, R is the isotropic hardening term and σ^0 is the initial yield stress.

Thermomechanical Hardening Rules

The nonlinear isotropic hardening rule was presented by Lemaitre and Chaboche [53] and Huber and Tsakmakis [54]. The isotropic hardening (R), which describes the expansion of the yield surface, is defined as an exponential function of accumulated plastic strain.

$$R = A(1 - e^{-b\bar{\epsilon}^{pl}}) \quad (2.11)$$

where $\bar{\epsilon}^{pl}$ is the equivalent plastic strain, while A and b are material parameters to be identified. A is the maximum change in the size of the yield surface, and b is the rate at which the size of the yield surface changes with changing plastic strains.

A nonlinear kinematic hardening rule proposed by Armstrong and Frederick [55] has been used to capture nonlinear hardening behavior and the smooth transition from elastic

to plastic deformation. The evolution of back stress tensor ($\underline{\dot{\alpha}}$) is given by

$$\underline{\dot{\alpha}} = \frac{S}{\gamma}(\underline{\sigma} - \underline{\alpha})\bar{\epsilon}^{pl} - \gamma\underline{\alpha}\bar{\epsilon}^{pl} \quad (2.12)$$

where S and γ are the material parameters, which can be identified from cyclic testing. The γ term determines the rate at which the saturation value of kinematic hardening decreases with increasing plastic deformation. S is the kinematic shift of the yield surface.

Chapter 3

Discussions of Simulation Results and Bonding Mechanism

3.1 Thermomechanical Simulation Results with Friction Heat Generation Only

3.1.1 Effect of Friction on Temperature

Figure 3.1 shows the coordinate plane projections of temperature distribution on the contact surface (bond interface) at the 1500th vibration cycle. With friction heat, a uniform peak temperature state exists on the contact surface, and the temperature decreases away from the contact surface into the substrate. The temperature distribution shows a cylindrical symmetry about the center of the contact surface. The average temperature due to friction at the contact surface increases rapidly with the number of vibration cycles in the initial period of bonding. As the welding time is longer (higher number of cycles), the temperature rise slows down, and seems to approach a steady state, as seen in Figure 3.2. The increase of interface temperature is influenced by $q_{friction} - q_{loss}$, where $q_{friction}$ is the friction heat generation rate, and $q_{loss} = k\Delta T$ is the heat dissipation rate by conduction. At the beginning of UW bonding, the rate of friction heat generation is much higher than that of heat loss, because initially the temperature difference (ΔT) is low. As the temperature becomes higher due to friction heating, two competing factors become significant: On the one hand, the decrease of Young's modulus (Table 2.1) and friction coefficient (Figure 2.5) will slow down the friction heat generation; on the other hand, the greater temperature difference enhances faster heat dissipation. Therefore, the temperature increase slows down. A balance between $q_{friction}$ and q_{loss} can be achieved eventually, giving rise to $q_{friction} - q_{loss} = 0$, i.e., a thermal steady state. Physically, if the heat loss is faster than the friction heat generation and the interface temperature decreases, the higher values of Young's modulus and friction coefficient at the lower temperature will automatically increase the

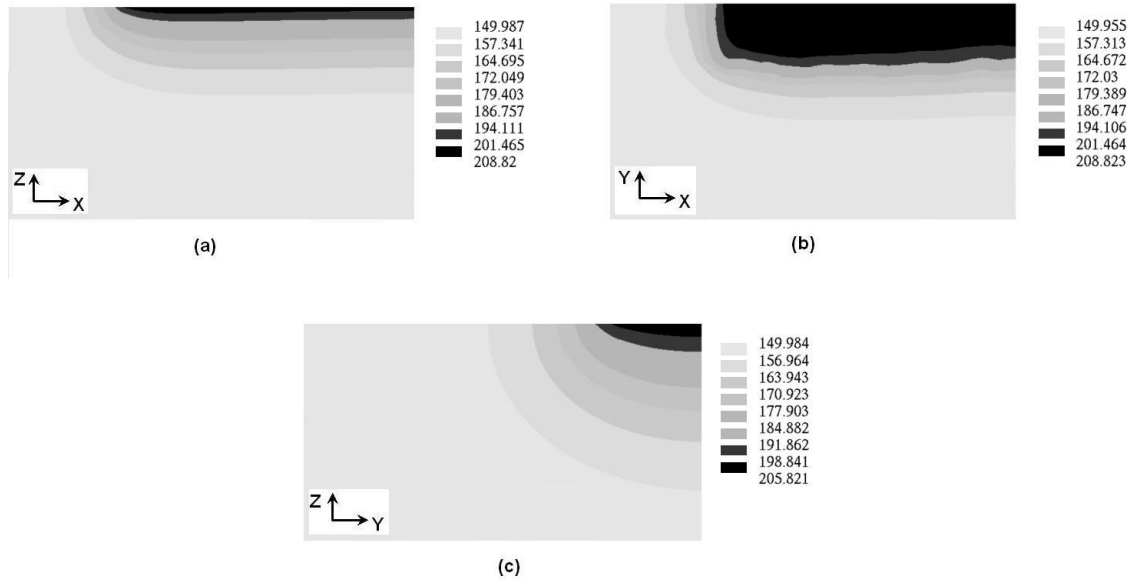


Fig. 3.1: Distribution of temperature at the 1500th vibration cycle: (a) front view, (b) top view, and (c) side view.

$q_{friction}$ until a new dynamic heat balance is achieved.

3.1.2 Effect of Friction on Plastic Deformation

Three paths (A-B, A-C, and A-D) have been selected to show the three-dimensional distributions of von Mises strain around the contact surface (Figure 3.3a) where the bonding process happens. Figure 3.3b shows the von Mises strain distributions along path A-B (the vibration direction) as a function of the number of vibration cycles. At the beginning of bonding (50th cycle), the plastic strain is relatively small. The highest plastic strain occurs at the edge of the contact surface near point B. By the 350th cycle, the plastic strain in the center of the contact surface exceeds the level of plastic strain at the edge. After the 500th vibration cycle, the value and distribution of plastic deformation reach a steady state, when there is no change in von Mises strain with further increase of the vibration cycle. The maximum plastic strain is located at point A, the center of the contact surface. A wavy-shaped plastic strain distribution caused by the cyclic vibration is apparent. The reason for such a characteristic wave shape distribution of deformation has been studied

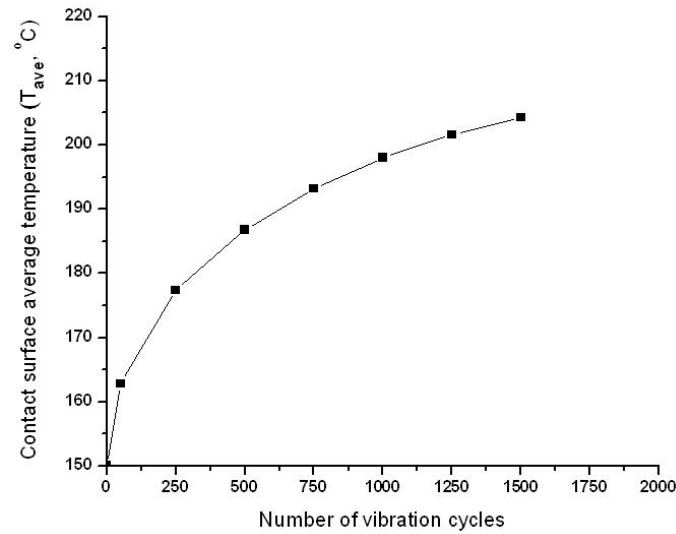


Fig. 3.2: Average temperature at the contact surface as a function of the vibration cycle.

and explained by superposition of vibration waves [56].

A similar plastic strain distribution can be seen along path A-C (the depth direction). In Figure 3.3c, a small region close to point A has a plastic strain level higher than 1 after 500 vibration cycles. The depth direction plastic strain in all other regions along path A-C is lower than 0.1. Since the depth direction plastic strain is the easiest to measure experimentally, it will be used for verification of the model later.

Path A-D is perpendicular to the sonotrode vibration; therefore, there is no waveshape deformation distribution (Figure 3.3d) as seen along path A-B. A steady increase of plastic strain levels can be seen as the number of cycles increases. The steady-state strain distribution indicates a peak level close to point A.

The evolution of the average von Mises plastic strain on the contact surface as a function of time (or vibration cycle) is shown in Figure 3.4. It has the same trend as the temperature vs. vibration cycle in Figure 3.2. The average plastic strain also approaches a steady state at the 1000th vibration cycle when the contact surface average temperature is close to the steady state. The time to reach the steady state for the average plastic strain is later than that for the three paths defined in Figure 3.3.

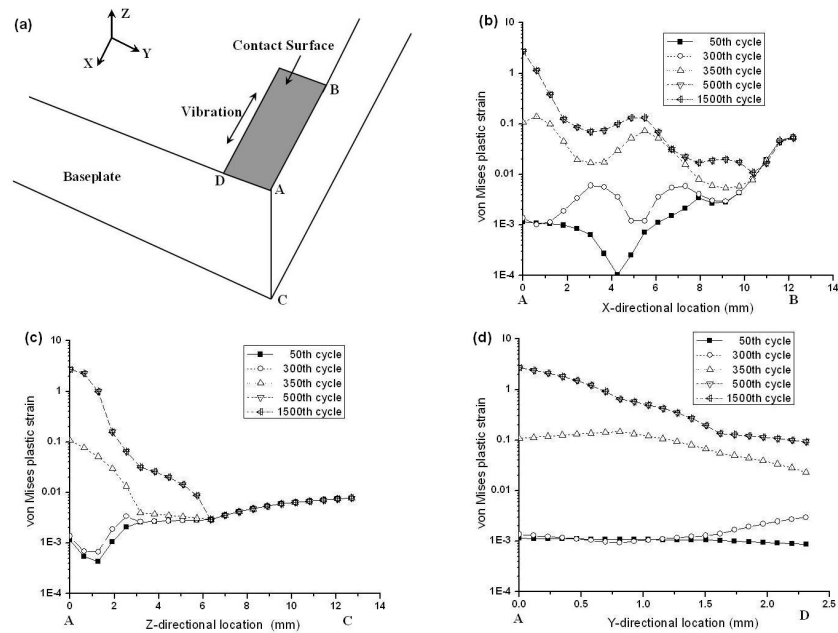


Fig. 3.3: Distribution of plastic strain as a function of location and vibration cycle.

In the initial period when the temperature is increasing, the yield strength of material in the contact area decreases. When the von Mises stress at any location reaches the yield condition at a certain specific temperature, the material yields and plastic deformation occurs. When the friction heat generation on the interface balances the heat loss and the average temperature approaches the steady state, the increase in plastic deformation stops and a similar steady state is also achieved. The number of cycles to reach this steady state may decrease (or the time may be shorter) when plastic deformation heat is considered in addition to the friction heat.

One important experimental observation in UW is that the strongest bonding is only achieved with an optimized travel velocity of the sonotrode. Too fast travel velocity does not provide enough dwell time (or number of vibration cycles) for bonding; while too slow a travel velocity does not result in good bonding either. The evolution of temperature and plastic strain obtained in this study can help quantitatively understand the phenomena, and provide control strategies for stronger bonding.

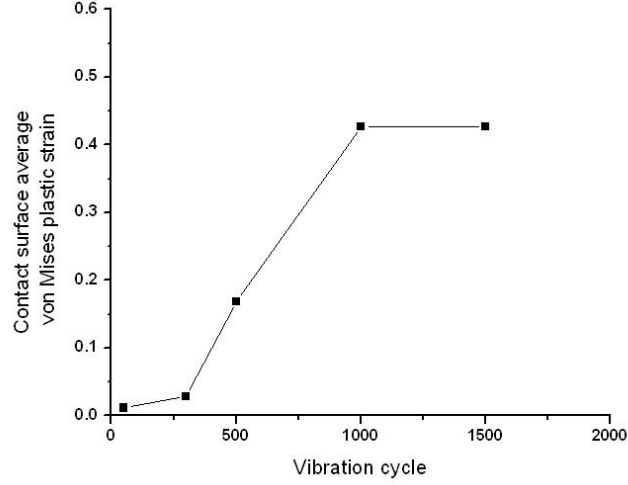


Fig. 3.4: Average plastic strain at the contact surface as a function of vibration cycle.

3.1.3 FEM Model Validation

Definition of Z-directional Compressive Equivalent Plastic Strain in Experiment

An experimental validation of the finite element results shown in Figure 3.3c was conducted. In UW, the height of deposit is smaller than the sum of the foil layers' original thickness, and the relative ratio changes with the number of layers. Figure 3.5 shows the difference in substrate height before bonding (a) and after bonding (b). Thus, the Z-directional compressive equivalent plastic strain can be defined as follows:

$$\epsilon_{z,eq}^{pl} = \frac{(H_a - H_b)}{H_b} \quad (3.1)$$

Samples of various numbers of layers are made on the UW machine with the same process parameters used in the simulations. The sample heights after UW bonding are measured with a precision height gauge with a 0.001-mm resolution. The substrate's height before bonding is calculated by 0.1 mm (foil thickness) \times 8 (layers number). Thus, the Z-directional compressive equivalent plastic strain from the experiment can be calculated by Equation 3.1.

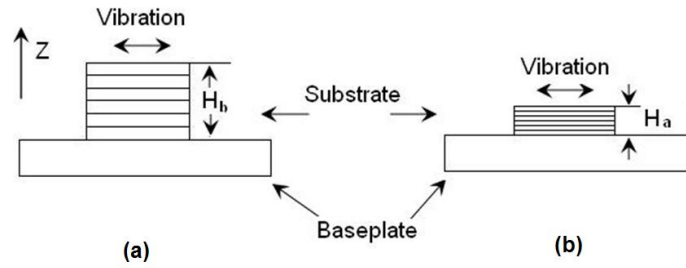


Fig. 3.5: Schematic of changing of substrate's height.

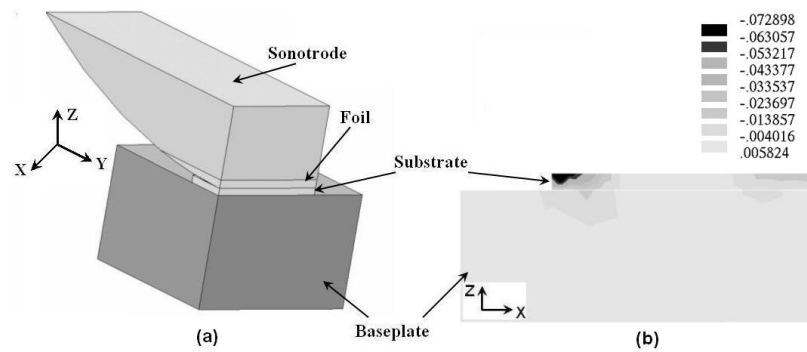


Fig. 3.6: 3-D coupled-field model for validation (a) and distribution of Z-directional plastic strain (b).

Definition of Z-directional Compressive Equivalent Plastic Strain in Simulation

For the purpose of finite element model validation, a new finite element model is developed to incorporate a rectangular substrate 12.2 mm long and 5.0 mm wide (Figure 3.6a). The same simulation strategy is applied to the new simulation with various buildup layers. A cut section plane parallel to the vibration direction and located at the center of the substrate is chosen for extracting the Z-directional plastic strain. Figure 3.6b shows the distribution of Z-directional plastic strain in the cut section for the eight layers of the substrate. The average value is calculated and defined as the Z-directional compressive equivalent plastic strain in simulation.

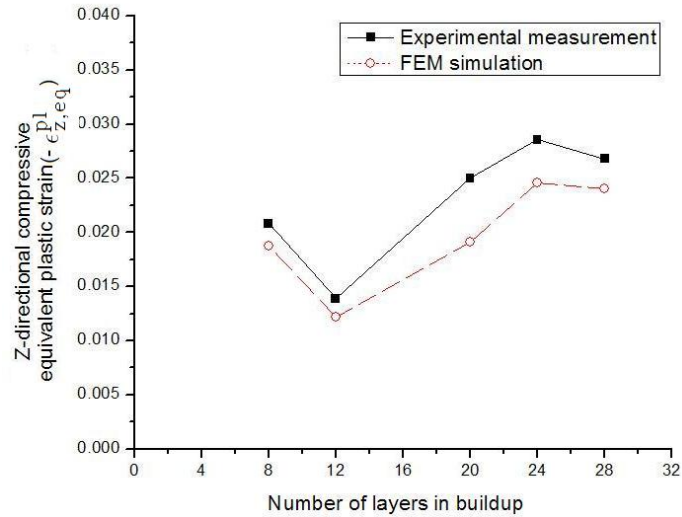


Fig. 3.7: Comparison of experimental measurements and simulation result for the average Z-direction compressive plastic strain.

Comparison

Figure 3.7 shows the comparison of Z-directional compressive equivalent plastic strain between experiment measurement and finite element simulation. It can be seen that the finite element model in this paper is able to correctly predict the trend and magnitude of the experimental data. The error may be very likely due to the omission of contribution of plastic heat flux to plastic strain. Another error source is the use of solid substrate in the validation model instead of the true substrate that actually has a layered structure and different mechanical properties. When the temperature increases, the tendency for microstructure change, such as recovery, which is not included in the model, may also contribute to the error.

3.2 Thermomechanical Simulation Results with Both Friction and Plastic Deformation Heat Generation

3.2.1 Distributions and Evolution of In-process Variables on the Contact Surface

Compressive Normal Stress

To help visualize the stress and strain distributions at the contact interface, paths A-B and C-D are defined in Figure 3.8, with path C-D along the edge of the contact area, and path A-B along the centerline underneath the sonotrode. The distributions and evolution of compressive normal stress (σ_z) have been extracted and shown in Figures 3.9 and 3.10 along paths A-B and C-D. The distributions of the stress are not uniform, but rather exhibit a wave feature. As the vibration cycle increases, the compressive normal stress fluctuates around 180MPa along path A-B, and 140MPa along path C-D. Under the clamping force, path A-B has the maximum Z-directional displacement, while path C-D has the minimum displacement. The σ_z for paths parallel to and in-between A-B and C-D is in the range between 180MPa and 140MPa.

The significance of σ_z and its variation is, that combined with the friction coefficient, it determines the friction stress on the contact surface based on Coulomb's law of friction. However, to determine the frictional heat, the shear stress and slide distance are required.

Slide Distance

Along the sonotrode vibration direction, the relative slide movement at the interface is the result of a superposition of sinusoidal vibration of the sonotrode and vibration of the baseplate. This relative motion of the contact interface for each vibration cycle can be defined as the slide distance $L_{slide} = L_{st} - L_{bp}$, where L_{st} is the sonotrode moving distance ($L_{st} = 32\mu\text{m}$ when the sonotrode makes a peak-to-peak movement with a $16\mu\text{m}$ amplitude), and L_{bp} is the moving distance of the baseplate during the same period. Depending on the magnitude and direction of the baseplate movement, L_{slide} can be greater or smaller than

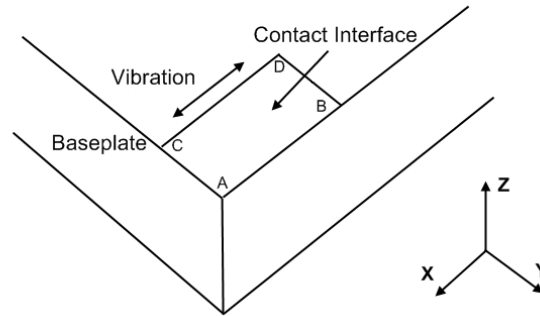


Fig. 3.8: Definition of paths A-B and C-D.

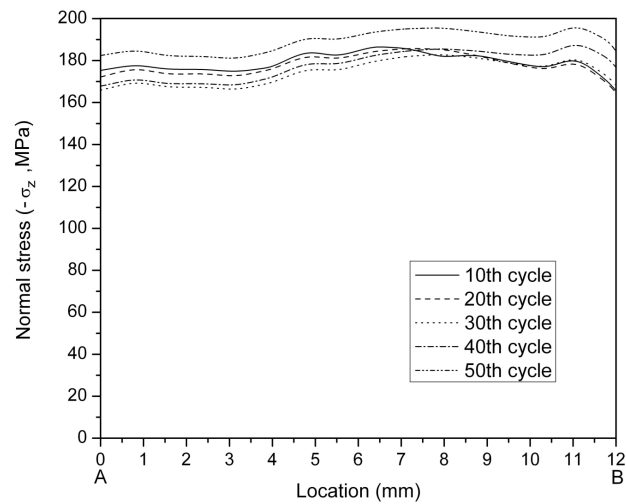


Fig. 3.9: Compressive normal stress distributions and evolution along path A-B.

$$L_{st} = 32\mu\text{m}.$$

The distributions and evolution of the relative slide distance along paths A-B and C-D on the contact surface are plotted in Figures 3.11 and 3.12. Initially, the slide distance is

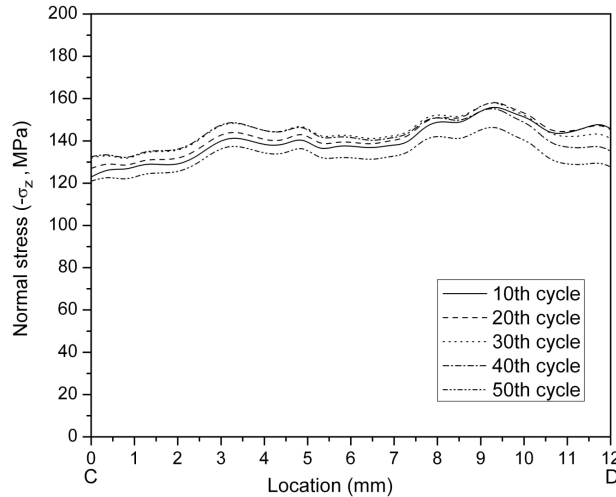


Fig. 3.10: Compressive normal stress distributions and evolution along path C-D.

smaller than the $32 \mu\text{m}$ sonotrode distance. As the number of vibration cycles increases, the slide distance increases. The region close to edge B-D of the contact area has a slide distance higher than $32 \mu\text{m}$. This means the baseplate is vibrating in the opposite direction against the sonotrode in that region. The same region also experiences the greatest increase in the slide distance as the number of cycles increases. As can be predicted, the B-D region will have the greatest frictional heat generation.

Shear Stress

The distributions of shear stress (τ_{xy}) at the contact interface show dramatic undulations. A shear stress peak exists at the center of the contact interface (near point A) on path A-B (Figure 3.13). A shear stress peak exists at the edge of the contact interface (near point D) on path C-D (Figure 3.14). The average shear stress on path C-D is greater than that on path A-B. On path C-D, it is also notable that the shear stress undergoes a greater increase as a function of vibration cycle. The undulations of shear stress at various locations will produce a similar distribution pattern for friction heat. That effect is clearly seen in the temperature distributions.

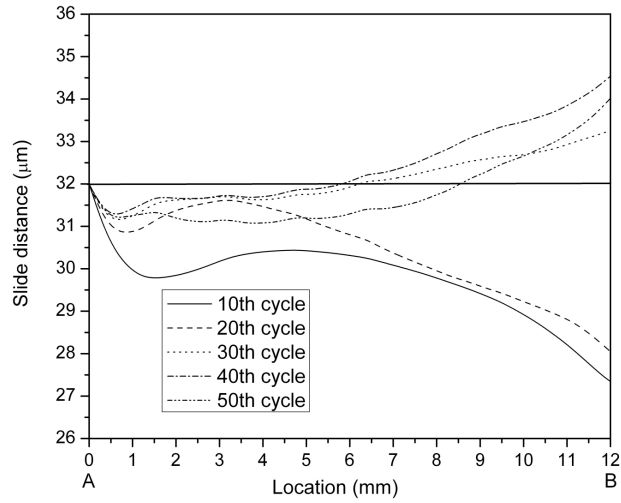


Fig. 3.11: Slide distance distributions and evolution along path A-B.

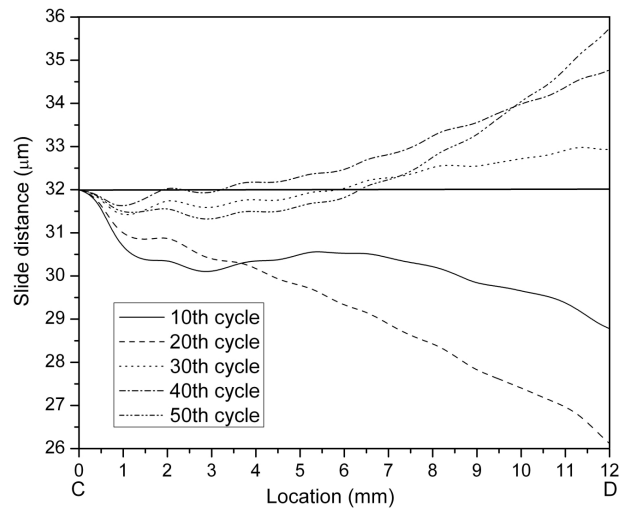


Fig. 3.12: Slide distance distributions and evolution along path C-D.

Temperature

Temperature distribution in the baseplate is shown in Figure 3.15. The peak temperature is located near the edges of the contact interface, and the temperature distribution in the depth direction is shallow, indicating the highly localized thermal phenomena in ul-

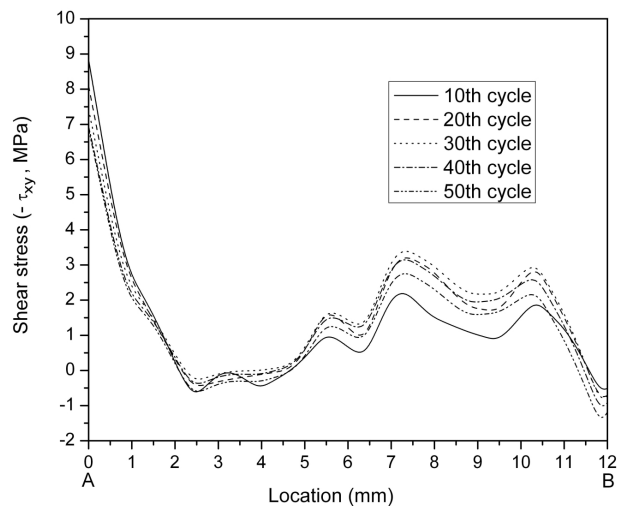


Fig. 3.13: Shear stress distributions and evolution along path A-B.

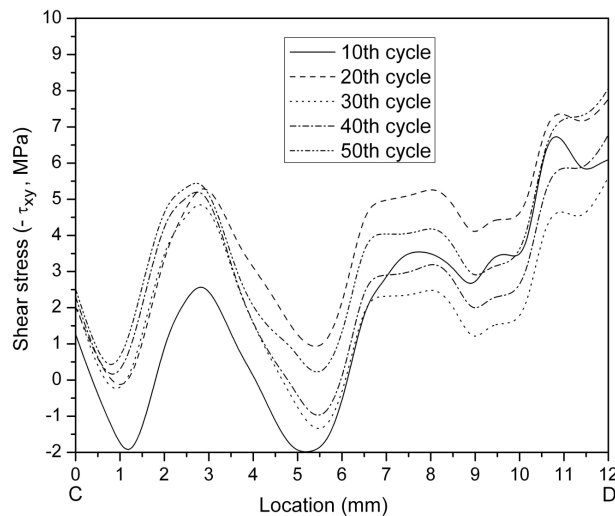


Fig. 3.14: Shear stress distributions and evolution along path C-D.

trasonic bonding. The evolution of the transient temperature field is shown in Figure 3.16 in the top-down view. The distributions and evolution of temperature along paths A-B and C-D have been plotted in Figures 3.17 and 3.18. The high temperature zone is limited to a small volume corresponding to the contact surface. This high temperature zone may

become larger with more heat generation on the contact surface with more vibration cycles. However, the growth of the temperature field will be limited because of the limitation of heat generation capacity and the loss of heat due to heat transfer. It can be clearly seen that the major heat source is located on the edges of C-D and B-D, and the maximum heat source is around point D. The undulations in temperature distribution along C-D correlate well with the undulations in shear stress shown earlier.

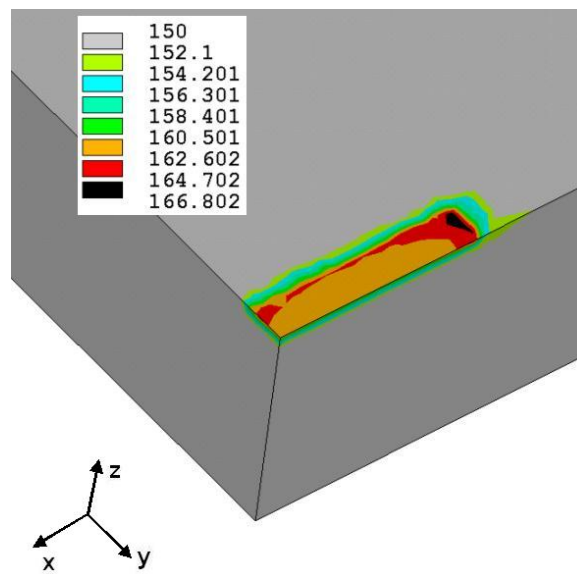
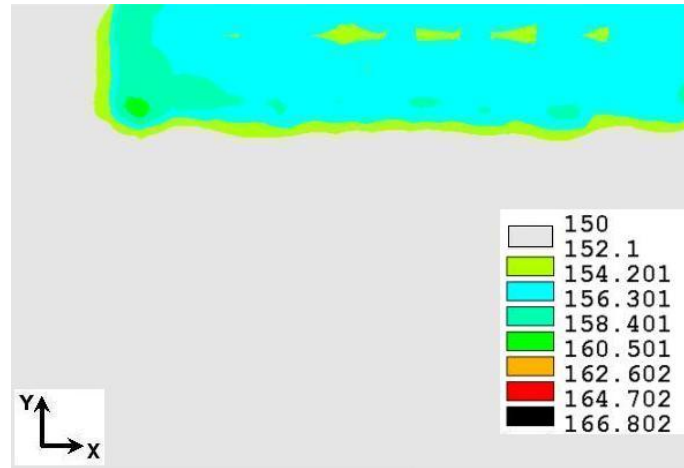
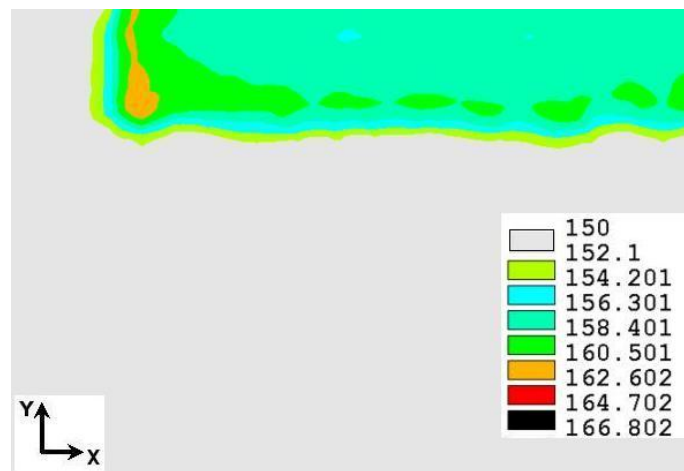


Fig. 3.15: The 3-D temperature distribution at the 50th vibration cycle (Unit: °C).

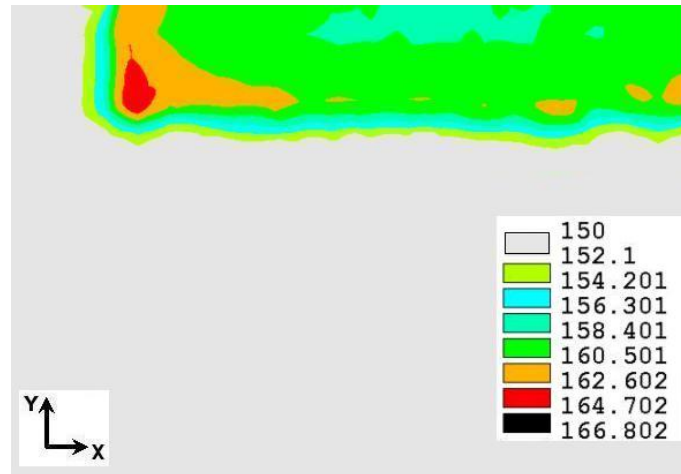


(a)

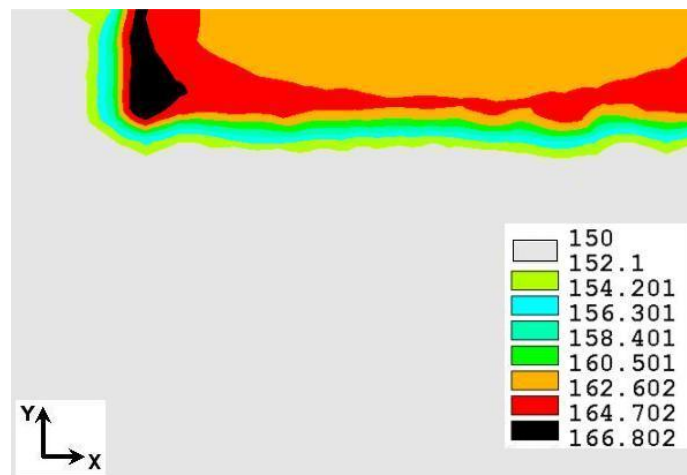


(b)

Fig. 3.16: Evolution of the top-down-view temperature distribution at : (a) the 10th vibration cycle, (b) the 20th vibration cycle, (c) the 30th vibration cycle, and (d) the 50th vibration cycle (Unit: $^{\circ}\text{C}$).



(c)



(d)

Fig. 3.16 continued.

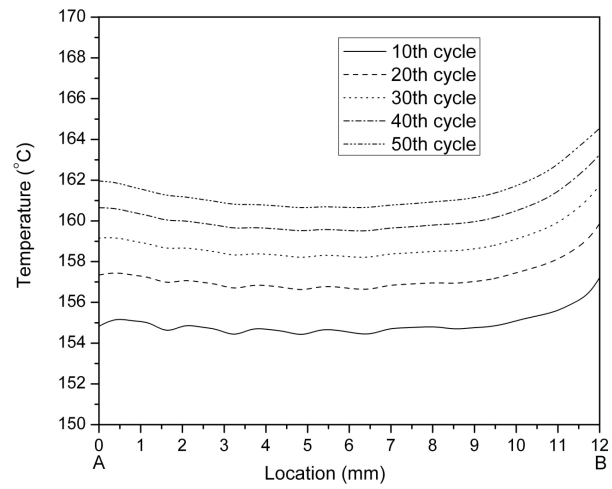


Fig. 3.17: Temperature distributions and evolution along path A-B.

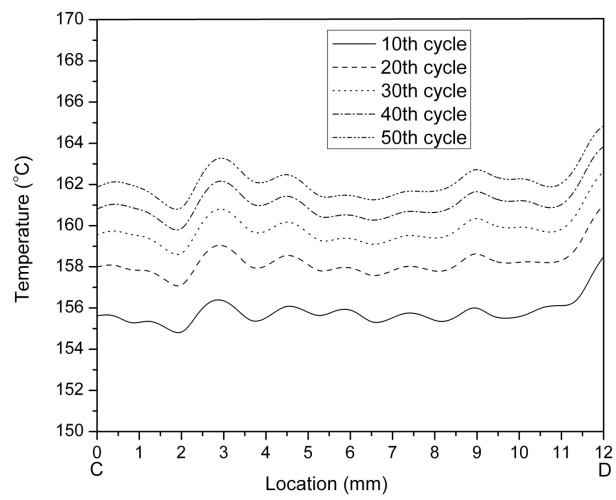


Fig. 3.18: Temperature distributions and evolution along path C-D.

Plastic Deformation

Figure 3.19 shows the 3-D von Mises plastic strain distribution in the baseplate at the 50th vibration cycle, and its evolution (top-down view) is shown in Figure 3.20. The distributions and evolution of von Mises plastic stain along the paths A-B and C-D have been plotted in Figures 3.21 and 3.22. The plastic deformation volume is also limited around the edges of the contact surface, where the conditions of high stresses and heat generation rate coexist. It is also found that the plastic deformation area extends toward the center of the contact surface with more vibration cycles. It can be predicted that the plastic deformation will cover the entire contact surface when more plastic deformation is generated with the sonotrode's vibration. A higher degree of plastic deformation has been found on path C-D than path A-B, and on edge B-D than edge C-D. After the 20th vibration cycle, the plastic deformation stops increasing, indicating a plastic saturation. A fluctuation in distribution of plastic deformation is also found to correlate with the shear stress distribution (Figure 3.22).

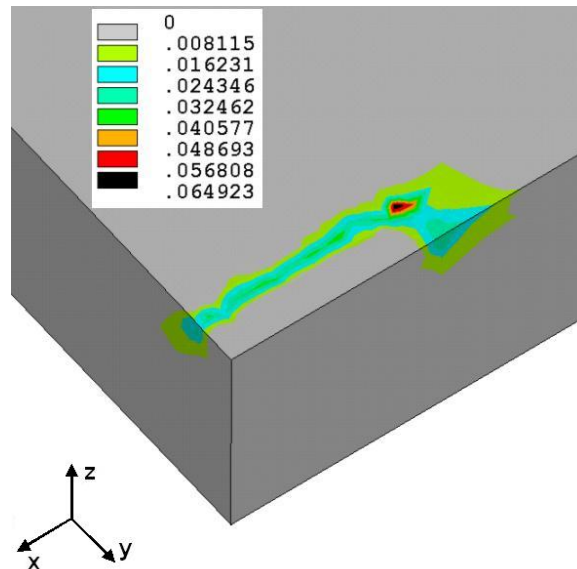
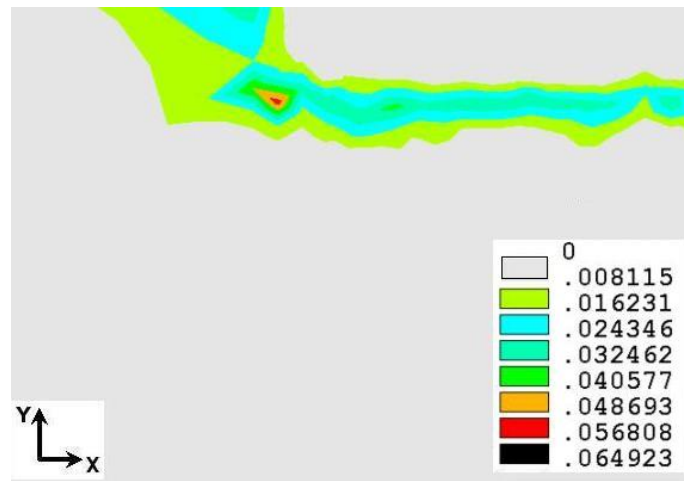
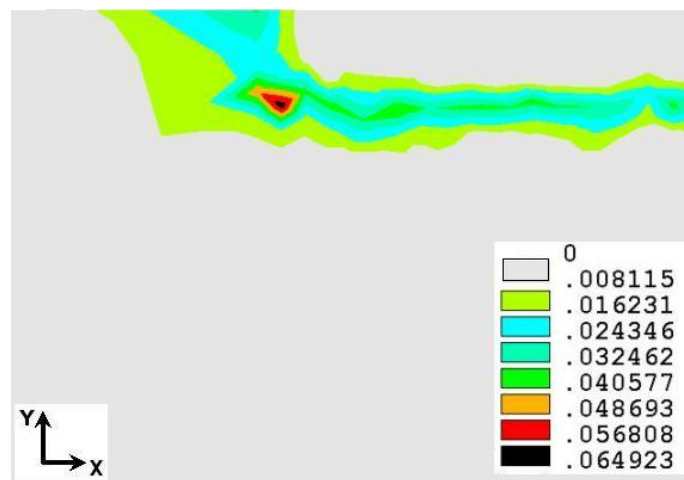


Fig. 3.19: 3-D von Mises plastic strain distribution at the 50th vibration cycle.

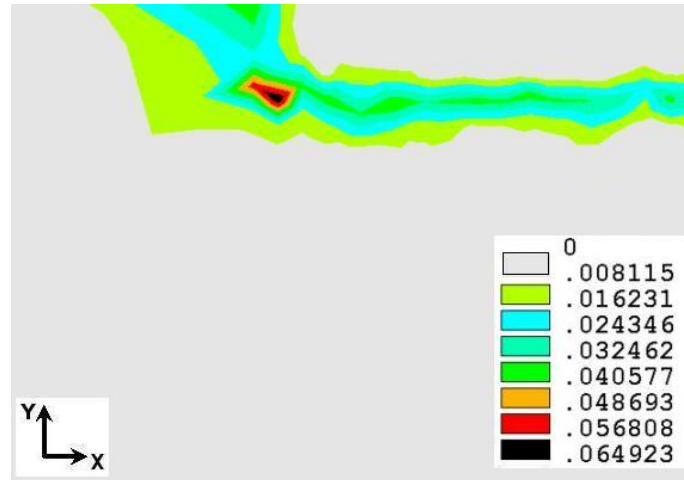


(a)

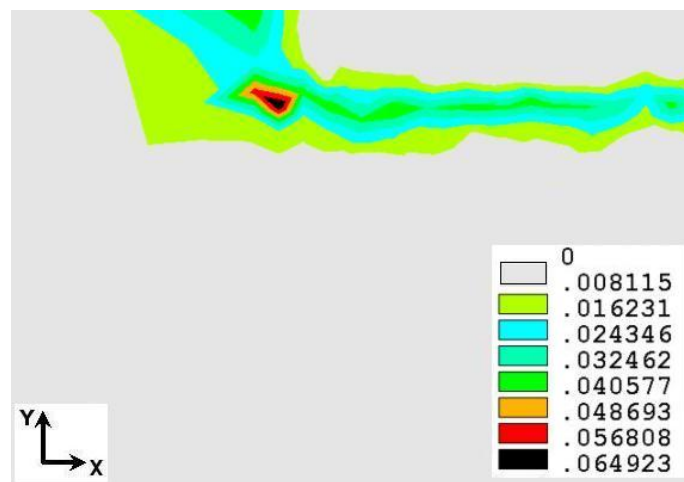


(b)

Fig. 3.20: Evolution of the top-down-view von Mises plastic strain distribution at: (a) the 1st vibration cycle, (b) the 3rd vibration cycle, (c) the 10th vibration cycle, and (d) the 50th vibration cycle.



(c)



(d)

Fig. 3.20 continued.

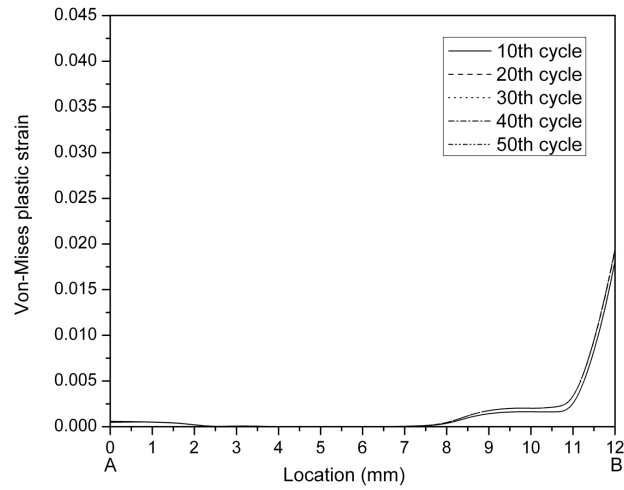


Fig. 3.21: Von Mises plastic strain distributions and evolution along path A-B.

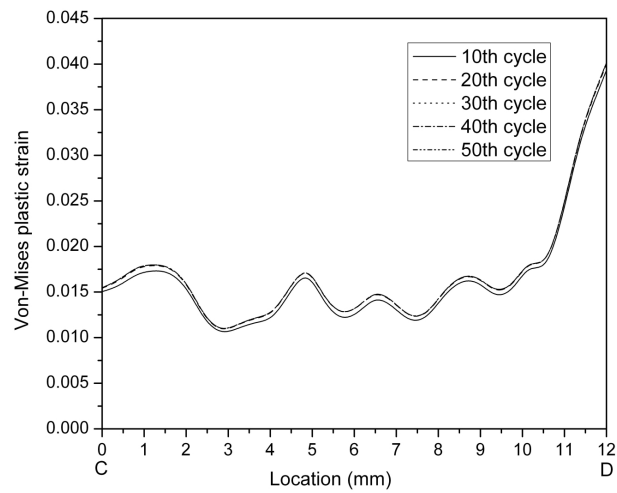


Fig. 3.22: Von Mises plastic strain distributions and evolution along path C-D.

3.2.2 Evolution of Average In-process Variables in the Bond Zone

To help understand the mechanism of ultrasonic bond formation and to study the interactions between the thermal and structure fields, the average in-process variables on the contact surface have been calculated. The surface heat density generation rate by friction is calculated by $\tau_{xy}L_{slide}/t_p$, where τ_{xy} is shear stress, and t_p is the period of vibration. The volume heat density generation rate by plastic deformation is calculated by $\Sigma\sigma_j \cdot \Delta\epsilon_j^p/t_p$, where σ_j is the principal stresses, $\Delta\epsilon_j^p$ is the change in the principal plastic strains, and j can be 1, 2 or 3. The equivalent volume heat density generation rate can be obtained by $C\rho\Delta T/t_p$, where C is specific heat, ρ is the density, and ΔT is the change in temperature. The net heat accumulation on the contact surface is the total heat generation subtracted by heat dissipation. Table 3.1 summarizes the average in-process variables in different periods. The first 50 vibration cycles of ultrasonic welding have been divided into three periods according to Figure 3.30.

(1) 1st - 10th vibration cycle: When the sonotrode starts to vibrate, a high stress state (Figures 3.23 and 3.24) is generated because of the high material stiffness at a lower temperature. Therefore, a large plastic deformation increase rate (Figure 3.30) with a high heat generation rate (Figures 3.26-3.28) is caused by the high stress state. In ultrasonic welding, two stress components, the normal stress σ_z and shear stress τ_{xy} , are more important than others. It is the normal stress σ_z that builds up a high stress state and yields the material, which contributes to plastic heating (Figure 3.23). The shear stress τ_{xy} mainly contributes to the generation of frictional heat. This observation can be verified by comparing the variations of frictional heat and that of shear stress τ_{xy} (Figures 3.24 and 3.26). Toward the end of this period, a large temperature increase rate (Figure 3.29) decreases the material properties rapidly. Consequently, the stresses and heat generation rates reduce fast (Figures 3.23, 3.24, and 3.26-3.28).

(2) 11th - 20th vibration cycle: This period is characterized by a lower heat generation rate than the previous period (Figures 3.26-3.28). The lower rate of temperature increase (Figure 3.29) slows down the rate of increase for the normal and shear stresses (Figures 3.23

and 3.24). Therefore, although the plastic deformation keeps increasing, its rate of increase is lower (Figure 3.30).

(3) 21st - 50th vibration cycle: The temperature continues to increase (Figure 3.29), indicating a positive heat accumulation at the bond interface - the heat generation on the contact surface is greater than the heat dissipation, although the equivalent heat generation rate decreases during this period (Figure 3.28). Therefore, the normal stress σ_z keeps decreasing (Figure 3.23) because of the decreased material properties. Toward the end of this period, the plastic deformation (Figure 3.30) has stopped increasing, arriving at a saturation state.

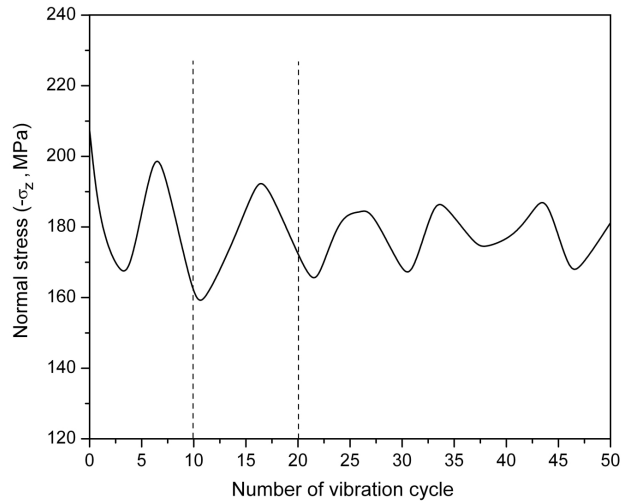


Fig. 3.23: Average compressive stress on the contact surface vs. vibration cycle.

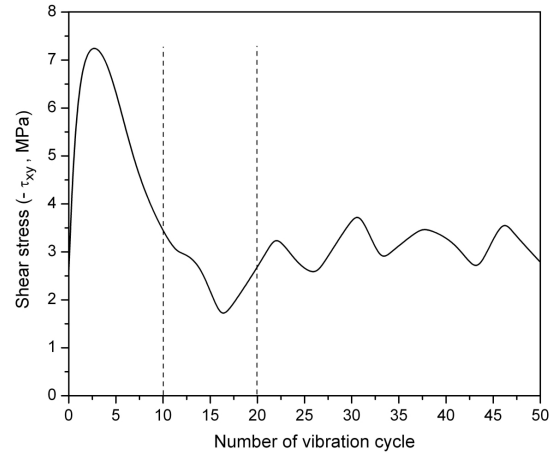


Fig. 3.24: Average shear stress on the contact surface vs. vibration cycle.

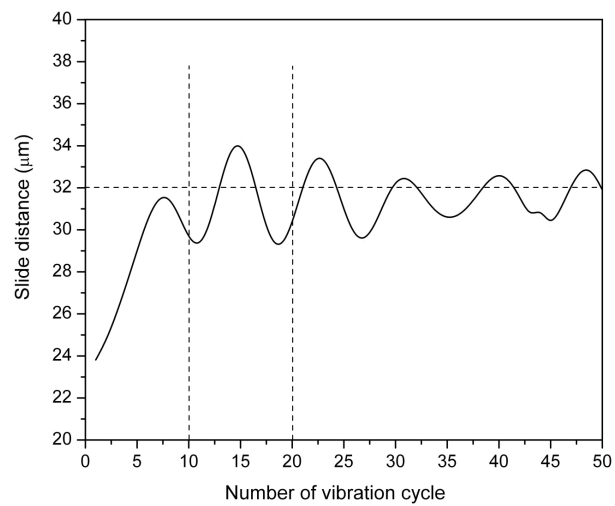


Fig. 3.25: Average slide distance on the contact surface vs. vibration cycle.

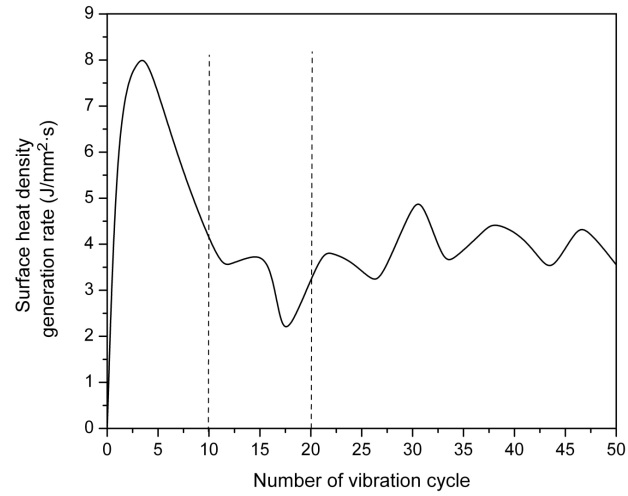


Fig. 3.26: Average heat generation rate by friction on the contact surface vs. vibration cycle.

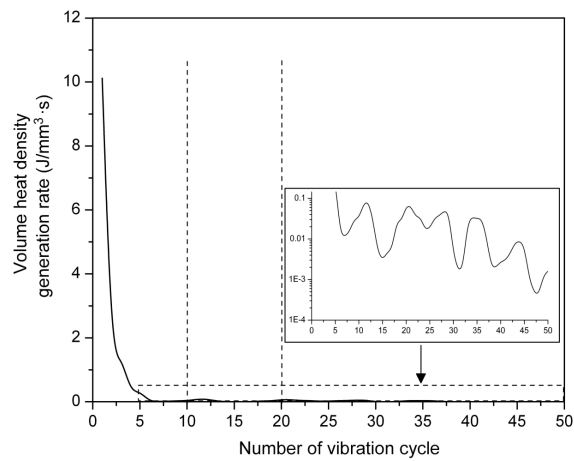


Fig. 3.27: Average heat generation rate by plastic deformation on the contact surface vs. vibration cycle.

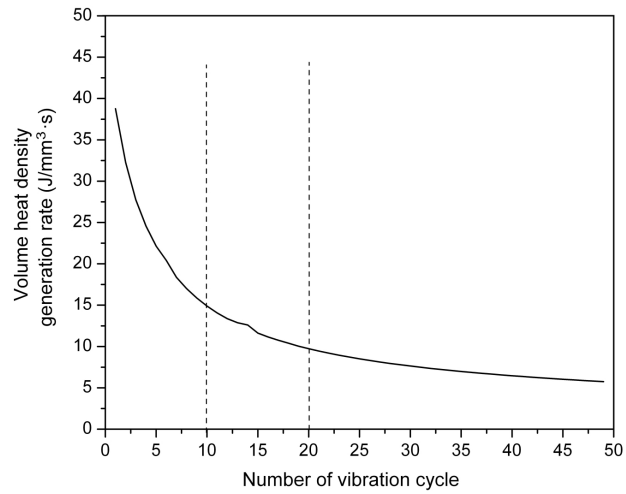


Fig. 3.28: Average equivalent heat generation rate on the contact surface vs. vibration cycle.

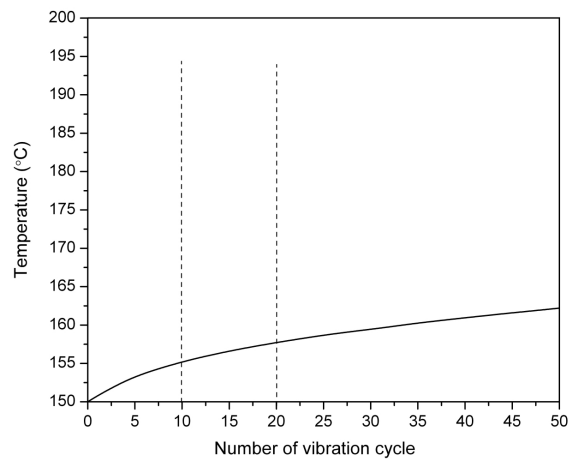


Fig. 3.29: Average temperature on the contact surface vs. vibration cycle.

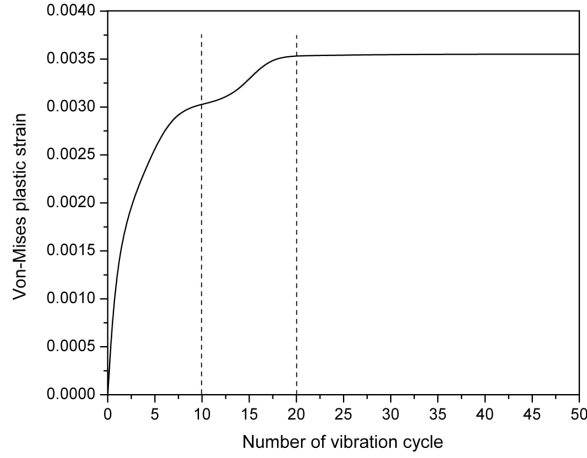


Fig. 3.30: Average von Mises plastic strain on the contact surface vs. vibration cycle.

Table 3.1: Evolution of Average In-process Variables on the Contact Surface

In-process variables	1st - 10th cycle	11th - 20th cycle	21st - 50th cycle
Normal stress $-\sigma_z$, (MPa)	182.63	180.41	176.87
Shear stress $-\tau_{xy}$, (MPa)	5.49	2.43	3.13
Slide distance L_{slide} , (μm)	28.50	31.90	31.60
Friction heat generation rate $\Delta H_f / \Delta t \cdot area$, ($J/mm^2 \cdot s$)	7.24	3.14	3.94
Plastic heat generation rate $\Delta H_p / \Delta t \cdot vol$, ($J/mm^3 \cdot s$)	1.37	0.03	0.02
Net heat accumulation rate $\Delta H_t / \Delta t \cdot vol$, ($J/mm^3 \cdot s$)	23.21	11.66	7.19
Temperature T, ($^{\circ}C$)	153.72	156.32	160.22
Changing rate of temperature $\Delta T / \Delta t$, ($^{\circ}C/s$)	1.20e4	4.80e3	3.00e3
Plastic deformation (ϵ_{von}^p)	2.32e-3	3.32e-3	3.55e-3
Changing rate of plastic deformation $\Delta \epsilon_{von}^p / \Delta t$, (1/s)	5.52	1.00	6.66e-3

3.2.3 Correlation of Simulation and Experimental Results

It is difficult to validate the numerical models of ultrasonic welding directly because of the difficulties of experimental measurement due to process characteristics of ultrasonic welding. The contact surface where the bonding occurs is invisible to observers, and measurement devices or sensors tend to damage the contact surface and interfere with the bonding process. Thermal mechanical processes that are highly localized (a few millimeters), transient (under 20KHz vibration), and require high spatial resolution (the amplitude of vibration is 16 μm), pose challenges for measurement instrumentation.

However, aluminum foils bonded under different process parameters have been tested and the typical fractured surfaces are shown in Figure 3.31. The distribution of areas that are not bonded correlates with the distributions of shear stress and temperature as predicted by the simulation. The bonded area can be used to evaluate the bond strength.

A rectangular region along path C-D is selected and divided into 16 areas with the same dimensions. The bonded areas for these divisions have been measured and plotted in Figure 3.32. von Mises plastic strain has been selected among the in-process variables from simulation results, because it is a variable that includes the accumulated effect of both thermal and mechanical processes at the interface. Based on the simulation results, average von Mises plastic strains on these same 16 divisions for the 50th vibration cycle have been extracted and plotted in Figure 3.32. The von Mises plastic strains has the same trend as the bonded area along the chosen stripes on path C-D. A strong linear correlation between them is also shown in Figure 3.33. Large von Mises plastic strain will result in more bonded areas. Therefore, it seems reasonable that von Mises plastic deformation from the simulation model can be quantitatively correlated with the bonded area (bond strength) from experiments.

Another positive correlation between the simulation and experimental observations is that our simulation results are able to provide new insights to, and explanations for, the relations of the in-process variables and their effects on bond formation.

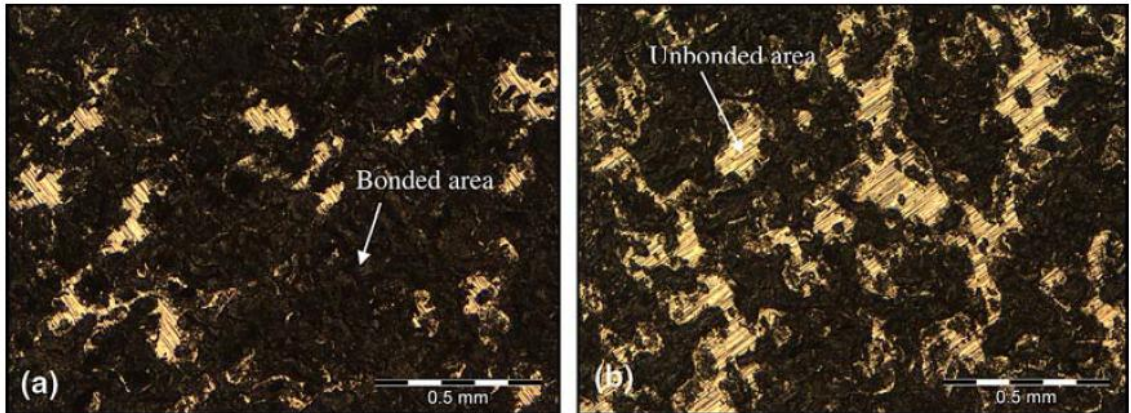


Fig. 3.31: Typical features of fractured surfaces: (a) 87.4% bonded, and (b) 76.4% bonded.

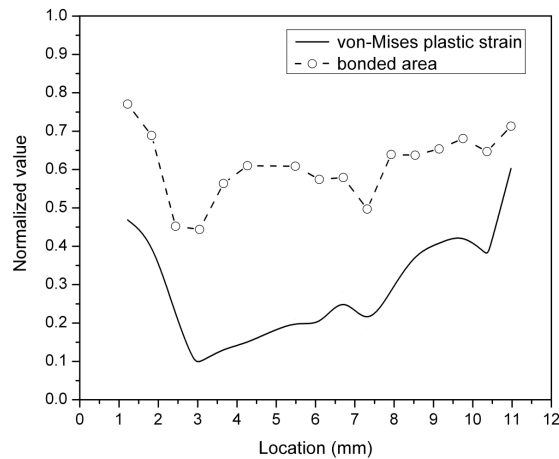


Fig. 3.32: Distributions of normalized bonded area and von Mises plastic strain along path C-D.

3.2.4 A Mechanism of Ultrasonic Welding

When the ultrasonic wave is transferred from the sonotrode to the contact surface of the baseplate, it travels along the X, Y, and Z directions. Each point on the contact interface can be seen as a wave source. The source and reflected waves interfere with each other, and a complicated superposition will be generated. The stress state that is determined

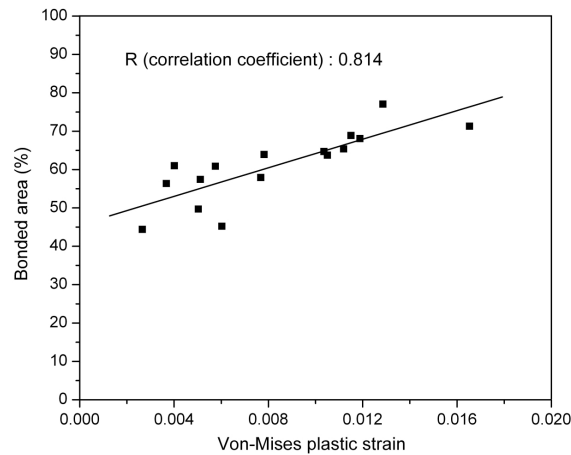


Fig. 3.33: Correlation of von Mises plastic strain and bonded area.

by the normal stress and friction conditions at the contact surface will alter the wave superposition. However, the characteristic wave-like distributions of in-process variables clearly are the result of wave superposition and vibration modes. A quantitative treatment of such dynamic effects is given elsewhere [57]. The fracture surface morphology on the ultrasonic bonding interface (Figure 5.8) is another indication of the ultrasonic vibration modes.

It appears that the temperature-dependent material stiffness (E) and yield stress (σ_y) have different effects on the in-process variables. The material stiffness will mostly affect the stress state and frictional heat generation, while the yield stress will mostly affect the generation of plastic strain and plastic heat. Based on the simulation results, a mechanism for bond formation in ultrasonic welding of an FCC metal (Al) is proposed as follows. (1) The normal clamping force and sonotrode vibration generate the shear stress (τ_{xy}) and shear strain. The shear strain displaces the oxide layer on the contact surfaces, to expose clean metal. (2) Shear stress and strain generate the frictional heat, and the plastic strain generates plastic heat. (3) A temperature rise will result from the equivalent heat generation determined by the combined effect of the frictional and plastic heat fluxes, and heat dissipation; (4) The rising temperature increases the plastic strain at the interface by lowering

the material mechanical properties, and changing friction conditions. This iterative process continues until the plastic deformation achieves saturation, at which point the freshly exposed metal surfaces achieve a metallurgical bonding because of close atomic proximity. Based on this bonding mechanism, plastic deformation affects the bond formation in the following ways: (1) it will disperse the surface oxide layer; (2) it will generate heat to make the atoms more thermally active; (3) it will deform the crystals and grains on the interfaces and generate a high dislocation density, which will also provide potential energy for grain boundary migration.

Chapter 4

Effect of Substrate Dimensions on Ultrasonic Bonding

A previous study shows that there is a maximum height-to-width ratio of free standing structures that can be built using ultrasonic consolidation. If the height-to-width ratio is higher than 1:1, no ultrasonic bonding can be achieved [58]. The objective of this work is to analyze the UW process by studying the bonding interface and the substrate, and to show how the ultrasound waves traverse in the substrate and its effect on the interface. Another objective is to develop the mechanics for ultrasonic consolidation that can explain the experimental results and point to the direction for ways for new development. An FEM model is set-up and the distributions of strain/stress in the substrate, and friction stress/displacement on the interface are analyzed. The FE simulation model has been validated by an ultrasound wave model.

4.1 A Dynamic Model for Ultrasonic Consolidation

Previous research on ultrasonic bonding has focused on the interface area where the bonding occurs [23, 59]. However, we found that wave interference within the substrate has a significant effect on joint formation. Therefore, a dynamic FEM model is developed to include the sonotrode and the entire substrate for a fundamental study of the time-dependent mechanical behavior of the substrate in ultrasonic consolidation.

Figure 4.1 shows the 2-D dynamic model involving the sonotrode that is pressed to contact with and vibrates against the substrate. The model domain is an analytical plane that is perpendicular to the welding direction and through the center of the sonotrode. The chosen model plane satisfies the plane strain condition. The base plate is not included in the model and its influence is simulated by adding a fixed boundary condition at the bottom surface of the substrate. The titanium sonotrode surface has been roughened so

that it is assumed that there is no slipping between the top surface of the aluminum foil and the sonotrode during vibration. A very large friction coefficient is assumed between the top aluminum foil and the titanium sonotrode to simulate the no slipping condition for the sonotrode. Any relative sliding happens between the bottom surface of the top foil and the substrate. The substrate is assumed to be homogeneous and isotropic. The Coulombs friction law, that friction is proportional to normal pressure and independent of slide velocity, is employed in simulation to determine the contact status (sticking or sliding) and shear stress at the contact interface. It should also be pointed out that the heating effect is not considered in this model.

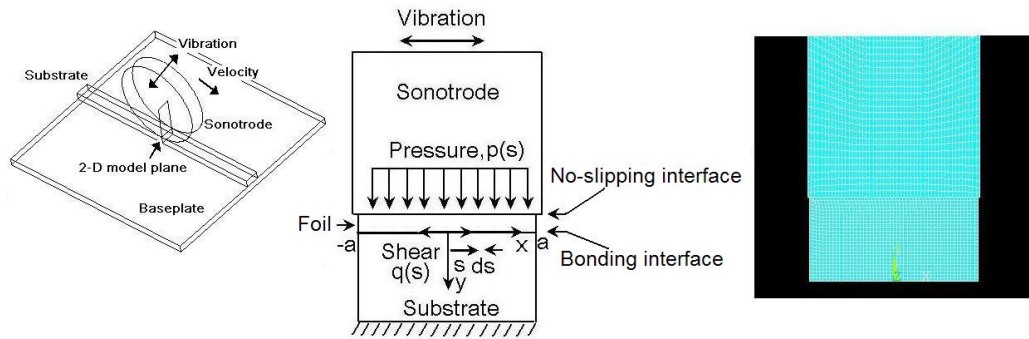
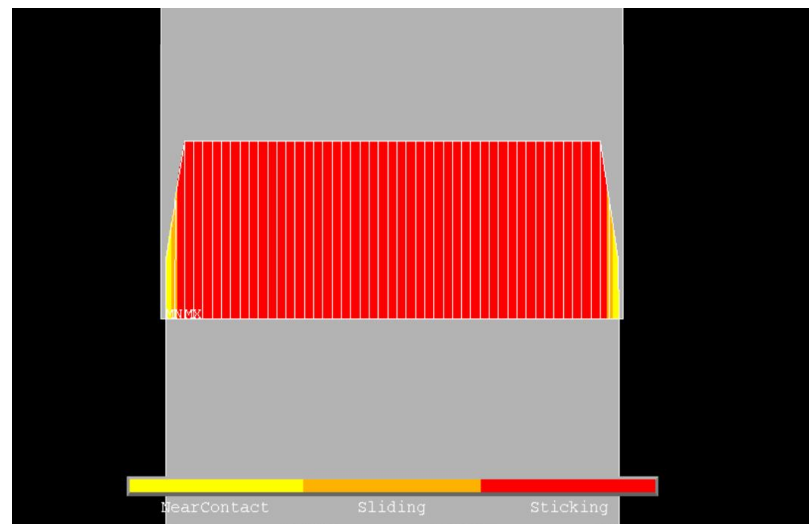


Fig. 4.1: 2-D dynamic model for ultrasonic consolidation.

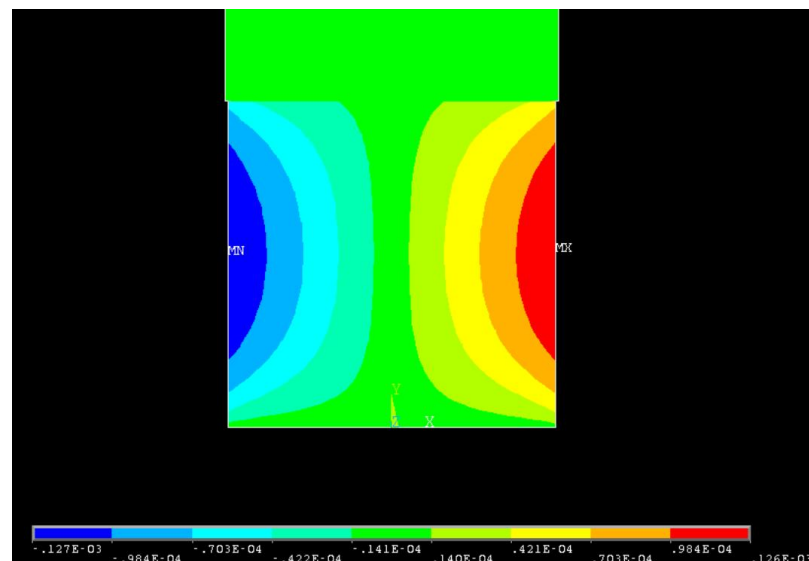
4.2 Static Analysis

Upon the initial application of the pressure on the substrate, and before the first vibration cycle is applied, the substrate exhibits a static distribution of stress and strain. For the substrate with an 1.0 height-to-width ratio, the distribution of mechanical status is shown in Figure 4.2. The friction status at the contact surface is mostly “sticking,” i.e., there is no relative motion in the horizontal direction between the sonotrode and the substrate. Near the two edges, however, there are small areas where “sliding” occurs. Right at the edges, the status of the interface is “near contact.” In Figure 4.2b, the horizontal displacement of the substrate shows a typical barreled shape: the mid-height of the substrate flows out undisturbed, while near the contact interface the frictional force opposes the outward flow

of the material. This effect is also seen in Figure 4.2c, in which the horizontal displacement at the contact interface reaches its maximum at the boundary between the “sticking” and “sliding” areas and decreases to zero at the center and the edges of the contact interface.



(a)



(b)

Fig. 4.2: Static distribution of mechanical status for substrate with a height-to-width ratio of 1.0: (a) contact interface friction status, (b) substrate’s horizontal displacement (inch), (c) contact interface horizontal displacement (inch), and (d) contact interface frictional stress (psi).

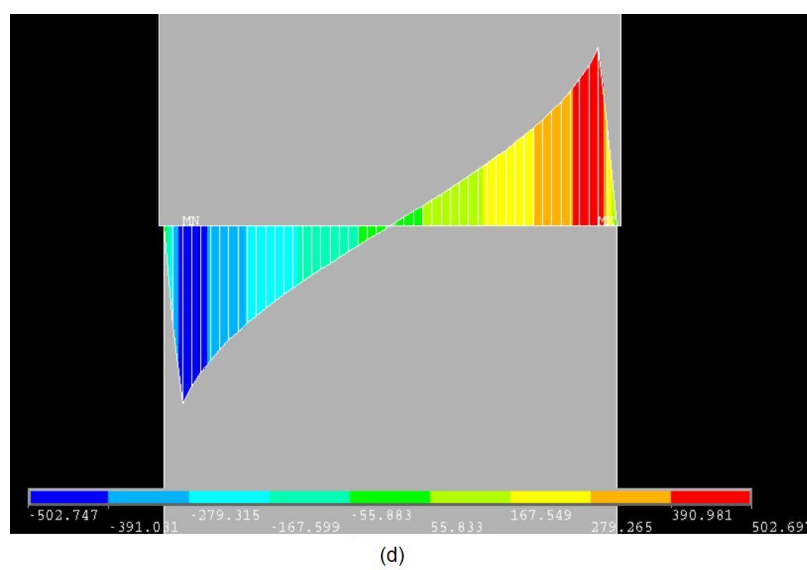
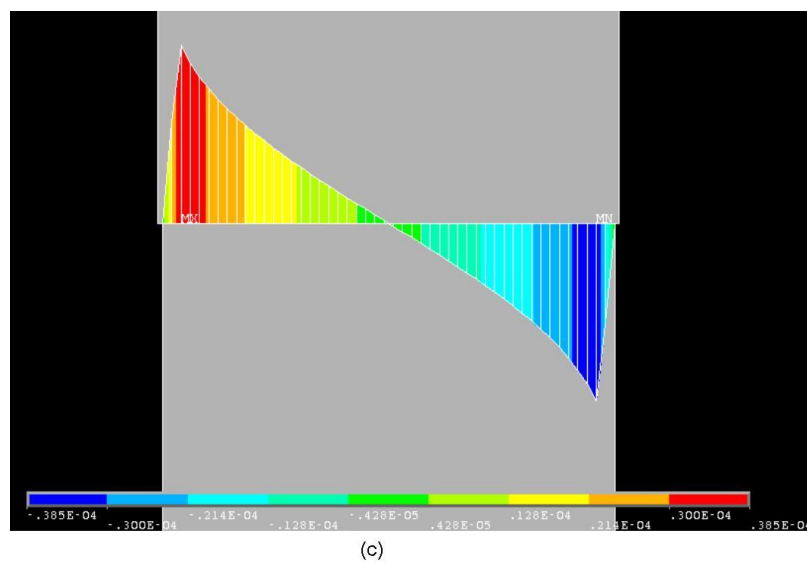


Fig. 4.2 continued.

Figure 4.2d shows the frictional stress distribution at the contact interface, respectively. The frictional stress points to the opposite directions for the two sides of the substrate, with a gradual decrease of intensity to the center of the interface, where the zero-friction point is located. A common feature for these stress distributions is that the maximum values are located at the boundary between the “sticking” and “sliding” areas. For the simulation conditions, there is no plastic deformation in the substrate under the initial static loading. Substrates with different height-to-width ratios have been simulated and the results are summarized in Figure 4.3. The frictional stress at the contact interface is shown in Figure 4.3a. It can be clearly seen that as the height-to-width ratio (H/W) is increased, the peak frictional stress decreases. The horizontal displacement at the interface has an inverse trends relative to the frictional stress. In Figure 4.3b, the substrate with the lowest H/W ratio has the lowest interface displacement. As the H/W ratio is increased, the interface displacement increases. More interface displacement means the contact interface is more “sliding” than “sticking.” The way in which “sticking” and “sliding” areas change as a function of H/W ratio is shown in Table 4.1. When other conditions remain the same, a higher substrate has a greater “sliding” area at the contact interface.

In Figure 4.3, all curves share a common trend: the maximum (or minimum) values in the curve move inwards towards the center of the contact interface as H/W ratio increases. Since the maximum values occur at the location where the “sticking”-to-“sliding” boundary is, this feature means that as the H/W ratio increases, more portions of the contact interface have become “sliding.”

Table 4.1: The Proportion of the Sticking Area at Contact Surface as a Function of the H/W Ratio

H/W ratio	0.25	0.5	0.75	1.0	1.5	2.0
$\frac{A_{sticking}}{A_{total}}$ (%)	100.0	93.5	95.9	95.5	91.4	87.0

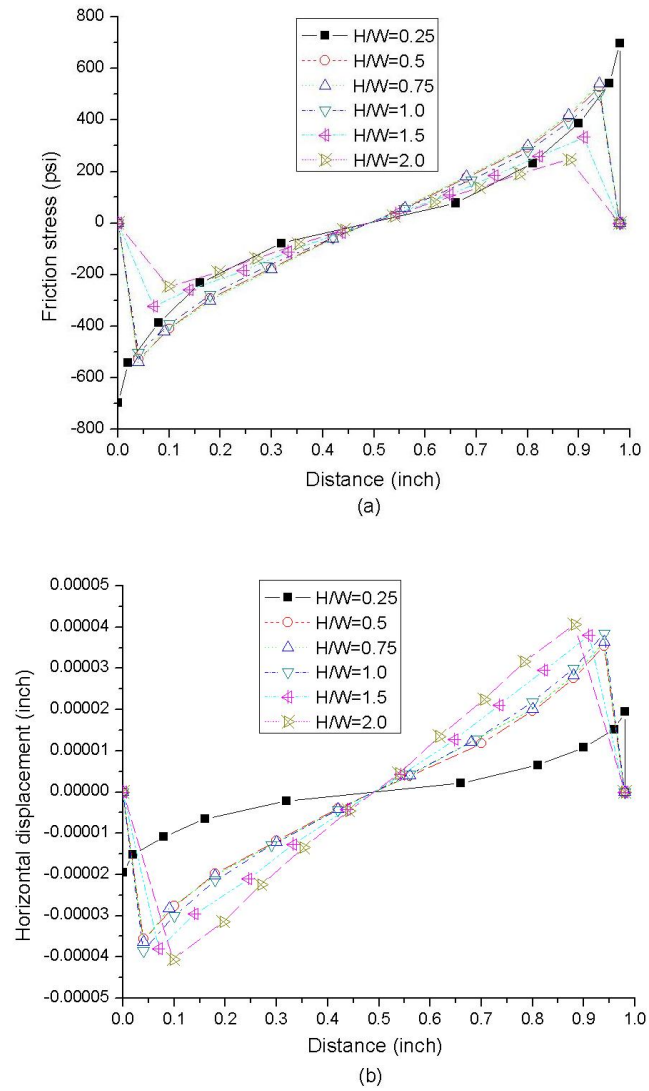


Fig. 4.3: Static distributions of stresses and displacement as a function of the H/W ratio: (a) friction stress (psi), and (b) horizontal displacement (inch).

4.3 Dynamic Analysis

When the sonotrode moves in opposite directions, the distributions of the contact interface displacement retain the same features as the static case. However, the zero-value point vibrates with the sonotrode (Figure 4.4). The displacement value at each point of contact exhibits a fluctuating variation with the vibration of the sonotrode. As the sonotrode moves to the left, the zero-value point also moves to the left.

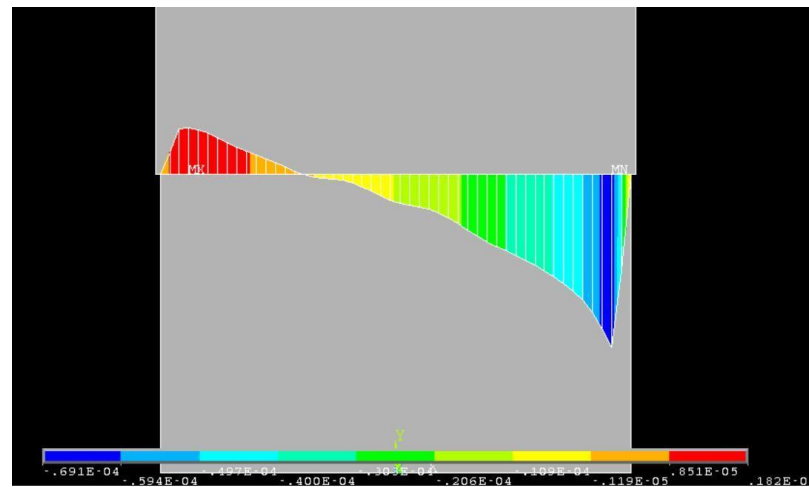
The magnitude of fluctuation in the displacement levels for a given point at the contact interface is defined as the “amplitude.” The distribution of interface displacement amplitude with different substrate H/W ratios and with different cycles is shown in Figure 4.5. With the H/W ratio increasing, the displacement amplitude increases (Figure 4.5a). This result agrees with the static analysis discussed earlier. The maximum variation of displacement occurs at the edge areas of the substrate, and the central areas have less variation.

Figure 4.5b shows the amplitude of displacement as a function of number of vibration cycles for a substrate with an H/W ratio of 0.75. For the same substrate height, as the number of cycles increases the displacement increases initially, but stabilizes to a “saturated” distribution after about 250 cycles. This result indicates the sonotrode traverse velocity (i.e., the welding speed) should be adjusted to ensure enough vibration cycles for the contact interface, but too slow a velocity will not increase the displacement action beyond the saturated level, in addition to decreasing the efficiency.

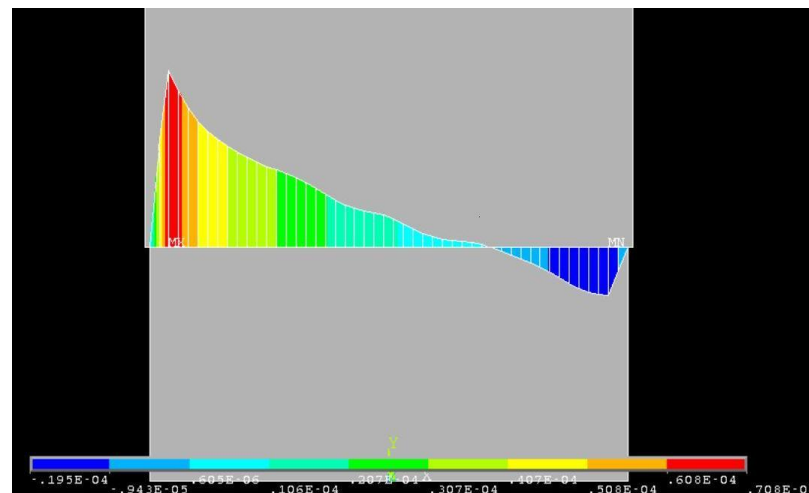
Away from the contact interface, the strain distribution within the substrate will influence the bond formation. Figure 4.6 shows the shear strain distributions for the substrate with a H/W ratio of 0.5 at different vibration cycles. There are four strain concentration regions that are located at four corners of the substrate. With increasing number of cycles, the originally concentrated strain regions start to spread and branch toward the inner regions of the substrate. After about 250 cycles, the shear strain distribution approaches a steady state. This confirms the earlier observed contact friction and displacement saturation at the 250th cycles for this specific experimental setup. When the sonotrode moves in the opposite direction, the pattern of the strain distribution does not change, except in an

inverse image (compare Figure 4.6e with f).

Figures 4.7 and 4.8 show the shear strain distributions at the 750th vibration cycle in substrates and the interface average shear strain respectively with various H/W ratios. The shear strain measures the degree of elastic/plastic deformation, which measures the potential for bond formation. For build H/W ratios less than 1.0, greater levels of shear strain ($\geq 2.4 \times 10^{-4}$) exist near the bond interface. In addition, significant shear strain exists inside the builds. Such internal shear strain appears to distribute in horizontal bands, apparently due to the interference of the traveling vibration waves. For build H/W ratios higher than 1.0, the shear strain has a lower level ($\leq 2.4 \times 10^{-4}$) at the bond interface. At these build heights, the internal shear strain bands are much less effective in producing bonding.

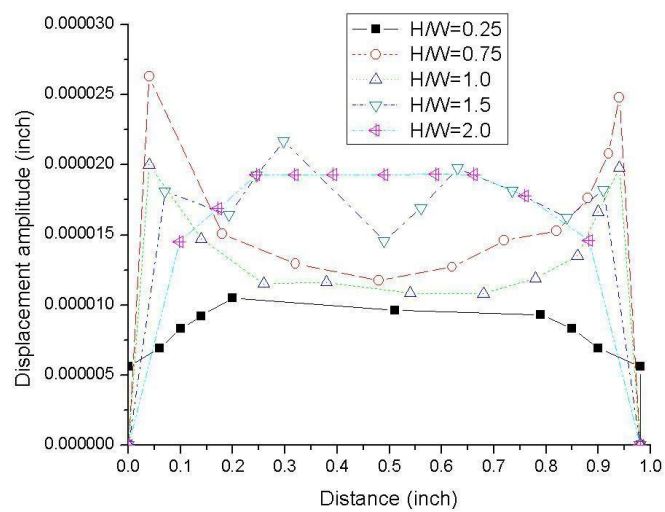


(a)

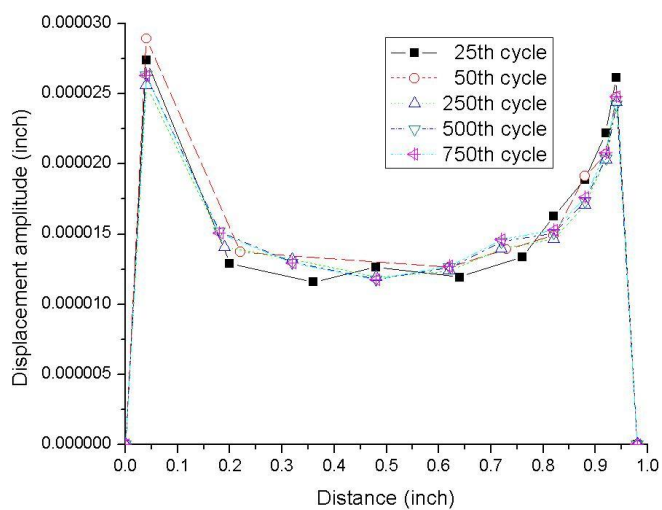


(b)

Fig. 4.4: Distribution of contact displacement (inch) at the 750th cycle with a H/W ratio of 0.75 when the sonotrode moves to opposite directions: (a) left direction displacement, and (b) right direction displacement.



(a)



(b)

Fig. 4.5: Displacement amplitude (inch): (a) at the 750th cycle for substrates with different height-to-width ratios, and (b) for substrate with a height-to-width ratio of 0.75 at different cycles.

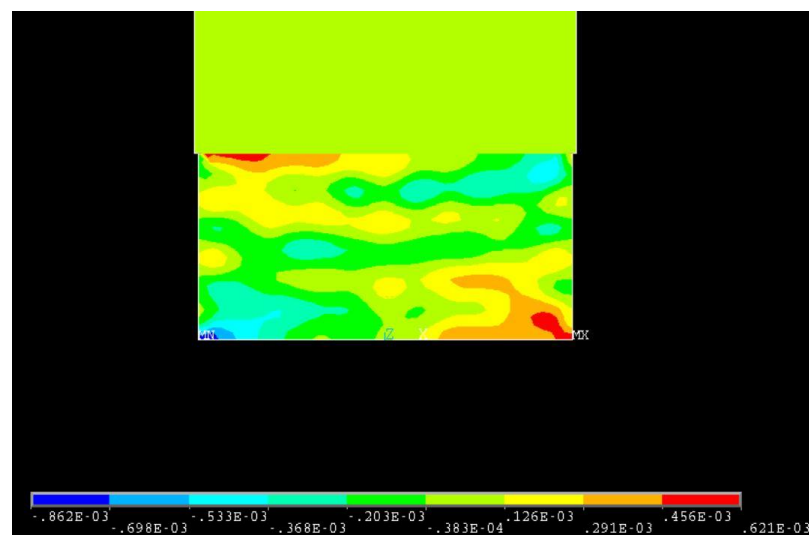
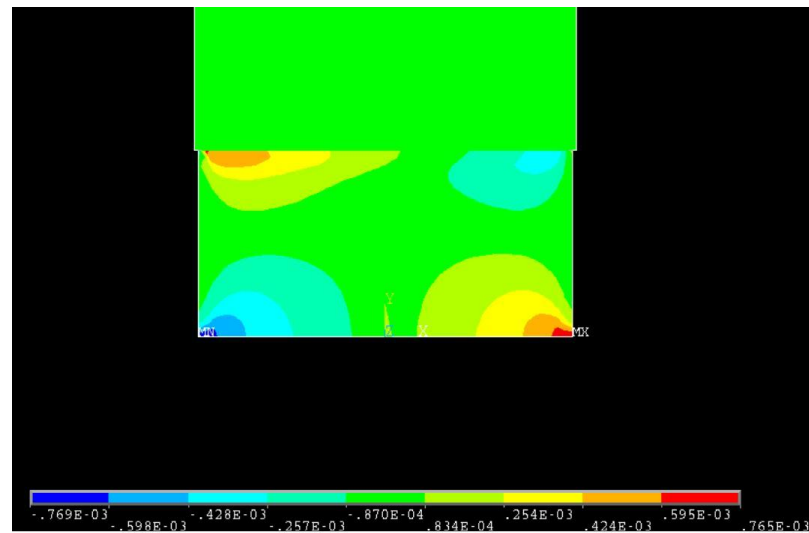


Fig. 4.6: Distribution of shear strain for the substrate with a 0.5 height-to-width ratio: (a) 1st cycle (right), (b) 25th cycle (right), (c) 50th cycle (right), (d) 250th cycle (right), (e) 750th cycle (right), and (f) 750th cycle (left).

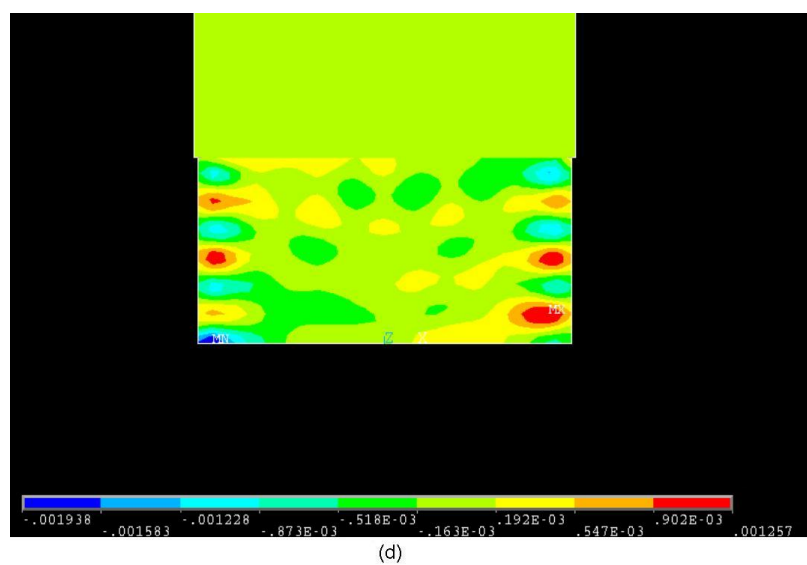
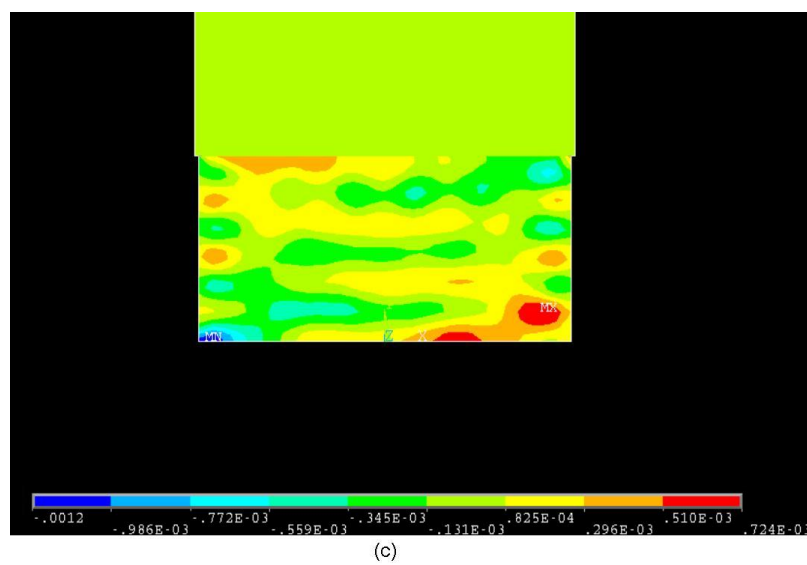


Fig. 4.6 continued.

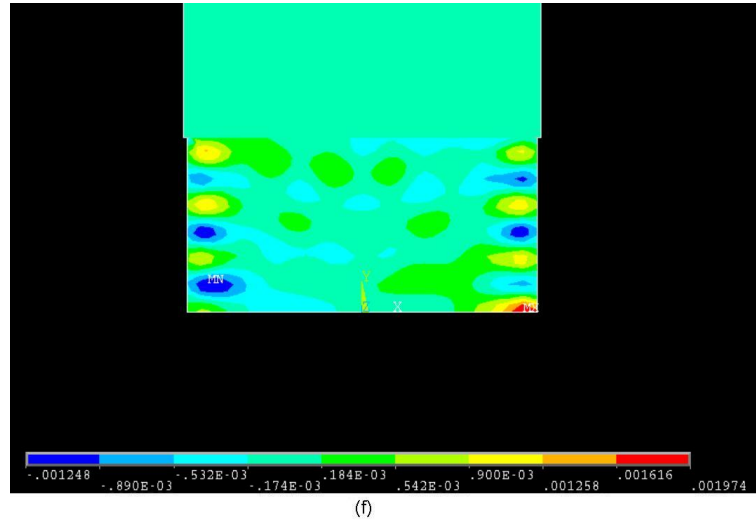
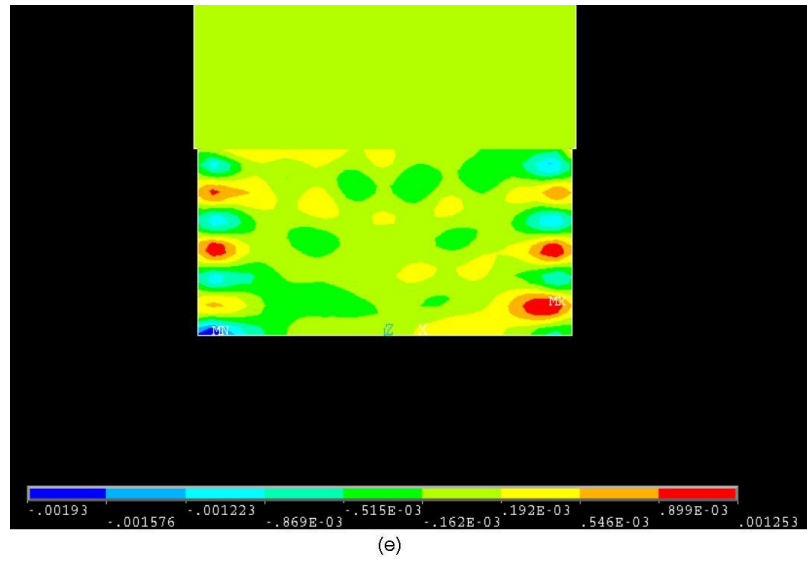
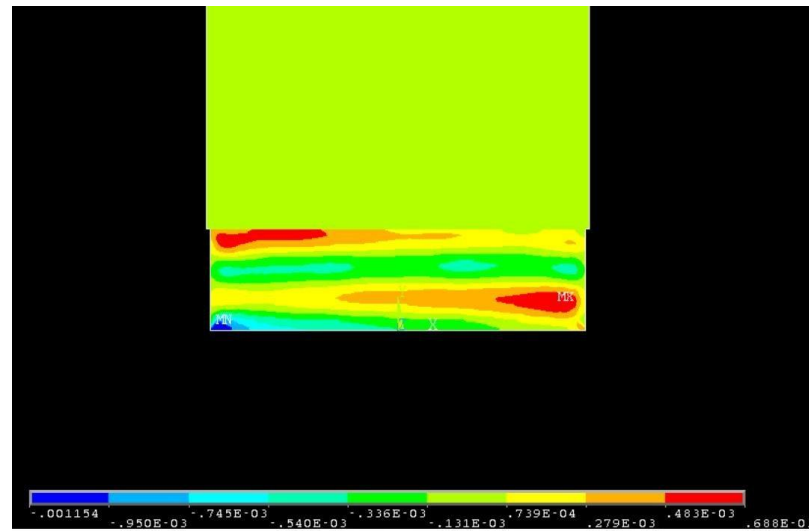
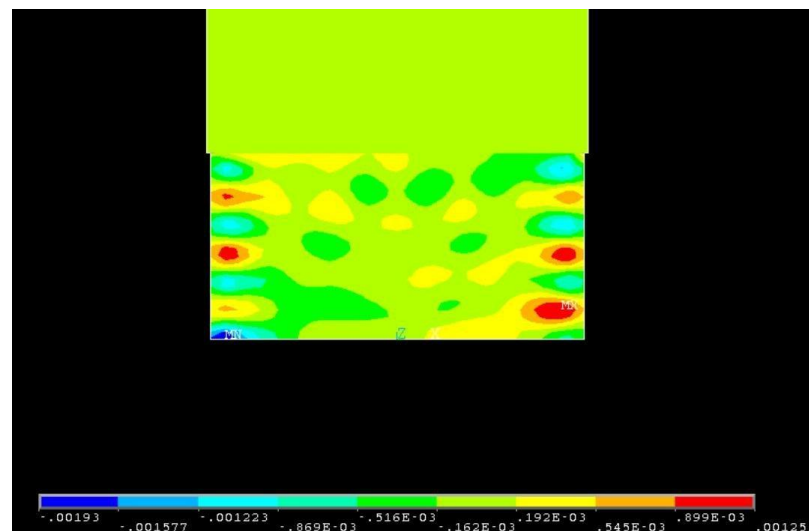


Fig. 4.6 continued.

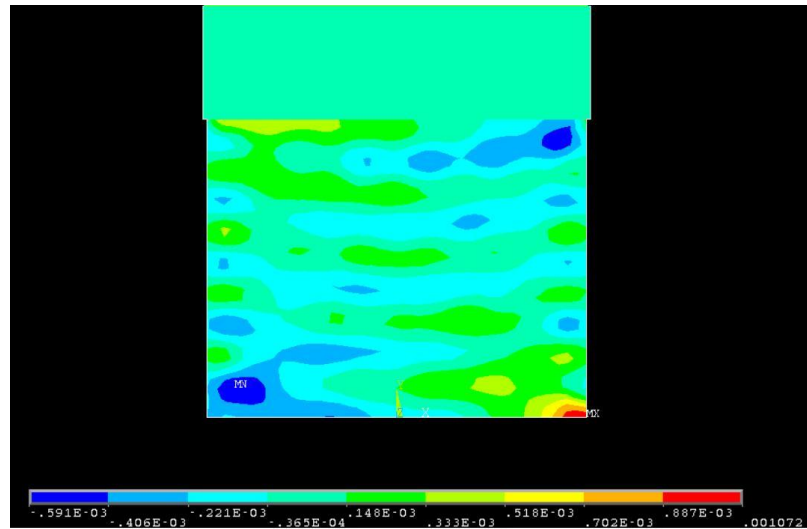


(a)

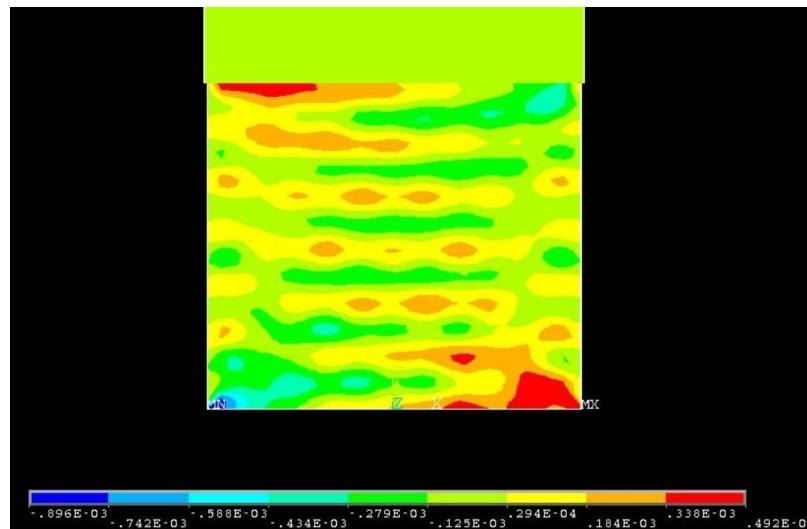


(b)

Fig. 4.7: Distribution of shear strain at the 750th cycle for substrate with different height-to-width ratios (H/W): (a) H/W=0.25, (b) H/W=0.5, (c) H/W=0.75, (d) H/W=1.0, (e) H/W=1.5, and (f) H/W=2.0.

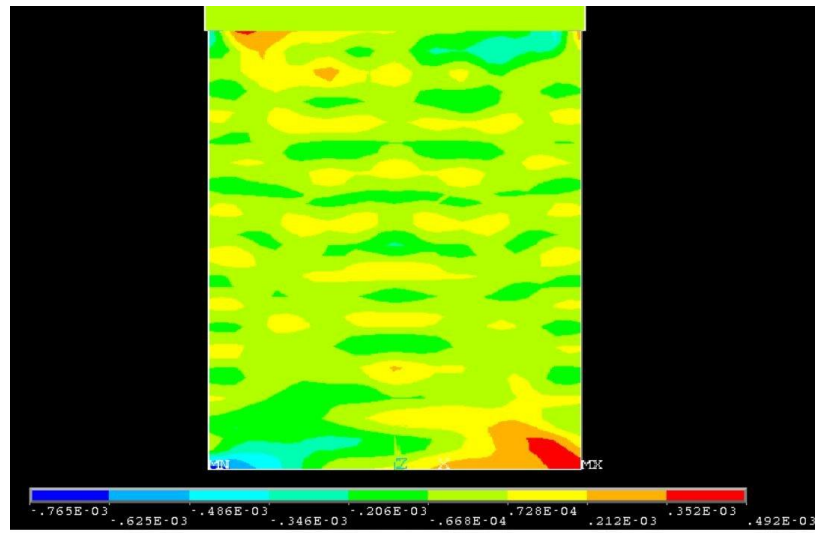


(c)

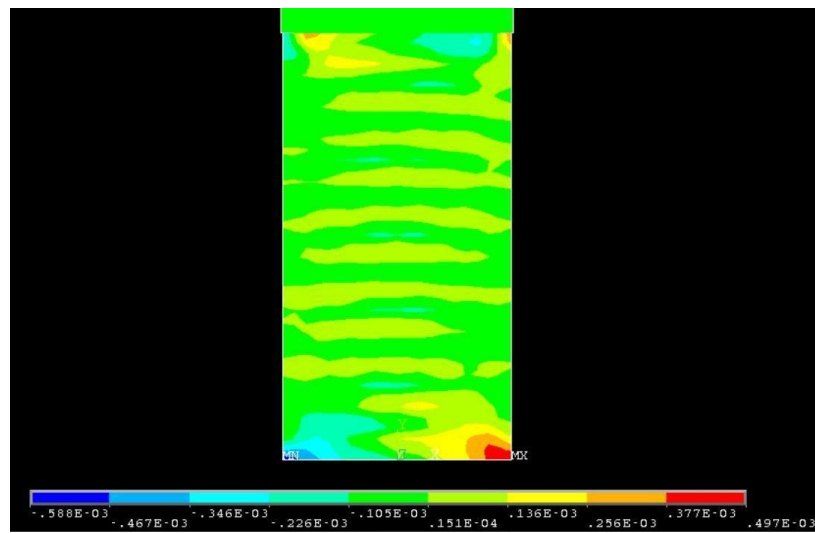


(d)

Fig. 4.7 continued.



(e)



(f)

Fig. 4.7 continued.

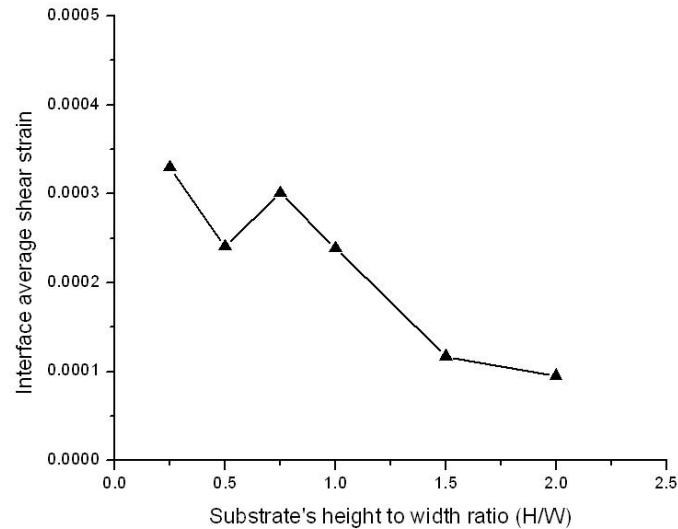


Fig. 4.8: Interface average shear strain vs. height-to-width ratio at the 750th cycle.

4.4 Validation

The simulated friction distributions have confirmed and explained the experimental observation that ultrasonic consolidation cannot be achieved for free standing substrates with a height-to-width ratio higher than 1.0. However, the simulation has not answered why the friction stress decreases to below the critical level for bond formation if the height of the substrate is greater than a certain value. In addition, the numerical simulation has not provided an explanation for the banded stress/strain distributions in the substrate.

To understand the wave interference during ultrasonic consolidation and to provide a validation of the FEM model, a vibration analysis of a 2-D ultrasound wave model is conducted. Consider a square domain with dimensions of $a \times a$, where a is the half width of the substrate in the FEM model (Figure 4.9). Let u and v represent the displacements in the x and y directions, respectively. The edge ($y = 0$) is fixed, whereas the edge ($x = 0$) is free. Based on the FE simulation results, the edge ($x = a$) is anti-symmetric. The edge ($y = 1$) is forced to oscillate along the X-direction at a frequency of 20 kHz with different vibration amplitudes. To simplify this vibration problem, a free vibration condition with a

uniform vibration amplitude ($16 \mu\text{m}$) is applied by ignoring the boundary condition on the edge ($y = a$) and assuming v equals to zero anywhere.

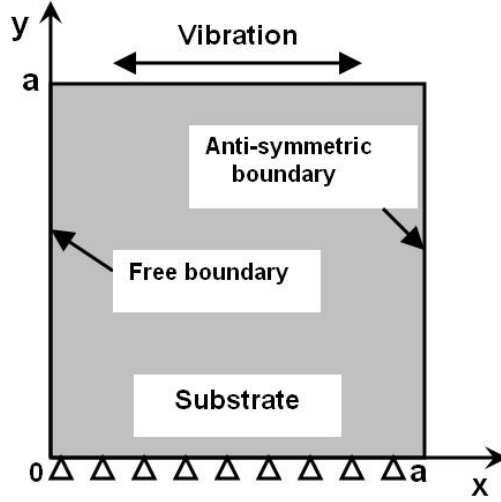


Fig. 4.9: A model for the analysis of wave traveling in the substrate.

The governing vibration equation can be solved using the Ritz method, in which the displacements u can be represented by algebraic polynomials, shown as follows:

$$u(x, y, t) = U(x, y)A_m \sin(\omega t), \quad (4.1)$$

$$U(x, y) = \sum_{i=0}^I \sum_{j=1}^J A_{ij} x^i y^j, \quad (4.2)$$

where A_m is the oscillation amplitude ($16 \mu\text{m}$), ω is the angular frequency ($2\pi f$), f is the ultrasound frequency (20KHz), t is time, A_{ij} are the coefficients. The boundary conditions on the edge ($y = 0$) are exactly satisfied by such a choice of functions.

The vibrational kinetic energy of this square domain is given by

$$T = \frac{\rho}{2} \int_0^a \int_0^a \dot{u}^2 dx dy, \quad (4.3)$$

where ρ is the mass per unit volume and $\dot{(\cdot)}$ represents the time derivative. Substituting Equation 4.1 into 4.3, the maximum kinetic energy during a full vibration cycle is

$$T_{max} = \frac{A_m^2 \rho \omega^2}{2} \int_0^a \int_0^a U^2 dx dy. \quad (4.4)$$

The strain energy during oscillation is

$$V = \frac{1}{2} \int_0^a \int_0^a (\sigma_x \epsilon_x + \tau_{xy} \gamma_{xy}) dx dy, \quad (4.5)$$

where σ_x and ϵ_x are the normal stress and strain, τ_{xy} and γ_{xy} are the shear stress and strain, respectively.

Constitutive stress-strain and strain-displacement equations are

$$\sigma_x = \frac{\nu E \epsilon_x}{(1 + \nu)(1 - 2\nu)}, \quad \tau_{xy} = G \gamma_{xy}, \quad (4.6)$$

$$\epsilon_x = \frac{\partial u}{\partial x}, \quad \gamma_{xy} = \frac{\partial u}{\partial y}, \quad (4.7)$$

where ν is Poisson's ratio, E is Young's modulus, G is the shear modulus. The maximum strain energy during an oscillation cycle can be given as a function of $U(x, y)$ by using Equations 4.1, and 4.5-4.7.

The functional L_{max} ($\equiv T_{max} - V_{max}$) is minimized by

$$\frac{\partial L_{max}}{\partial A_{ij}} = 0 \quad (i = 0 \dots I; j = 1 \dots J). \quad (4.8)$$

After the minimization of L_{max} , the yielded set of $(I + 1)J$ homogeneous, linear, algebraic equations are solved for the nontrivial solutions under the condition of $I = J = 9$.

The obtained displacement u is applied to calculate the shear strain, plotted in Figure 4.10. The vibration model has provided a steady distribution of shear strain for the substrate with a 0.5 H/W ratio, which corresponds to that in Figure 4.6e and f. By comparison, two common characters are found: (1) Both positive and negative shear strain

concentrations occur in the substrate. (2) After half of the vibration period, those strain concentrations will switch their locations, while keeping their magnitude almost unchanged. A good agreement (with the maximum difference of 13.0%) is also found by a quantitative comparison of those strain concentrations closest to the contact surface for the FEM and vibration models, shown in Table 4.2. The banded shear strain distributions and cyclic strain concentrations are indeed caused by the ultrasound wave interference in the substrate. This wave-based understanding helps to explain many process problems and reveal the mechanisms of ultrasonic consolidation.

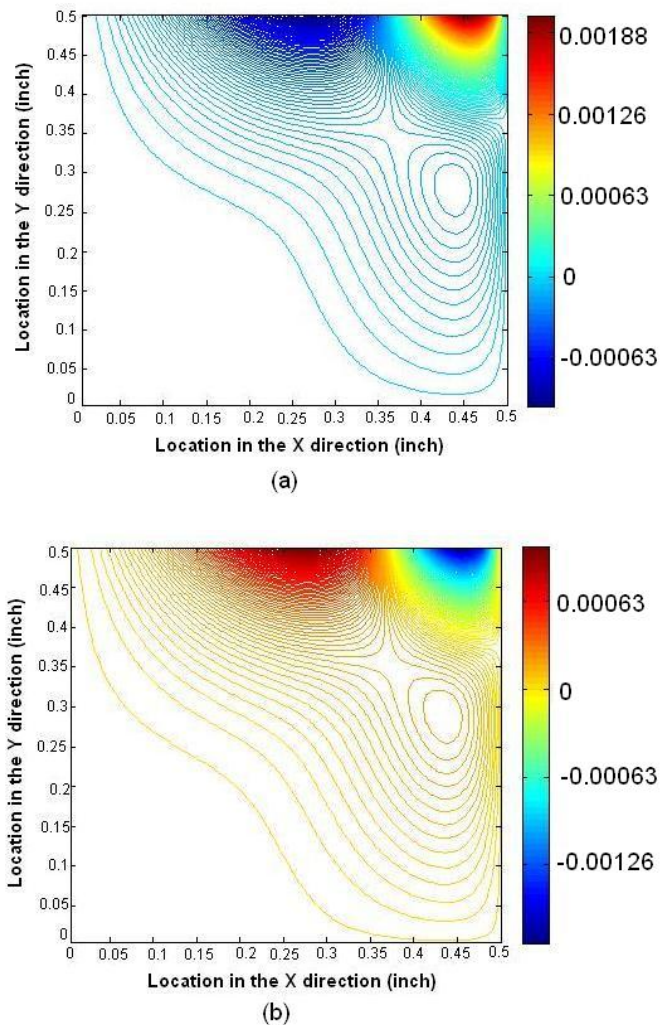


Fig. 4.10: Distribution of shear strain for the substrate with a 0.5 height-to-width ratio: (a) sonotrode moves to the right, and (b) sonotrode moves to the left.

Table 4.2: Comparison of Shear Strain Between FEM and Vibration Models (H/W=0.5)

	Shear Strain (right)	Shear Strain(left)
FEM Model	-0.00104	+0.00108
Vibration Model	-0.00094	+0.00094
Difference (%)	9.6	13.0

Chapter 5

Bond Strength Characterization and Improvement of UW Parts

5.1 Background of Bond Strength Determination

The overall strength of a laminated structure depends critically on the bond strength between laminated layers. Bond strength measurement has been widely employed to evaluate strength for thin film coating, dental adhesion, and composite fabrication. There are several methods to experimentally evaluate the bond strength for laminated structures. These methods can be broadly classified into five categories depending upon the debonding mode: (a) tensile test (pull-off test), (b) bending test, (c) peel test, (d) scratch test, and (e) ultrasonic test.

The tensile test [60] is the most commonly used method for the bond strength measurement of thin layers. An increasing tensile force is applied perpendicular to the bonding interface until the interface fails. Some criteria, such as the maximum force or area of detachment, are measured and used to evaluate the bond strength. Two requirements of this method include: (1) the interface bond strength must be lower than the strength of the adhesive media with which an extension can be joined to the thin layer for gripping and (2) a uniform load should be applied across the interface. Two drawbacks for the direct tensile test are: (a) a tensile test involves a complex mixture of tensile and shear forces that make the results difficult to interpret and (b) a tensile test is limited by the strength of the adhesive media.

In the bend test [61, 62], a load is applied to a sample fixed in a support. Crack nucleation and propagation are detected from the disruption of the load-deflection curve or by an acoustic emission detector. In the peel test [63, 64], a layer is peeled from a substrate

by a peeling load. For both the bend and peel tests, alignment of the specimen with the axis of the testing machine is not required, but both tests are limited to soft coatings or films. The peel test has less substrate distortion resulted in the test than the tensile and bend test. However, the results of peel tests are difficult to interpret, and are not directly comparable with the other bond strength testing methods, unless a layer can be completely debonded from the substrate. This limits the peel test to be suitable for layers with poor bond strength.

In the scratch test [65–67], an increasing load is applied to scratch a layer by an indenter tip. The critical load for failure is recorded for the analysis of bond strength. In the ultrasonic test [68–72], the ultrasonic wave is reflected from the interface when a bond is not perfect. The amplitude of the reflected wave is measured and correlated with the bond strength. The difficulty with these two methods is the interpretation of results and qualitative correlation to the bond strength. The difficulty arises from the many variables involved in the measurement and analysis procedures. Some other methods, such as lateral force-sensing microindentation, and laser spallation tests have also been developed for the evaluation of bond strength [73, 74].

We propose a new method for bond strength measurement, the push-pin experiment and finite element simulation (PEFE) for laminated structures. PEFE has the following advantages: (a) It is simple to set up, and efficient to test. (b) No adhesive is required for specimen extension, therefore there is no limitation due to the adhesive media or interface strength. (c) For multilayer structures, it has the capability of performing the bond strength measurement of each layer. (d) It is straight forward to interpret the experimental results by decoupling the shear force and obtaining the bond strength, with the help of finite element analysis. (e) A master calibration curve can be prepared to identify bond strength.

Bond strength is a critical issue that can greatly influence the development of UW and its expansion in existing and new fields. However, there have been no bond strength data from direct measurements on UW structures. The reason seems to be the lack of a suitable method to evaluate the bond strength between bonded layers. Attempts to conduct tensile

testing for the bond strength have not been successful, because the specimen fails at the grips, and it is difficult to control the failure location to a specific bond interface. The motivation for the current work, as well as the first application, has been to measure the bond strength of structures built by UW.

5.2 Methodology of PEFE

PEFE is defined as push-pin experiment and finite element simulation to determine the bond strength measurement for laminated structures. PEFE involves the following steps: (1) conduct the push-pin experiment and record the force and displacement data; (2) develop a layer-structured FE model with the same dimensions to simulate the push-pin experiment; (3) try different property coefficients for the bond zone until the curve of force versus displacement from the simulation match those from the experiment; and (4) pick the maximum stress normal to the bonded areas from simulation results as the bond strength.

5.2.1 Push-pin Experiment

Experiment Setup

Figure 5.1 shows the schematic and setup of the push-pin experiment on the GleebleTM 1500D thermomechanical simulator. In the push-pin test, the left surface of the baseplate is fixed and a hole is machined from the right surface of the baseplate along the direction normal to the bond areas. The depth of the hole is determined by the depth of the interface of which the bond strength is to be measured. A uniform load is applied with a given strain rate on the specimen by a push rod until the specimen fails. The history of force and displacement is recorded.

Specimen Design and Test Procedure

The push-pin experiment is designed to study the bond strength of UW, as influenced by process parameters in UW. The controlled process parameters include the sonotrode vibration amplitude, normal pressure, and sonotrode's travel velocity. For this study, four

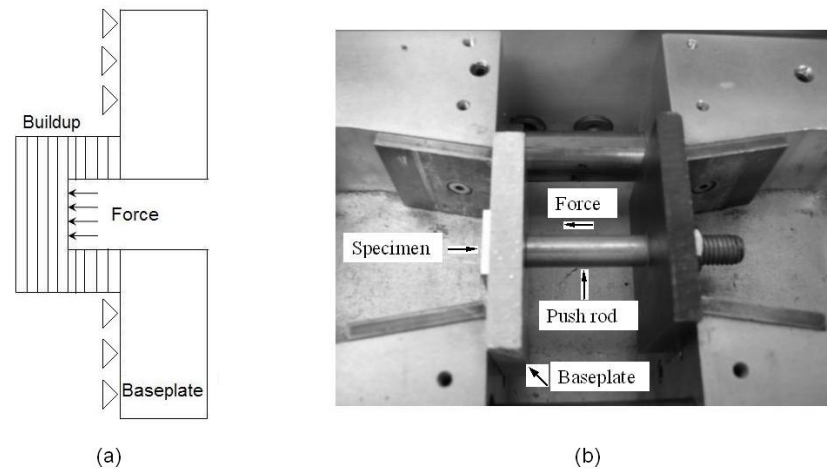


Fig. 5.1: Schematic of the push-pin experiment (a), and setup for the push-pin experiment on the Gleeble (b).

levels are chosen for each process parameter (Table 5.1), while keeping other process parameters constant. Samples are built on the SolidicaTM Formation UW system. The geometry of the multilayered buildup for the process parameter study is 22.9 mm wide, 25.4 mm long, and 16-layers high. The raw material in foil format is Al3003-H18 aluminum strip 0.1 mm thick, and 23.9 mm wide. A hole with 12.4 mm diameter was machined into the specimen. The diameter of push rod is 12.3 mm. The displacement rate used in the push-pin test is 0.42 mm/s. The tests are conducted on the GleebleTM 1500D thermomechanical simulator. To provide the necessary data for simulation, the mechanical properties and coefficient of friction of the raw aluminum foil Al3003-H18 have also been measured on the Gleeble.

PEFE allows bond strength measurement on bond interfaces through different depths of the hole machined into the substrate. This unique characteristic of PEFE provides a reliable way to perform bond strength measurement for parts with a very small number of layers. One layer buildup is able to be tested, with a backup metal plate glued to the top surface of the layer to prevent the push through by the push pin.

Table 5.1: Process Parameters for Ultrasonic Welding

Parameters	Level 1	Level 2	Level 3	Level 4
Sonotrode vibration amplitude (μm)	12	16	20	24
Sonotrode normal pressure (N)	800	1000	1400	1800
Sonotrode travel velocity (m/min)	1.02	1.27	1.42	1.53

Note: The pressure and velocity are 1800 N and 1.42 m/min for the study of amplitude; the amplitude and velocity are 16 μm and 1.42 m/min for the study of pressure; the pressure and amplitude are 1800 N and 16 μm for the study of velocity.

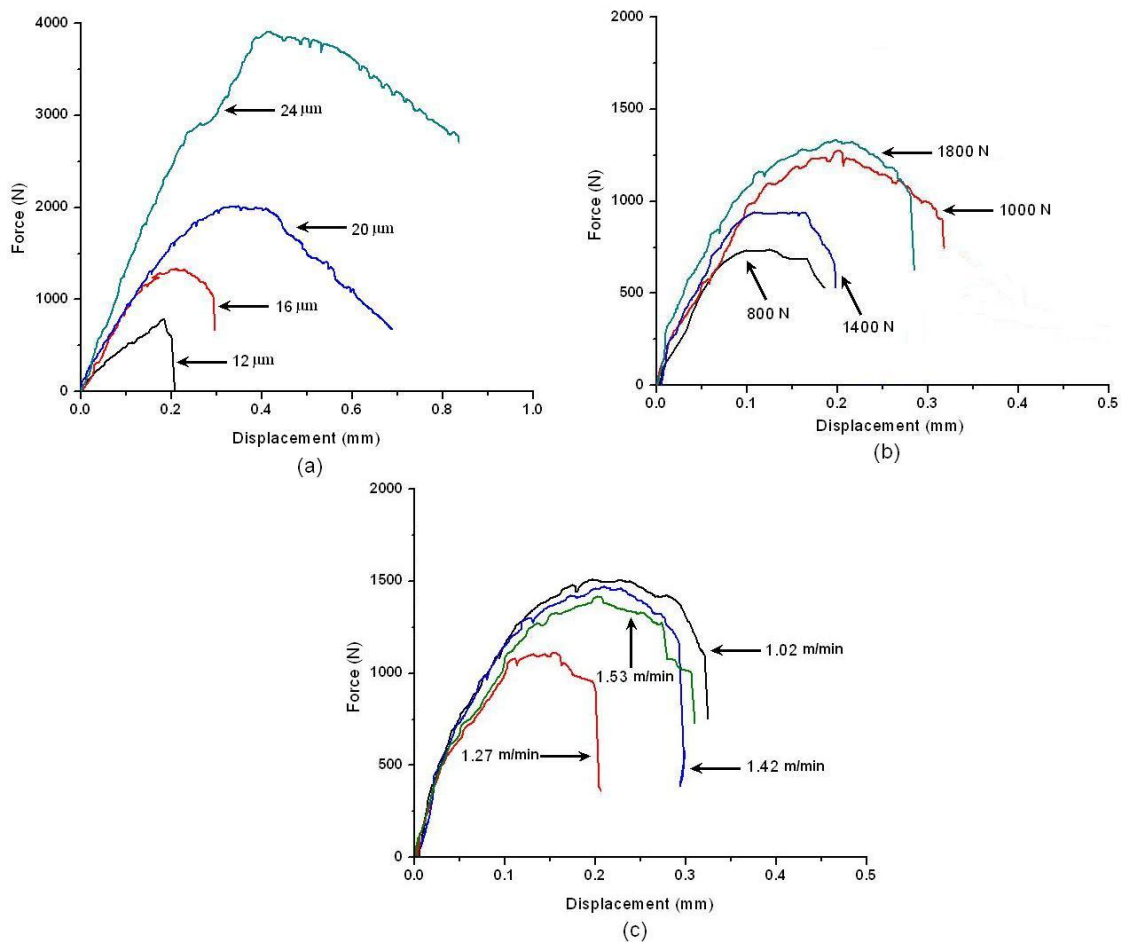


Fig. 5.2: Force vs. displacement curves from push-pin experiments on specimens made with varying UC process parameters: (a) vibration amplitude, (b) normal pressure, and (c) sonotrode's travel velocity.

Experimental Results

The results of push-pin experiments for the study of UW process parameters are shown in Figure 5.2. It can be seen that applied force increases nonlinearly with displacement. The force reaches its peak value and drops down quickly. Parts built by UW show brittle fracture normal to bonded areas with low ductility. A typical tested push-pin specimen is shown in Figure 5.3. Most of the specimens fracture along the bond line by separation of the foil interface. A fractography with porous features can be observed, indicating only a fraction of the bond interface has been bonded. The slopes and peak values on these curves of force versus displacement are revealing the real property differences among samples. Higher slope and peak value represent greater material stiffness and fracture strength. These results indicate that the push-pin experiment is a sensitive method for bond strength measurement.

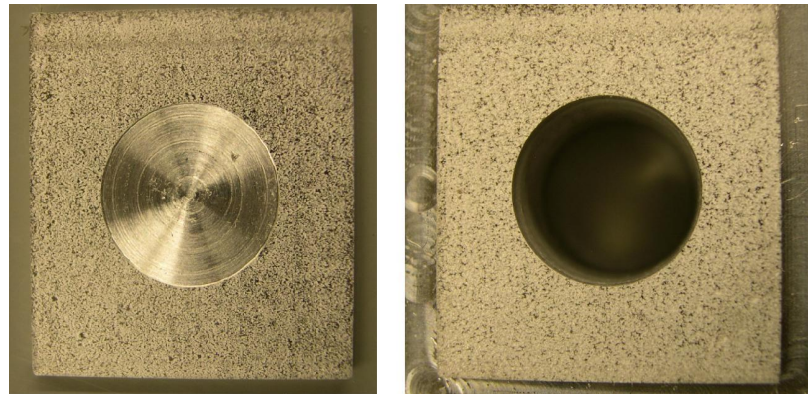


Fig. 5.3: Typical tested push-pin specimen.

The push-pin experiment also improves the understanding of process parameters on bonding strength. Among the selected process parameters, vibration amplitude has the most significant influence on the slope and peak value of force versus displacement curve. That means vibration amplitude is a critical factor that determines properties of parts built by the UW process. Larger vibration amplitude resulted in a stronger UW bond, which is shown by the slope and peak value increasing with vibration amplitude. In the selected range of normal pressure and sonotrode's travel velocity, the curves of force versus

displacement show similar slopes, but different peak values. This implies that the major effect of normal pressure and travel velocity is on the fracture strength, not the structure stiffness. The peak values showed a nonlinear relation with parameters of normal pressure and travel velocity. Experiment and simulation of the UW process showed that too large or too small of a normal pressure, and too high of a travel velocity will produce defective UW bonding. On the other hand, too slow of a velocity will reduce the efficiency of the manufacturing process. That is shown in the results (Figure 5.2b-c). It is seen that the peak force values for 800 N pressure and 1.53 m/ min velocity are smaller than the optimized process parameters, which are 1800 N and 1.42 m/min, respectively.

5.2.2 Finite Element Simulation

The Finite Element Model

Due to the symmetry of specimens, a 2D FE model with a layered structure has been developed, based on the commercial ANSYS software, to simulate the push-pin experiment. Figure 5.4 shows the layered and meshed models. In the layered part, narrow strips are located between layers to simulate the bond zones. The layered part and bond zones have been finer meshed. The dimensions and number of layers are the same as those in push-pin specimen in order to match the experimental and simulation results. The bond zone has been estimated to be 10-20 μm thick based on microstructure observation. Since the bond zones are not as strong as the layers, they are assumed to have a proportion of the properties of raw material. Proportional coefficients of properties are used to correlate the unknown properties of the bond zones with those of the raw material.

In simulation, the fixed and axial symmetric boundary conditions are used on the edges of AB and DE, respectively. A bilinear kinematic hardening rule has been used to simulate the nonlinear material behavior. A load with the same loading rate as in the push-pin experiment is uniformly applied on the edge of CD (Figure 5.4). Different combinations of property coefficients of elastic modulus, yield stress, and work hardening for the bond zones were tested until the curve of force versus displacement from FE simulation matches that

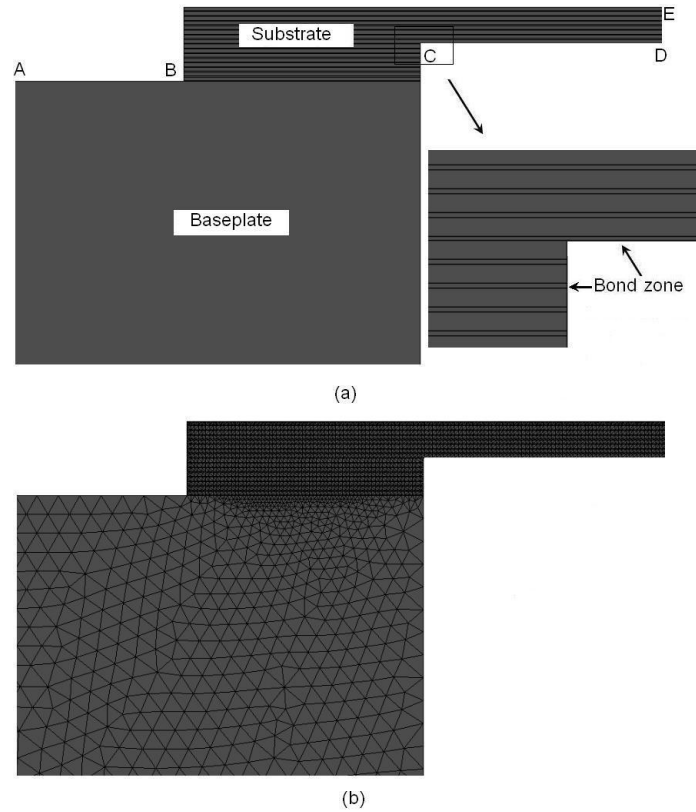


Fig. 5.4: Finite element model of the push-pin specimen (a), and meshed model (b).

from push-pin experiment. Therefore, the bond strength can be identified by this FE-based “experiment”.

Simulation Results

From the FE simulation results, the stress distribution along the edge CD is calculated for each displacement. Then multiplying the cross section area of the push rod, the average force on edge CD is calculated for all the displacement levels. The coefficients of elastic modulus used in simulations for different parameter sets are shown in Table 5.2, which is found to be the most sensitive in changing the shape of the force versus displacement curve. Figure 5.5 shows the comparison of force versus displacement from experiment and simulation for all the parameter sets. It is clearly seen that, when the appropriate bond strength properties are fed to the FE model, a matching curve can be achieved. The

simulation stops at the peak force, based on the assumption that the peak force is the onset of failure. After the peak force is reached, the crack will propagate quickly and the resistance will decrease.

Table 5.2: Coefficients of Elastic Modulus for the Bond Layer

Vibration Amplitude (μm)	12	16	20	24
Coefficient ($\times 10^{-4}$)	1.1	2.8	3.3	5.0
Normal Pressure (N)	800	1000	1400	1800
Coefficient ($\times 10^{-4}$)	1.9	2.4	2.5	2.8
Travel velocity of sonotrode (m/min)	1.02	1.27	1.42	1.53
Coefficient ($\times 10^{-4}$)	2.9	2.4	2.8	2.5

The von Mises yield criterion is used in this nonlinear simulation of FE model and the typical distribution of von Mises strain is shown in Figure 5.6a. The maximum von Mises strain is located at the corner of the machined inner hole. In the push-pin experiment, all specimens fail from the corner of the machined inner hole. Therefore, the depth of the hole can accurately determine the location of the failure. The distribution of von Mises strain confirms this experimental observation.

The debonding process is a complicated effects of normal and shear stresses on the fracture plane. However, these effects can be simplified by PEFE, because the stress normal to the bond areas is more importantly than the other stress components. This means that this stress will most likely dominate the crack propagation process. This can be proven by the stress distribution normal to the bond areas, shown in Figure 5.6b. This stress distribution is similar to that of the von Mises plastic strain, and the maximum tensile stress is also located at the corner bonding interface. Therefore, the stress data obtained from the FE simulation, which match the push-pin experiment data, can be used to evaluate the bond strength.

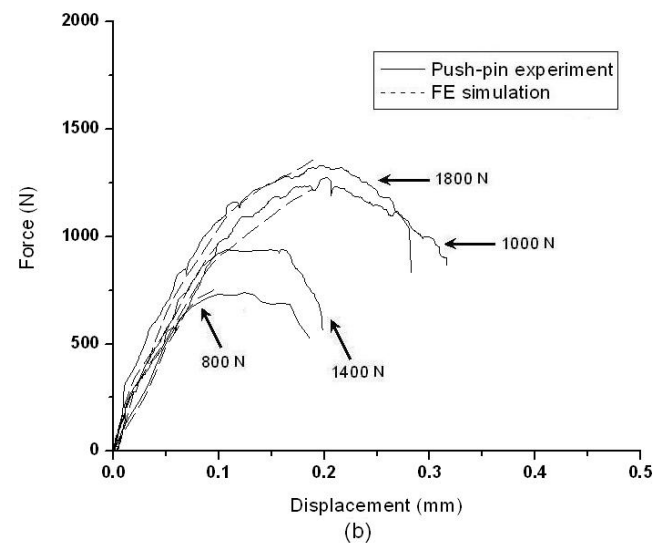
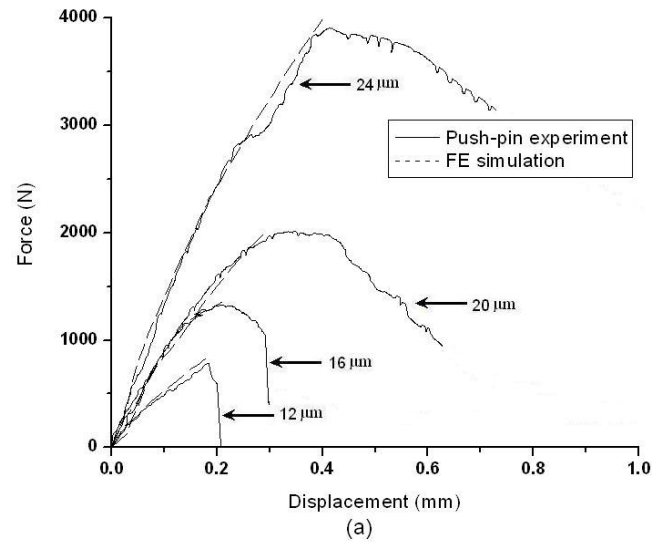


Fig. 5.5: Comparison of force vs. displacement curves from the push-pin experiment and from the FE simulation for specimens made using various process parameters: (a) vibration amplitude, (b) normal pressure, and (c) sonotrode's travel velocity.

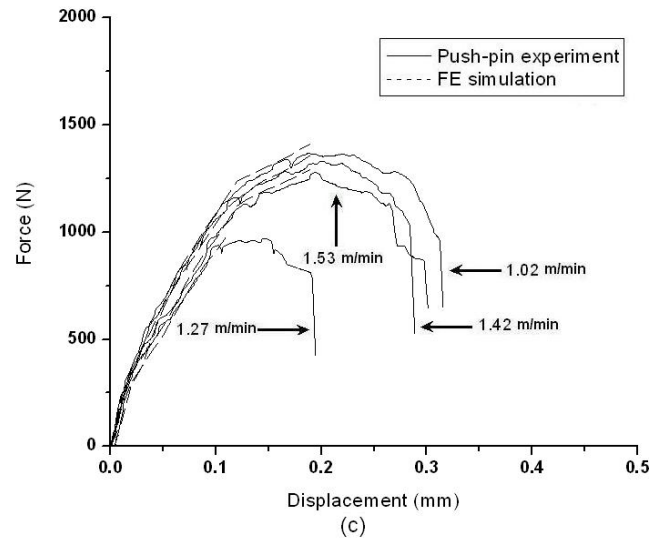


Fig. 5.5 continued.

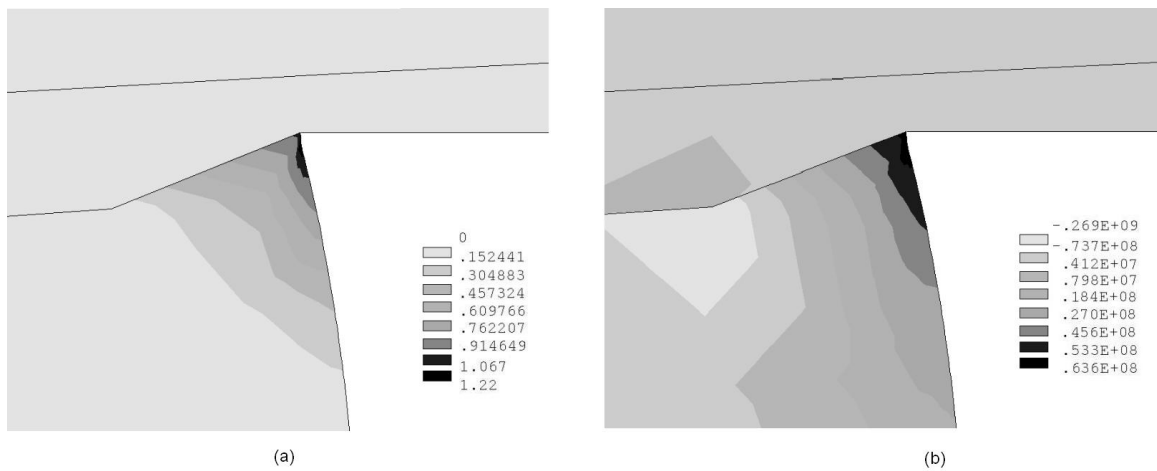


Fig. 5.6: A typical von Mises strain distribution near the corner of the push-pin hole (a) and a typical z-direction (specimen height direction) stress σ_{zz} distribution at the same location (Unit : Pa) (b).

5.2.3 Determination of Bond Strength by PEFE

The debonding at the interface involves the coupling of stresses that make the failure analysis for layered structures difficult. Brewer and Lagace [75] proposed a general criterion, the quadratic delamination criterion (QDC), for the prediction of interface debond that takes the coupling effect of normal and shear stresses into account. Based on QDC, the onset of interfacial failure can be predicted with

$$\left(\frac{\bar{\sigma}_{xz}}{Z^{sx}}\right)^2 + \left(\frac{\bar{\sigma}_{yz}}{Z^{sy}}\right)^2 + \left(\frac{\bar{\sigma}_{zz}^t}{Z^t}\right)^2 + \left(\frac{\bar{\sigma}_{zz}^c}{Z^c}\right)^2 = 1 \quad (5.1)$$

where $\bar{\sigma}_{ij}$ = average of a stress component;

Z^t = tensile interlaminar normal strength;

Z^c = compressive interlaminar normal strength;

Z^{sx} = interlaminar shear strength for σ_{xz} stresses; and

Z^{sy} = interlaminar shear strength for σ_{yz} stresses.

When one stress component, normal or shear stress, is dominant in the failure process, Equation 5.1 can be significantly simplified. We assume that the debond failure in UW specimens is caused mainly by normal tensile stress at the failure interface. Therefore, all terms on the left hand side of Equation 5.1 can be ignored, except the Z^t term. With this simplification Equation 5.1 becomes $\bar{\sigma}_{zz}^t = Z^t$, the maximum stress failure criterion. Compared with QDC, the maximum stress failure criterion [76,77] is an effective and popular approach for engineering failure analysis. In the push-pin experiment, the peak force has been found and taken as the starting point for debonding failure. The maximum stress σ_{zz}^t corresponding to the peak force is calculated from the FE simulation and taken as the bond strength Z^t .

Figure 5.7 shows the determined bond strength Z^t of UW specimens using the proposed PEFE method. The highest bond strength for parts made with the parameters discussed in this study is identified to be 180 MPa, which is 75% of the UTS of the raw material. For the first time, the bond strength of UW processed components can be positively determined using PEFE.

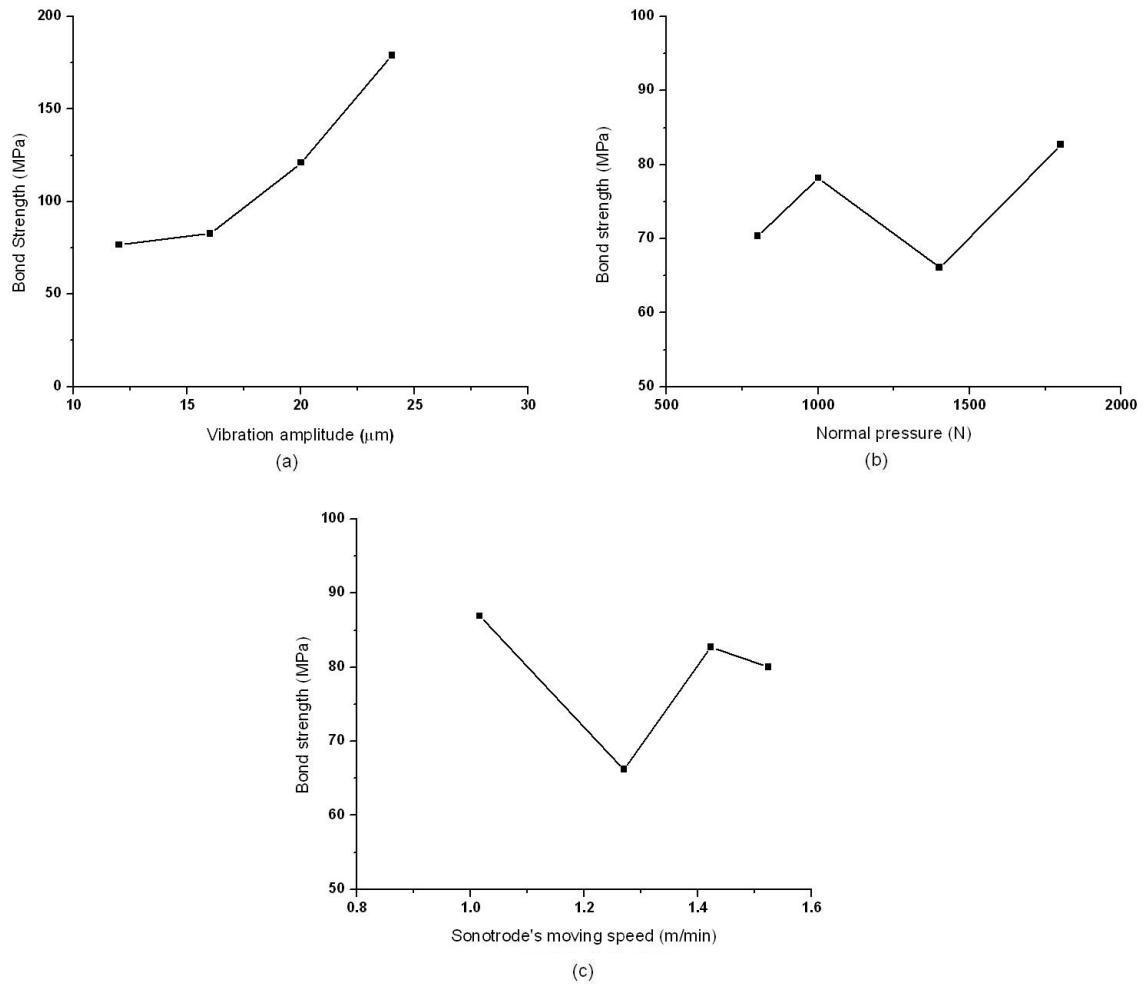


Fig. 5.7: Bond strength evaluated by PEFE varies with process parameters: (a) vibration amplitude, (b) normal pressure, and (c) sonotrode's travel velocity.

It should be noted that while the bond strength identified by PEFE does seem to be reasonable, PEFE can only be considered as a semi-quantitative method. The process involves inversely identifying, through simulation, at least two bond strength parameters (strength and modulus), therefore, there is always the “uniqueness” question of the parameters identified. However, since there is no better way for evaluating the bond strength for UW, the bond strength identified by PEFE is better than other qualitative methods, such as the peeling test. One way to overcome this semi-quantitative limit of the PEFE method is to calibrate the bond strength identified by another direct measurement method.

Based on Figure 5.7a, an increasing vibration amplitude increases the bond strength and it has more significant effect on bond strength/formation than the other two parameters (pressure and travel velocity). In the actual manufacturing process of UW, the optimum parameter set used were 16 μm vibration amplitude, 1800 N normal pressure, and 1.42 m/min travel velocity. In Figure 5.7b, normal pressure of 1800 N resulted in the best bond strength, while the lowest pressure of 800 N resulted in the worst bond strength. Based on our understanding of the UW process, to a high normal pressure will decrease bond strength, because the vibrational motion of the parts is constrained, and not enough friction heat will be generated. Thus, intermediate normal pressure, in the range of 1800 N, is the optimal for higher quality UW bonding. Figure 5.7c shows that the lowest travel velocity has the highest bond strength. Although the lowest moving velocity is optimal for bond formation, it also reduces the working efficiency. Therefore, considering the bond strength and working efficiency, 1.42 m/min is the optimal choice for the range of parameters considered in this study.

5.2.4 Correlation of Bond Strength by PEFE with Percentage of Bonded Area

Figure 5.8 shows the typical morphology of fractured surfaces for both high and low quality UW bonding conditions discussed in this study. A difference in the percentage of bonded area exists for these two bonding conditions. Most of the area is bonded for the better UW bonds, while large un-bonded areas exist in the low quality UW bonds. The percentage of bonded area is a reliable way of evaluating the bond formation and bond strength. The bulk bond strength can be predicated by the averaged percentage of bonded area of the entire bonded surfaces. The bond strength in a local region can also be evaluated from the percentage of bonded area of that local region. The bulk averaged percentage of bonded areas have been measured with digital images of fractured surfaces, and converted to bond strength. The bond strength from area measurements is defined as the averaged percentage of bonded area times raw material's UTS. The results of bond strength evaluated by percentage of bonded areas are plotted in Figure 5.9. Comparing Figures 5.7 and 5.9, the bond strength provided by these two approaches are showing the same trend for all three

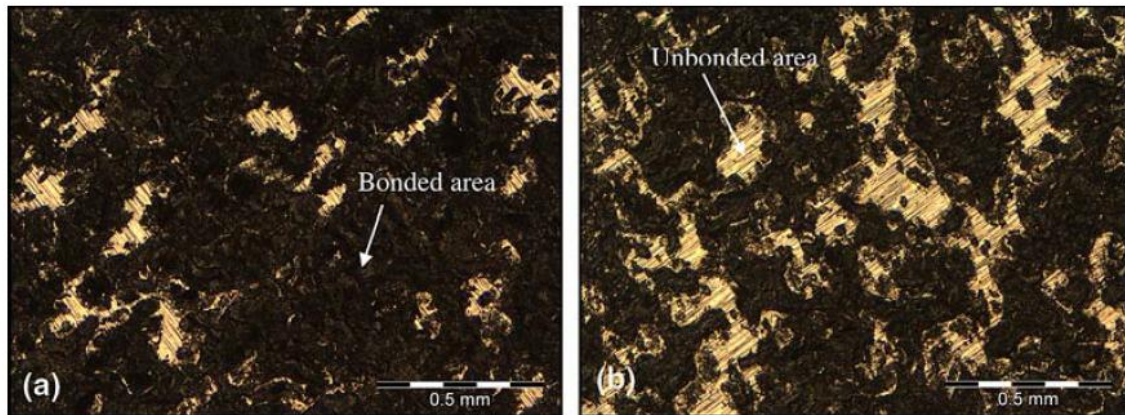


Fig. 5.8: Typical feature of fractured surfaces: (a) 91.6% bonded area, and (b) 74.9% bonded area.

process parameters. The similarity between these two sets of results is a strong validation for the PEFE method. The bond strength is probably over-evaluated by the percentage of bonded area method. It should be recognized that the UTS of the raw material used in the bond strength calculation from area is higher than the real bond strength.

Figure 5.10a shows the relation between the bond strength by PEFE and percentage of bonded area. With this curve, an estimate of bond strength can also be obtained by conducting a bonded area measurement of UW specimens. The bond strength increases rapidly when the percentage of bonded area is over 90%. This trend can be explained by a basic fracture mechanics analysis. The general form of the Griffith [78] equation is $\sigma_f = \sqrt{\frac{2Ew_f}{\pi a}}$, where σ_f is failure stress, E is elastic modulus, w_f is fracture energy, which could include plastic, viscoelastic, or viscoplastic effects, depending on materials, and a is half of the crack length. The failure stress is proportional to $\sqrt{\frac{1}{a}}$, when assuming that E and w_f are constants for a given material. In PEFE, the un-bonded area can be considered as a crack, so an equivalent crack length (a_{equiv}) is calculated from the percentage of bonded area by

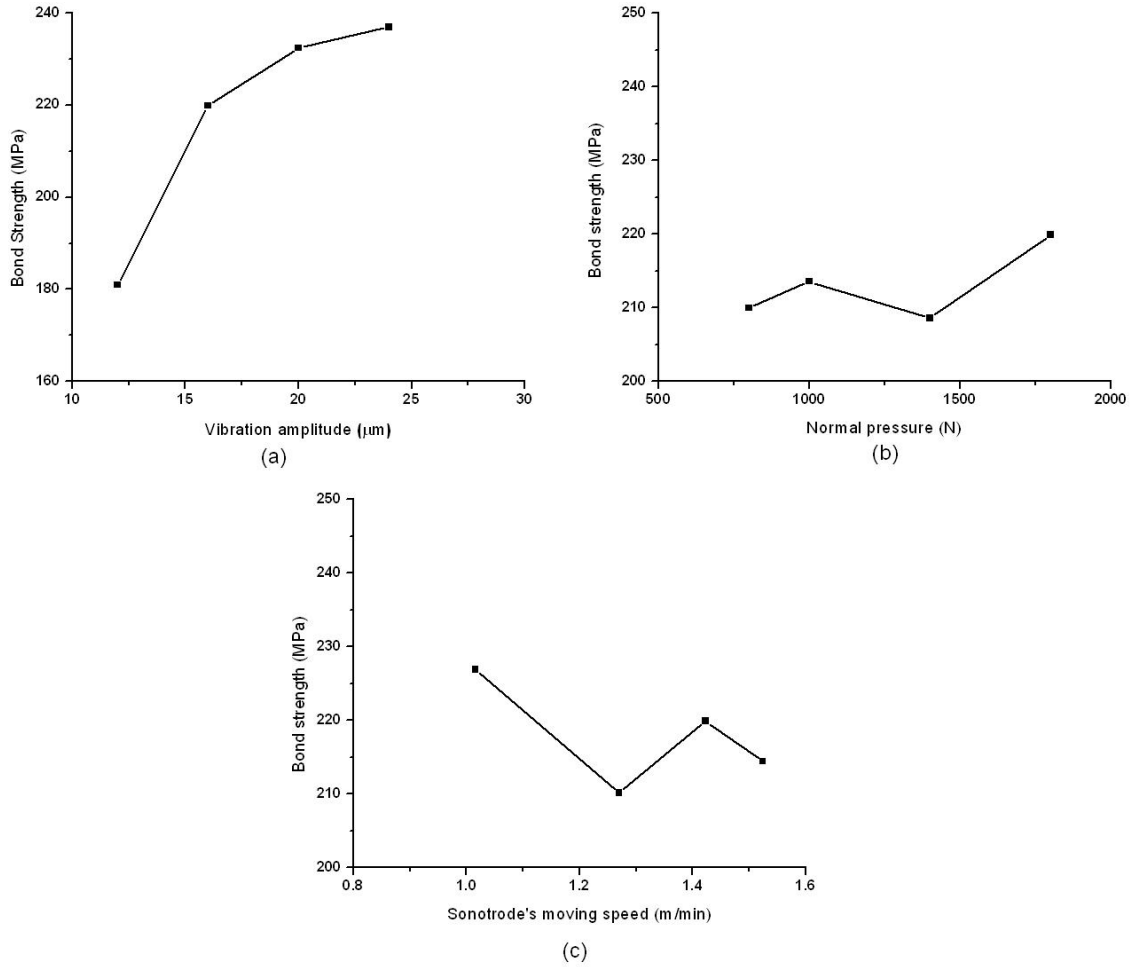


Fig. 5.9: Bond strength evaluated by bonded area to total area ratio varying with process parameters: (a) vibration amplitude, (b) normal pressure, and (c) sonotrode's travel velocity.

$$a_{eqv} = \frac{1}{2L} \sqrt{A(1 - r_b)} \quad (5.2)$$

where A is total cross section area, r_b is percentage of bonded area, and L is width of cross section.

Substituting the crack length a in the Griffith equation with a_{eqv} , the equivalent fracture strength for the defect-containing UW specimens can be written as

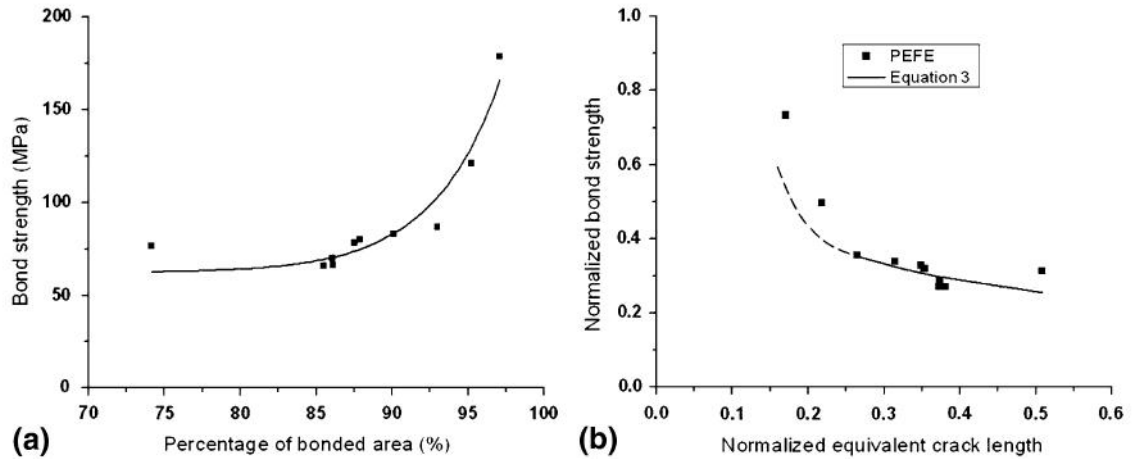


Fig. 5.10: Bond strength by PEFE vs. percentage of bonded area (a), and bond strength as a function of the equivalent crack length, with a comparison between strength obtained using PEFE and strength predicted using Equation 5.3 (b).

$$s_f = \sqrt{\frac{2Ew_f}{\pi a_{eqv}}} \quad (5.3)$$

If a constant fracture energy w_f is assumed, the curve of s_f (i.e., bond strength) versus a_{eqv} can be fitted to correlate with the bond strength determined by PEFE (Figure 5.10b). When the normalized equivalent crack length is longer than 0.25, the bond strength governed by Equation 5.2 matches well with that by PEFE. Although assumed as a constant, the fracture energy w_f should be viewed as a parameter, which may affect the bond strength.

5.3 Pressurized Post-weld Heat Treatment (PWHT)

Attempting to solve the problem of poor vertical bond strength of UW parts, a pressurized post-weld heat treatment is proposed. First, a systematic experimental matrix of PWHT is designed and conducted to study the coupled influences of temperature and time on bond strength. Based on the results of traditional PWHT and diffusion bonding mechanism, a series of pressurized PWHT tests have been conducted to find out the potential of bond strength improvement and the optimum parameter set, including pressure,

temperature and time.

5.3.1 Experiment Setup

Samples are built on a SolidicaTM Formation ultrasonic welding system under the same condition of 1800 N clamping force, 150 °C preheating temperature, 16 μm vibration amplitude, 20 kHz vibration frequency and 1.42 m/min traveling speed. The raw material in foil format is Al3003-H18 aluminum strip 0.1 mm thick, and 23.9 mm wide. The geometry of the multilayered buildup for the pressurized PWHT study is 22.9 mm wide, 25.4 mm long, and 16-layers high. Samples will be evenly heated and soaked for a certain period. 7×7 temperature-time matrix is designed and shown in Table 5.3 to study their coupling effects on bond strength improvement. All these samples are cooled in air after heat treatment. The bond strength of samples after PWHT are determined by the bond strength tests described above.

For the existence of unbonded areas on the bonding interface of UW, typical PWHT method may not be good enough to solve the weak-bonding problem of UW. Based on a diffusion bonding mechanism, the pressurized PWHT approach is employed to realize the full potential of bond strength improvement of UW. Since higher temperature and longer time facilitate the diffusion of bonding interface, the upper bounds of aluminum annealing temperature and time (450 °C and 3 h) are chosen and fixed to study the influence of applied pressure on bond strength. Once the optimum pressure is determined, it is applied to the pressurized PWHT tests for reducing the elevated temperature and shortening the time while maintaining the level of bond strength. Pressurized PWHT tests proceed on the GleebleTM 1500D thermomechanical simulator, and the PEFE test is used to estimate the bond strength.

Table 5.3: Temperature-time Matrix of PWHT Tests

Level number	1	2	3	4	5	6	7
Temperature (°C)	220	260	300	350	370	390	415
Time (h)	0.25	0.5	1	3	9	15	24

5.3.2 Experiment Results

PWHT

Based on the obtained systematic bond strength data, a 2-D map is plotted (Figure 5.11), which clearly exhibits the influence of temperature and time of PWHT on the bond strength. The lower left corner represents the low bond strength zone, corresponding to the low-temperature and short-period PWHT conditions. On the contrary, the upper right corner is the high bond strength zone, resulting from the opposite conditions. With the aid of this 2D map, the quantitative bond strength can be easily determined for an arbitrary combination of temperature and time within the studied range. The bond strength is increased by 35% from the as-welded condition using the best PWHT condition.

Since the UW bonding interface is severely plastic-deformed and residual stress/strain concentrated, the annealing influence of PWHT causes the stress relief and reduces the effect of strain hardening, which result in the decrease of hardness, but the improvement of ductility. After the classification of failure mode for the PWHT samples after PEFET testing, a 2D map of failure mode has been plotted, shown in Figure 5.12, and the critical PWHT condition of failure mode change is marked by a dash line. The PWHT condition of higher temperature and longer time has more potential to generate the ductile failure. Typical evolutions of peak load vs. displacement for the brittle and ductile failure modes are also shown in Figure 5.12. For the UW parts with a laminate structure, the ductile failure mode is characterized by a debonding surface across several layers, while the brittle failure mode occurs on a single bonding layer.

Pressurized PWHT

Although the bond strength of UW parts can be improved by the traditional PWHT method (35% from the as-welded condition), it is still far below the strength of the base material (only 48%). To achieve even higher bond strength, pressurized PWHT tests are conducted, and results are shown in Figures 5.13-5.16. Under the PWHT condition of 450 °C and 3 h, the bond strength increases with applied pressure (≤ 2.5 MPa) initially,

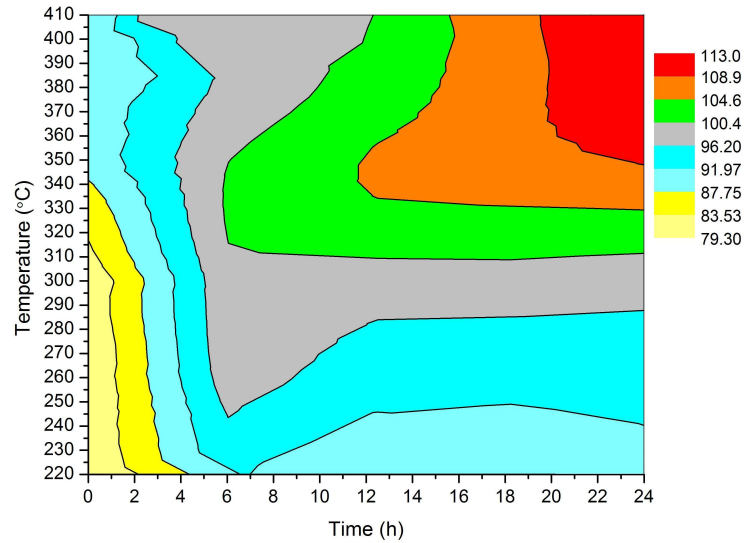


Fig. 5.11: 2D temperature-time bond strength map after PWHT (Unit: MPa).

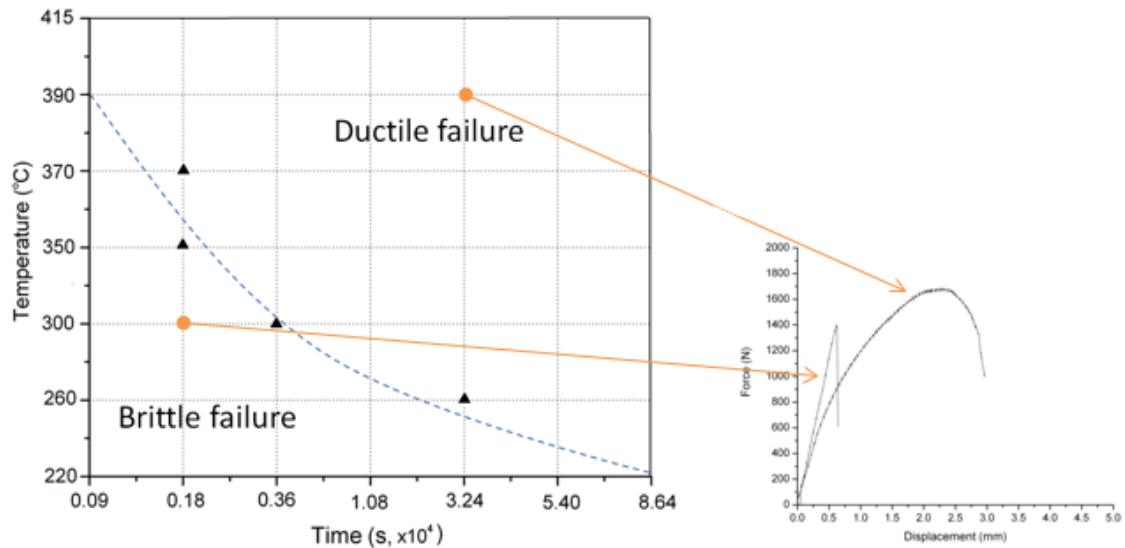


Fig. 5.12: 2D temperature-time failure mode map after PWHT.

and remains constant at higher stresses (>2.5 MPa), illustrated in Figure 5.13. Through applying pressure to the UW parts during PWHT tests, the bond strength can be modified significantly, up to 96% of the strength of the base material. When the applied pressure is above 2.5 MPa, the sample geometry starts to be destroyed permanently (Figure 5.14).

Considering both the bond strength and sample geometry, the pressure of 2.5 MPa is chosen as the optimum pressure value, which is applied to the rest of the tests determining the optimum temperature and time. The bond strength keeps increasing with temperature (Figure 5.15). Therefore, 450 °C is a good choice as the optimum temperature value. In Figure 5.16, the PWHT time can be shortened from 3 h to 1.5 h, while maintaining the bond strength level. Finally, the optimum parameters of pressurized PWHT tests are 2.5 MPa pressure, 450 °C temperature and 1.5 h time.

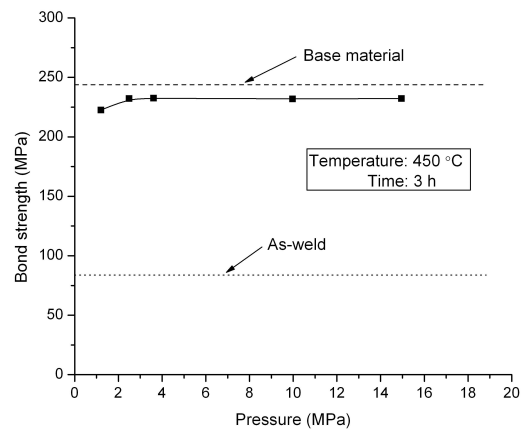


Fig. 5.13: Variation of bond strength after pressurized PWHT with pressure.

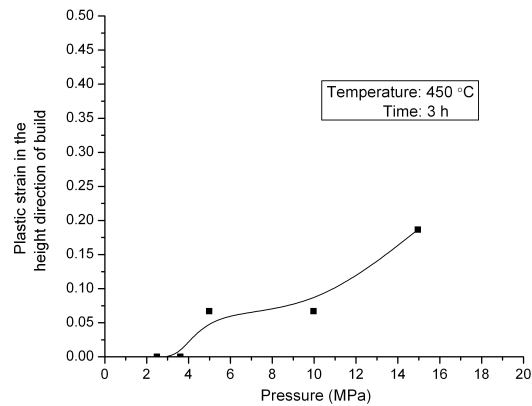


Fig. 5.14: Variation of buildup deformation after pressurized PWHT with pressure.

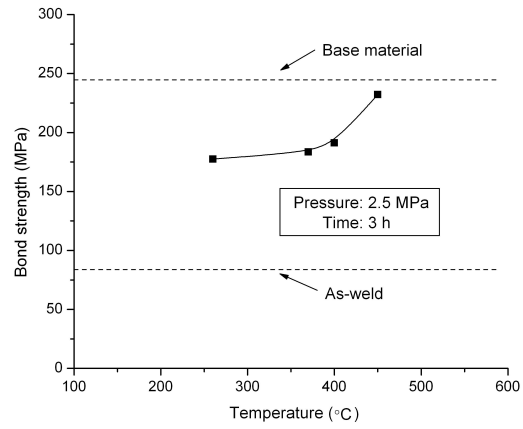


Fig. 5.15: Variation of bond strength after pressurized PWHT with temperature.

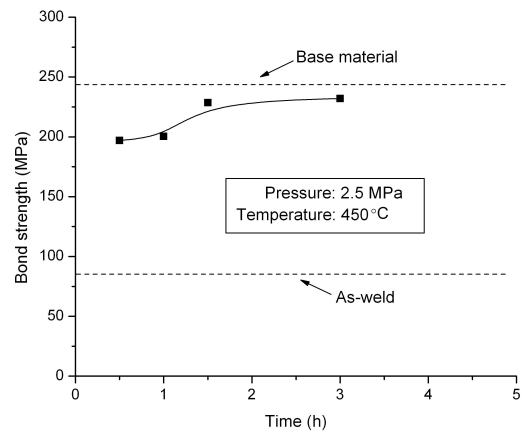


Fig. 5.16: Variation of bond strength after pressurized PWHT with time.

Chapter 6

Conclusions and Future Work

1. A new fully coupled thermomechanical finite element model has been developed to simulate the ultrasonic metal welding process. Insight has been gained on the evolution and interactions between the in-process thermal- mechanical variables. The model is able to explain experimental observations on the effect of process parameters on bond formation. The von Mises plastic strain obtained from the simulation has been quantitatively verified by experimentally measured data and correlated with the measured bonded area of ultrasonic joints. A possible mechanism for bond formation is proposed. Ultrasonic welding forms the bond through the following process: ultrasonic vibration generates shear stress, sliding, and heat at the interface through friction; the heat lowers the mechanical properties of the material, which enhances localized plastic deformation; and more plastic deformation generates more heat in addition to frictional heat. It is the severe, localized, plastic deformation at the bond region that is believed to be the major phenomenon for bond formation in ultrasonic welding.
2. With the increased substrate height, the amplitude of contact frictional stress decreases, while that of contact interface displacement increases. The reason for the decrease in the frictional stress at the contact interface for certain substrate heights is caused by the wave interference occurring in the substrate, which has been validated by an analytical vibration model. The formation of banded and cyclic stress-strain maxima in the substrate is also caused by superposition of vibrations.
3. A push-pin type, combined experimental and numerical method for evaluation of bond strength of laminated structures has been developed. The new method has been

validated by experimental bonded area data, and has been successfully applied to laminated structures produced using UW.

4. Bond strength for different UW parameter combinations has been evaluated. The best bond strength produced using the set of ultrasonic process parameters in this study is 75% of the UTS of the base material. An increase of vibration amplitude of the sonotrode will increase the bond strength. There is an optimized sonotrode pressure level for bond strength; both low and high pressures will cause a decrease in bond strength. Lower travel velocity of the sonotrode will generate higher bond strength.
5. PWHT tests have been conducted on the UW parts to explore the potential of bond strength improvement. Influences of PWHT temperature and time on bond strength are analyzed and illustrated in a 2D map of bond strength. Under the optimum PWHT condition (temperature and time only), the bond strength can be improved by 35% from the as-welded condition. The pressurized PWHT approach is originally proposed to improve the bond strength of UW parts. The results show that the modified bond strength, up to 96% of the strength of the base material, can be achieved under the optimum parameters of 2.5 MPa pressure, 450 °C temperature and 1.5 h time.

The results in this dissertation provide a strong foundation for future work in the mechanism study of ultrasonic metal welding. One area is in the long-term thermomechanical fully coupled simulation of bond formation process, in which the simulation of minimum 3000 vibration cycles is suggested. The results of plastic deformation at the bonding surface for the long-term simulation should be compared and correlated with the bonded area ratio. Another area is in applying the theory of dynamic recrystallization to quantitatively explain the generation of nano-sized grains at the bonding interface, combined with thermo-mechanical numerical results from FE simulation. It is also required to develop a method for directly measurements of the strain/stress or temperature during the bonding process to validate the numerical model of ultrasonic welding and to help the interpretation of the bonding mechanisms.

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