



An Introduction to Differential Geometry with Maple

PCMI - July 5, 2013

The Eguchi-Hanson metric

Synopsis

The Eguchi–Hanson metric is an example of a gravitational instanton (complete, Riemannian, Ricci-flat). In this worksheet we shall:

- define the Eguchi-Hanson metric and build an orthonormal frame for calculating its properties
- show that the metric is Ricci-flat and self-dual.
- show that the holonomy is $so(3)$
- explicitly prove that the metric admits a quaternionic-Kähler structure.
- calculate the Kähler potential for the metric

The Eguchi-Hanson metric

```
[M > restart:
```

```
> with(DifferentialGeometry): with(Tensor): with(LieAlgebras):
```

Define the coordinates with DGsetup.

```
> DGsetup([R, theta, phi, psi], M0);  
frame name: M0 (1.1)
```

Define a set of three 1-forms used to construct an orthogonal frame for the metric.

```
> s1 := evalDG(sin(psi)*dtheta - cos(psi)*sin(theta) *dphi);  
s1 := sin(ψ) dθ - cos(ψ) sin(θ) dφ (1.2)
```

```
> s2 := evalDG(cos(psi)*dtheta + sin(psi)*sin(theta)*dphi);  
s2 := cos(ψ) dθ + sin(ψ) sin(θ) dφ (1.3)
```

```
> s3 := evalDG(dpsi + cos(theta)*dphi);  
s3 := cos(θ) dφ + dψ (1.4)
```

Define an [anholonomic frame](#) on M_0 . We shall do all our calculations with respect to this frame.

```
[M > FD := FrameData([dR, s1, s2, s3], M);  
FD := [dθ1 = 0, dθ2 = -θ3 ∧ θ4, dθ3 = θ2 ∧ θ4, dθ4 = -θ2 ∧ θ3] (1.5)
```

Initialize this frame with [DGsetup](#).

$$\left[\begin{array}{l} \text{M} > \text{DGsetup}(\text{FD}, [\text{zeta0}, \text{zeta1}, \text{zeta2}, \text{zeta3}], [\text{sigma0}, \text{sigma1}, \\ \text{sigma2}, \text{sigma3}]); \\ \text{frame name: } M \end{array} \right. \quad (1.6)$$

Here is the Eguchi-Hansen metric in this frame

$$\left[\begin{array}{l} \text{M} > \text{g} := \text{evalDG}((1 - a/R^4)^{-1} * \text{sigma0} \&t \text{sigma0} + R^2/4 * (1 - \\ a/R^4) * \text{sigma3} \&t \text{sigma3} + R^2/4 * (\text{sigma1} \&t \text{sigma1} + \text{sigma2} \&t \\ \text{sigma2})); \\ g := \frac{R^4}{R^4 - a} \sigma_0 \otimes \sigma_0 + \frac{R^2}{4} \sigma_1 \otimes \sigma_1 + \frac{R^2}{4} \sigma_2 \otimes \sigma_2 + \frac{R^4 - a}{4 R^2} \sigma_3 \otimes \sigma_3 \end{array} \right. \quad (1.7)$$

Curvature

The Eguchi-Hansen metric is Ricci-flat and self-dual.

$$\left[\begin{array}{l} \text{M} > \text{RicciTensor}(g); \\ 0 \sigma_0 \otimes \sigma_0 \end{array} \right. \quad (2.1)$$

We calculate the [curvature tensor](#) and its [dual](#) and check that these are equal.

$$\left[\begin{array}{l} \text{M} > \text{C} := \text{CurvatureTensor}(g); \\ \text{M} > \text{C1} := \text{RaiseLowerIndices}(g, \text{C}, [1]); \\ \text{M} > \text{Cdual} := \text{DualCurvature}(g, \text{C}) \text{ assuming } R > 0; \\ \text{M} > \text{C1} \&minus \text{Cdual}; \\ 0 \sigma_0 \otimes \sigma_0 \otimes \sigma_0 \otimes \sigma_0 \end{array} \right. \quad (2.2)$$

Infinitesimal Holonomy

The [infinitesimal holonomy](#) is a Lie algebra of matrices which act on the tangent space of the manifold at a given point by isometries.

$$\left[\begin{array}{l} \text{M} > \text{H} := \text{InfinitesimalHolonomy}(g, [\text{R} = a, \text{theta} = \text{Pi}/2, \text{phi} = 0, \\ \text{psi} = 0]); \\ H := \left[\begin{array}{cccc} 0 & -\frac{a}{2 R^4} & 0 & 0 \\ \frac{2 a}{(R^4 - a) R^2} & 0 & 0 & 0 \\ 0 & 0 & 0 & -\frac{a}{R^5} \\ 0 & 0 & \frac{a}{(R^4 - a) R} & 0 \end{array} \right] \end{array} \right. \quad (3.1)$$

$$\left[\begin{array}{cccc} 0 & 0 & -\frac{a}{2R^4} & 0 \\ 0 & 0 & 0 & \frac{a}{R^5} \\ \frac{2a}{(R^4-a)R^2} & 0 & 0 & 0 \\ 0 & -\frac{a}{(R^4-a)R} & 0 & 0 \end{array} \right],$$

$$\left[\begin{array}{cccc} 0 & 0 & 0 & \frac{a(R^4-a)}{R^8} \\ 0 & 0 & \frac{2a}{R^5} & 0 \\ 0 & -\frac{2a}{R^5} & 0 & 0 \\ -\frac{4a}{(R^4-a)R^2} & 0 & 0 & 0 \end{array} \right]$$

Here are the structure equations for the holonomy algebra.

$$\begin{aligned} & \mathbf{M} > \mathbf{LD} := \mathbf{map}(\mathbf{simplify}, \mathbf{LieAlgebraData}(\mathbf{H}, \mathbf{hol})); \\ & \mathbf{LD} := \left[[e1, e2] = -\frac{a e3}{(R^4 - a) R}, [e1, e3] = \frac{4 a e2}{R^5}, [e2, e3] = -\frac{4 a e1}{R^5} \right] \end{aligned} \quad (3.2)$$

We can simplify these structure equations by scaling the vectors.

$$\begin{aligned} & \mathbf{M} > \mathbf{DGsetup}(\mathbf{LD}); \\ & \text{Lie algebra: hol} \end{aligned} \quad (3.3)$$

$$\begin{aligned} & \mathbf{hol} > \mathbf{LD2} := \mathbf{LieAlgebraData}(\mathbf{[r*e1, s*e2, t*e3]}, \mathbf{alg}); \\ & \mathbf{LD2} := \left[[e1, e2] = -\frac{r s a e3}{t (R^4 - a) R}, [e1, e3] = \frac{4 r t a e2}{s R^5}, [e2, e3] = -\frac{4 s t a e1}{r R^5} \right] \end{aligned} \quad (3.4)$$

$$\begin{aligned} & \mathbf{hol} > \mathbf{S} := \mathbf{solve}(\mathbf{subs}(\{-\mathbf{r*s*a}/(\mathbf{t*(R^4-a)*R}) = \mathbf{1}, \mathbf{4*r*t*a}/(\mathbf{s*R^5}) = \mathbf{-1}, \mathbf{-4*s*t*a}/(\mathbf{r*R^5}) = \mathbf{1}\}), \{\mathbf{r}, \mathbf{s}, \mathbf{t}\}, \mathbf{explicit}); \\ & \mathbf{S} := \left\{ r = \frac{\sqrt{R^4 - a} R^3}{2 a}, s = \frac{\sqrt{R^4 - a} R^3}{2 a}, t = -\frac{R^5}{4 a} \right\}, \left\{ r = -\frac{\sqrt{R^4 - a} R^3}{2 a}, s = \right. \\ & \left. -\frac{\sqrt{R^4 - a} R^3}{2 a}, t = -\frac{R^5}{4 a} \right\}, \left\{ r = \frac{\sqrt{R^4 - a} R^3}{2 a}, s = -\frac{\sqrt{R^4 - a} R^3}{2 a}, t = \frac{R^5}{4 a} \right\}, \left\{ r = \right. \\ & \left. -\frac{\sqrt{R^4 - a} R^3}{2 a}, s = \frac{\sqrt{R^4 - a} R^3}{2 a}, t = \frac{R^5}{4 a} \right\} \end{aligned} \quad (3.5)$$

$$\left[\begin{array}{l} \text{hol} > \text{subs}(S[1], \text{LD2}); \\ \quad \quad \quad [[e1, e2] = e3, [e1, e3] = -e2, [e2, e3] = e1] \end{array} \right. \quad (3.6)$$

We conclude that our metric has infinitesimal holonomy given by so(3).

Quaternionic-Kahler Structure

A quaternionic- Kahler structure consists of three covariantly constant $\begin{pmatrix} 1 \\ 1 \end{pmatrix}$ tensors J_1, J_2, J_3 , which satisfy the algebraic relations of the quaternions. Because these tensors are covariantly constant and square to -1, each define a Kahler structure.

It is possible to calculate the tensors J_1, J_2, J_3 directly with the command

CovariantlyConstantTensors but the calculation runs much faster if we first calculate the 1 -1 tensors which are point-wise invariant with respect to the holonomy matrices.

First, [generate](#) all the 1-1 tensors.

$$\left[\begin{array}{l} \text{hol} > \text{OneOneTensors} := \text{GenerateTensors}([\text{zeta0}, \text{zeta1}, \text{zeta2}, \\ \quad \quad \quad \text{zeta3}], [\text{sigma0}, \text{sigma1}, \text{sigma2}, \text{sigma3}]); \\ \text{OneOneTensors} := [\zeta_0 \otimes \sigma_0, \zeta_0 \otimes \sigma_1, \zeta_0 \otimes \sigma_2, \zeta_0 \otimes \sigma_3, \zeta_1 \otimes \sigma_0, \zeta_1 \otimes \sigma_1, \\ \quad \quad \quad \zeta_1 \otimes \sigma_2, \zeta_1 \otimes \sigma_3, \zeta_2 \otimes \sigma_0, \zeta_2 \otimes \sigma_1, \zeta_2 \otimes \sigma_2, \zeta_2 \otimes \sigma_3, \zeta_3 \otimes \sigma_0, \zeta_3 \otimes \sigma_1, \\ \quad \quad \quad \zeta_3 \otimes \sigma_2, \zeta_3 \otimes \sigma_3] \end{array} \right. \quad (4.1)$$

Find the 1-1 tensors which are [invariant](#) with respect to the holonomy.

$$\left[\begin{array}{l} \text{M} > \text{HolInvTensors} := \text{InvariantTensorsAtAPoint}(\text{H}, \text{OneOneTensors}); \\ \text{HolInvTensors} := \left[\zeta_0 \otimes \sigma_0 + \zeta_1 \otimes \sigma_1 + \zeta_2 \otimes \sigma_2 + \zeta_3 \otimes \sigma_3, \frac{R^4 - a}{2 R^3} \zeta_0 \otimes \sigma_1 \right. \\ \quad - \frac{2}{R} \zeta_1 \otimes \sigma_0 - \frac{R^4 - a}{R^4} \zeta_2 \otimes \sigma_3 + \zeta_3 \otimes \sigma_2, - \frac{R^4 - a}{2 R^3} \zeta_0 \otimes \sigma_2 \\ \quad - \frac{R^4 - a}{R^4} \zeta_1 \otimes \sigma_3 + \frac{2}{R} \zeta_2 \otimes \sigma_0 + \zeta_3 \otimes \sigma_1, - \frac{(R^4 - a)^2}{4 R^6} \zeta_0 \otimes \sigma_3 \\ \quad \left. + \frac{R^4 - a}{2 R^3} \zeta_1 \otimes \sigma_2 - \frac{R^4 - a}{2 R^3} \zeta_2 \otimes \sigma_1 + \zeta_3 \otimes \sigma_0 \right] \end{array} \right. \quad (4.2)$$

Find the [covariantly constant tensors](#) in the span of the holonomy invariant tensors.

$$\left[\begin{array}{l} \text{M} > \text{C} := \text{Christoffel}(g); \\ \text{M} > \text{CCOneOneTensors} := \text{CovariantlyConstantTensors}(\text{C}, \\ \quad \quad \quad \text{HolInvTensors}[2 \dots 4]); \end{array} \right. \quad (4.3)$$

$$\begin{aligned}
CCOneOneTensors := & \left[-\frac{R^4 - a}{4R^3} \zeta_0 \otimes \sigma_3 + \frac{1}{2} \zeta_1 \otimes \sigma_2 - \frac{1}{2} \zeta_2 \otimes \sigma_1 \right. \\
& + \frac{R^3}{R^4 - a} \zeta_3 \otimes \sigma_0, -\frac{\sqrt{R^4 - a}}{2R} \zeta_0 \otimes \sigma_2 - \frac{\sqrt{R^4 - a}}{R^2} \zeta_1 \otimes \sigma_3 \\
& + \frac{2R}{\sqrt{R^4 - a}} \zeta_2 \otimes \sigma_0 + \frac{R^2}{\sqrt{R^4 - a}} \zeta_3 \otimes \sigma_1, \frac{\sqrt{R^4 - a}}{2R} \zeta_0 \otimes \sigma_1 \\
& \left. - \frac{2R}{\sqrt{R^4 - a}} \zeta_1 \otimes \sigma_0 - \frac{\sqrt{R^4 - a}}{R^2} \zeta_2 \otimes \sigma_3 + \frac{R^2}{\sqrt{R^4 - a}} \zeta_3 \otimes \sigma_2 \right]
\end{aligned} \tag{4.3}$$

Now scale the results to get tensors with the correct algebraic properties.

$$\begin{aligned}
M > J0 := \text{evalDG}(\text{zeta0} \& \text{t sigma0} + \text{zeta1} \& \text{t sigma1} + \text{zeta2} \& \text{t sigma2} \\
& + \text{zeta3} \& \text{t sigma3}); \\
J0 := & \zeta_0 \otimes \sigma_0 + \zeta_1 \otimes \sigma_1 + \zeta_2 \otimes \sigma_2 + \zeta_3 \otimes \sigma_3
\end{aligned} \tag{4.4}$$

$$\begin{aligned}
M > J1 := 2 \& \text{mult evalDG}(-(\text{R}^4 - \text{a})/(\text{4}*\text{R}^3)*\text{zeta0} \& \text{t sigma3} + \\
& 1/2*\text{zeta1} \& \text{t sigma2} - 1/2*\text{zeta2} \& \text{t sigma1} + \text{R}^3/(\text{R}^4 - \text{a})* \\
& \text{zeta3} \& \text{t sigma0}); \\
J1 := & -\frac{R^4 - a}{2R^3} \zeta_0 \otimes \sigma_3 + \zeta_1 \otimes \sigma_2 - \zeta_2 \otimes \sigma_1 + \frac{2R^3}{R^4 - a} \zeta_3 \otimes \sigma_0
\end{aligned} \tag{4.5}$$

$$\begin{aligned}
M > J2 := (-1) \& \text{mult evalDG}(-\text{sqrt}(\text{R}^4 - \text{a})/(\text{2}*\text{R})*\text{zeta0} \& \text{t sigma2} \\
& - \text{sqrt}(\text{R}^4 - \text{a})/(\text{R}^2)*\text{zeta1} \& \text{t sigma3} + 2*\text{R}/\text{sqrt}(\text{R}^4 - \text{a})*\text{zeta2} \\
& \& \text{t sigma0} + \text{R}^2/\text{sqrt}(\text{R}^4 - \text{a})*\text{zeta3} \& \text{t sigma1}); \\
J2 := & \frac{\sqrt{R^4 - a}}{2R} \zeta_0 \otimes \sigma_2 + \frac{\sqrt{R^4 - a}}{R^2} \zeta_1 \otimes \sigma_3 - \frac{2R}{\sqrt{R^4 - a}} \zeta_2 \otimes \sigma_0 \\
& - \frac{R^2}{\sqrt{R^4 - a}} \zeta_3 \otimes \sigma_1
\end{aligned} \tag{4.6}$$

$$\begin{aligned}
M > J3 := \text{evalDG}(\text{sqrt}(\text{R}^4 - \text{a})/(\text{2}*\text{R})*\text{zeta0} \& \text{t sigma1} - 2*\text{R}/\text{sqrt} \\
& (\text{R}^4 - \text{a})*\text{zeta1} \& \text{t sigma0} - \text{sqrt}(\text{R}^4 - \text{a})/\text{R}^2 * \text{zeta2} \& \text{t sigma3} + \\
& \text{R}^2/\text{sqrt}(\text{R}^4 - \text{a})* \text{zeta3} \& \text{t sigma2}); \\
J3 := & \frac{\sqrt{R^4 - a}}{2R} \zeta_0 \otimes \sigma_1 - \frac{2R}{\sqrt{R^4 - a}} \zeta_1 \otimes \sigma_0 - \frac{\sqrt{R^4 - a}}{R^2} \zeta_2 \otimes \sigma_3 \\
& + \frac{R^2}{\sqrt{R^4 - a}} \zeta_3 \otimes \sigma_2
\end{aligned} \tag{4.7}$$

$$\begin{aligned}
M > QTensors := [J0, J1, J2, J3]; \\
QTensors := & \left[\zeta_0 \otimes \sigma_0 + \zeta_1 \otimes \sigma_1 + \zeta_2 \otimes \sigma_2 + \zeta_3 \otimes \sigma_3, -\frac{R^4 - a}{2R^3} \zeta_0 \otimes \sigma_3 \right.
\end{aligned} \tag{4.8}$$

$$\begin{aligned}
& + \zeta_1 \otimes \sigma_2 - \zeta_2 \otimes \sigma_1 + \frac{2R^3}{R^4 - a} \zeta_3 \otimes \sigma_0, \frac{\sqrt{R^4 - a}}{2R} \zeta_0 \otimes \sigma_2 \\
& + \frac{\sqrt{R^4 - a}}{R^2} \zeta_1 \otimes \sigma_3 - \frac{2R}{\sqrt{R^4 - a}} \zeta_2 \otimes \sigma_0 - \frac{R^2}{\sqrt{R^4 - a}} \zeta_3 \otimes \sigma_1, \\
& \frac{\sqrt{R^4 - a}}{2R} \zeta_0 \otimes \sigma_1 - \frac{2R}{\sqrt{R^4 - a}} \zeta_1 \otimes \sigma_0 - \frac{\sqrt{R^4 - a}}{R^2} \zeta_2 \otimes \sigma_3 \\
& + \frac{R^2}{\sqrt{R^4 - a}} \zeta_3 \otimes \sigma_2 \Big]
\end{aligned}$$

Check that all our 1-1 tensors are parallel.

$$\begin{aligned}
\mathbf{M} & > \text{map}(\text{CovariantDerivative}, \text{QTensors}, \mathbf{C}); \\
& [0 \zeta_0 \otimes \sigma_0 \otimes \sigma_0, 0 \zeta_0 \otimes \sigma_0 \otimes \sigma_0, 0 \zeta_0 \otimes \sigma_0 \otimes \sigma_0, 0 \zeta_0 \otimes \sigma_0 \otimes \sigma_0] \quad (4.9)
\end{aligned}$$

$$\mathbf{M} > \text{QMatrices} := \text{map}(\text{convert}, \text{map}(\text{convert}, \text{QTensors}, \text{DGArray}), \text{Matrix});$$

$$\text{QMatrices} := \left[\begin{array}{cccc} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{array} \right], \left[\begin{array}{cccc} 0 & 0 & 0 & -\frac{R^4 - a}{2R^3} \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ \frac{2R^3}{R^4 - a} & 0 & 0 & 0 \end{array} \right], \quad (4.10)$$

$$\left[\begin{array}{cccc} 0 & 0 & \frac{\sqrt{R^4 - a}}{2R} & 0 \\ 0 & 0 & 0 & \frac{\sqrt{R^4 - a}}{R^2} \\ -\frac{2R}{\sqrt{R^4 - a}} & 0 & 0 & 0 \\ 0 & -\frac{R^2}{\sqrt{R^4 - a}} & 0 & 0 \end{array} \right],$$

$$\begin{bmatrix} 0 & \frac{\sqrt{R^4 - a}}{2R} & 0 & 0 \\ -\frac{2R}{\sqrt{R^4 - a}} & 0 & 0 & 0 \\ 0 & 0 & 0 & -\frac{\sqrt{R^4 - a}}{R^2} \\ 0 & 0 & \frac{R^2}{\sqrt{R^4 - a}} & 0 \end{bmatrix}$$

Show that these 4 matrices define an [algebra](#) using ordinary matrix multiplication.

```
M > simplify(AlgebraData(QMatrices, proc(A,B) A.B end, alg))
      assuming R^4 - a > 0;
```

$$\begin{aligned} [e1^2 = e1, e1 \cdot e2 = e2, e1 \cdot e3 = e3, e1 \cdot e4 = e4, e2 \cdot e1 = e2, e2^2 = -e1, e2 \cdot e3 = e4, e2 \cdot e4 \\ = -e3, e3 \cdot e1 = e3, e3 \cdot e2 = -e4, e3^2 = -e1, e3 \cdot e4 = e2, e4 \cdot e1 = e4, e4 \cdot e2 = e3, e4 \\ \cdot e3 = -e2, e4^2 = -e1] \end{aligned} \quad (4.11)$$

The multiplication table matches exactly that of the quaternions!

```
M > AlgebraLibraryData("Quaternions", q);
```

$$\begin{aligned} [e1^2 = e1, e1 \cdot e2 = e2, e1 \cdot e3 = e3, e1 \cdot e4 = e4, e2 \cdot e1 = e2, e2^2 = -e1, e2 \cdot e3 = e4, e2 \cdot e4 \\ = -e3, e3 \cdot e1 = e3, e3 \cdot e2 = -e4, e3^2 = -e1, e3 \cdot e4 = e2, e4 \cdot e1 = e4, e4 \cdot e2 = e3, e4 \\ \cdot e3 = -e2, e4^2 = -e1] \end{aligned} \quad (4.12)$$

Kahler Metric, Kahler Potential, and Holomorphic Sectional Curvature

In this section we shall analyze the complex structure defined by the 1-1 tensor J1. First we calculate the eigenforms for this tensor and use these to define a new frame.

```
M > lambda, P := DGEigenTensors(J1, [sigma0, sigma1, sigma2,
      sigma3]);
```

$$\lambda, P := \begin{bmatrix} I \\ I \\ -I \\ -I \end{bmatrix}, \left[-\frac{2IR^3}{R^4 - a} \sigma_0 + \sigma_3, I\sigma_1 + \sigma_2, \frac{2IR^3}{R^4 - a} \sigma_0 + \sigma_3, -I\sigma_1 + \sigma_2 \right] \quad (5.1)$$

The 1st and 3rd forms are complex conjugates. The 2nd and 4th forms are complex conjugates.

$$\begin{aligned} \mathbf{N} > \text{DGconjugate}(\mathbf{P}[1]) \text{ \− } \mathbf{P}[3]; \\ & \quad 0 \sigma 0 \end{aligned} \quad (5.2)$$

$$\begin{aligned} \mathbf{M} > \text{DGconjugate}(\mathbf{P}[2]) \text{ \− } \mathbf{P}[4]; \\ & \quad 0 \sigma 0 \end{aligned} \quad (5.3)$$

$$\begin{aligned} \mathbf{M} > \text{FD2} := \text{FrameData}(\mathbf{P}, \mathbf{N}); \\ \text{FD2} := \left[d\Theta 1 = \frac{1}{2} \Theta 2 \wedge \Theta 4, d\Theta 2 = \frac{1 \Theta 1 \wedge \Theta 2}{2} - \frac{1 \Theta 2 \wedge \Theta 3}{2}, d\Theta 3 \right. \\ \left. = \frac{1}{2} \Theta 2 \wedge \Theta 4, d\Theta 4 = -\frac{1 \Theta 1 \wedge \Theta 4}{2} - \frac{1 \Theta 3 \wedge \Theta 4}{2} \right] \end{aligned} \quad (5.4)$$

$$\begin{aligned} \mathbf{M} > \text{DGsetup}(\text{FD2}, '[\mathbf{U1}, \mathbf{U2}, \mathbf{V1}, \mathbf{V2}]', '[\text{omega1}, \text{omega2}, \text{theta1}, \\ \text{theta2}]'); \\ & \quad \text{frame name: } \mathbf{N} \end{aligned} \quad (5.5)$$

Now we transform the metric and the complex structure J1 to the new frame.

$$\begin{aligned} \mathbf{N} > \text{IdMtoN} := \text{Transformation}(\mathbf{M}, \mathbf{N}, [\mathbf{R} = \mathbf{R}, \text{theta} = \text{theta}, \text{phi} = \\ \text{phi}, \text{psi} = \text{psi}]); \\ \text{IdMtoN} := [\mathbf{R} = \mathbf{R}, \theta = \theta, \phi = \phi, \psi = \psi] \end{aligned} \quad (5.6)$$

$$\begin{aligned} \mathbf{M} > \text{IdNtoM} := \text{Transformation}(\mathbf{N}, \mathbf{M}, [\mathbf{R} = \mathbf{R}, \text{theta} = \text{theta}, \text{phi} = \\ \text{phi}, \text{psi} = \text{psi}]); \\ \text{IdNtoM} := [\mathbf{R} = \mathbf{R}, \theta = \theta, \phi = \phi, \psi = \psi] \end{aligned} \quad (5.7)$$

Here in the Eguchi-Hanson metric in the new frame:

$$\begin{aligned} \mathbf{N} > \mathbf{gN} := \text{Pullback}(\text{IdNtoM}, \mathbf{g}); \\ \mathbf{gN} := \frac{R^4 - a}{8 R^2} \omega 1 \otimes \theta 1 + \frac{R^2}{8} \omega 2 \otimes \theta 2 + \frac{R^4 - a}{8 R^2} \theta 1 \otimes \omega 1 + \frac{R^2}{8} \theta 2 \otimes \omega 2 \end{aligned} \quad (5.8)$$

$$\begin{aligned} \mathbf{N} > \mathbf{JN} := \text{PushPullTensor}(\text{IdMtoN}, \text{IdNtoM}, \mathbf{J1}); \\ \mathbf{JN} := 1 \mathbf{U1} \otimes \omega 1 + 1 \mathbf{U2} \otimes \omega 2 - 1 \mathbf{V1} \otimes \theta 1 - 1 \mathbf{V2} \otimes \theta 2 \end{aligned} \quad (5.9)$$

Here is our complex structure.

$$\begin{aligned} \mathbf{N} > \mathbf{JN} := (1/\mathbf{I}) \text{ \&mult; } \mathbf{JN}; \\ \mathbf{JN} := \mathbf{U1} \otimes \omega 1 + \mathbf{U2} \otimes \omega 2 - \mathbf{V1} \otimes \theta 1 - \mathbf{V2} \otimes \theta 2 \end{aligned} \quad (5.10)$$

Here is the [Kahler form](#).

$$\begin{aligned} \mathbf{N} > \mathbf{Omega} := \text{KahlerForm}(\mathbf{gN}, \mathbf{JN}); \\ \mathbf{Omega} := \frac{R^4 - a}{8 R^2} \omega 1 \wedge \theta 1 + \frac{R^2}{8} \omega 2 \wedge \theta 2 \end{aligned} \quad (5.11)$$

This form is closed.

$$\begin{aligned} \mathbf{N} > \text{ExteriorDerivative}(\Omega); \\ & 0 \omega_1 \wedge \omega_2 \wedge \theta_1 \end{aligned} \quad (5.12)$$

Here is the [Kahler potential](#) for the metric.

$$\begin{aligned} \mathbf{N} > \text{rho} := \text{KahlerPotential}(\Omega, \mathbf{JN}, \text{ansatz} = \mathbf{F}(\mathbf{R})); \\ \rho := 2R^2 - 2\sqrt{a} \operatorname{arctanh}\left(\frac{R^2}{\sqrt{a}}\right) + _CI \end{aligned} \quad (5.13)$$

$$\begin{aligned} \mathbf{N} > \text{alpha} := \text{DolbeaultExteriorDerivative}(\text{rho}, [0, 1], \mathbf{JN}); \\ \alpha := \left(\frac{R^2}{2} + \frac{I R^2}{2}\right) \omega_1 + \left(\frac{R^2}{2} - \frac{I R^2}{2}\right) \theta_1 \end{aligned} \quad (5.14)$$

$$\begin{aligned} \mathbf{N} > \text{beta} := \text{DolbeaultExteriorDerivative}(\text{alpha}, [1, 0], \mathbf{JN}); \\ \beta := -\frac{\frac{1}{4}(R^4 - a)}{R^2} \omega_1 \wedge \theta_1 - \frac{1}{4} R^2 \omega_2 \wedge \theta_2 \end{aligned} \quad (5.15)$$

$$\begin{aligned} \mathbf{N} > \text{evalDG}(I/2 * \text{beta} - \Omega); \\ & 0 \omega_1 \wedge \omega_2 \end{aligned} \quad (5.16)$$

Finally, here is the holomorphic sectional curvature:

$$\begin{aligned} \mathbf{N} > \mathbf{C} := \text{CurvatureTensor}(\mathbf{gN}); \\ \mathbf{N} > \text{HolomorphicSectionalCurvature}(\mathbf{gN}, \mathbf{C}, \mathbf{U1} + \mathbf{V1}, \mathbf{JN}); \\ & -\frac{a(R^4 - a)}{R^8} \end{aligned} \quad (5.17)$$

Complex Coordinates

$$\begin{aligned} \mathbf{N} > \text{ExteriorDifferentialSystems}:-\text{FirstIntegrals}([\text{omega1}, \text{omega2}]) \\ ; \\ \left[I \ln\left(\frac{1 - \cos(\theta)}{\sin(\theta)}\right) + \phi, -I \ln(\sin(\theta)) - \frac{I \ln(R^4 - a)}{2} + \psi \right] \end{aligned} \quad (6.1)$$

$$\begin{aligned} \mathbf{N} > \text{ExteriorDifferentialSystems}:-\text{FirstIntegrals}([\text{theta1}, \text{theta2}]) \\ ; \\ \left[-I \ln\left(\frac{1 - \cos(\theta)}{\sin(\theta)}\right) + \phi, I \ln(\sin(\theta)) + \frac{I \ln(R^4 - a)}{2} + \psi \right] \end{aligned} \quad (6.2)$$